## Hadronic Light by Light Contribution to Muon $g$ - 2

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## . Introduction

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, "Old" Calculations: 1995-2001
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- New Short-Distance Constraints: 2003-2004
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-Comparison
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Hadronic light-by-light contribution to muon $g-2$


$$
\begin{gathered}
\mathcal{M}=|e|^{7} A_{\beta} \int \frac{\mathrm{d}^{4} p_{1}}{(2 \pi)^{4}} \int \frac{\mathrm{~d}^{4} p_{2}}{(2 \pi)^{4}} \frac{1}{q^{2} p_{1}^{2} p_{2}^{2}\left(p_{4}^{2}-m^{2}\right)\left(p_{5}^{2}-m^{2}\right)} \\
\times \underline{\Pi^{\rho \nu \alpha \beta}\left(p_{1}, p_{2}, p_{3}\right)}
\end{gathered}
$$

Need

$$
\begin{aligned}
\Pi^{\rho \nu \alpha \beta}\left(p_{1}, p_{2}, p_{3}\right)=i^{3} \int & \mathrm{~d}^{4} x \int
\end{aligned} \mathrm{~d}^{4} y \int \mathrm{~d}^{4} z \exp ^{i\left(p_{1} \cdot x+p_{2} \cdot y+p_{3} \cdot z\right)} \times
$$

with $V^{\mu}(x)=\left[\bar{q} \widehat{Q} \gamma^{\mu} q\right](x)$ and $\widehat{Q}=\frac{1}{3} \operatorname{diag}(2,-1,-1)$
full four-point function with $p_{3} \rightarrow 0 \bullet$
Using gauge invariance

$$
\Pi^{\rho \nu \alpha \lambda}\left(p_{1}, p_{2}, p_{3}\right)=-p_{3 \beta} \frac{\delta \Pi^{\rho \nu \alpha \beta}\left(p_{1}, p_{2}, p_{3}\right)}{\delta p_{3 \lambda}}
$$

one just needs derivatives at $p_{3}=0 \bullet$

Ł Many scales involved: Impose low energy and several OPE limits $\Rightarrow$ Not full first principle calculation at present $\bullet$

Large $N_{c}$ and CHPT counting:
Organizes different degrees of freedom contributions •
E. de Rafael

- Goldstone boson exchange: $\mathcal{O}\left(N_{c}\right)$ and $\mathcal{O}\left(p^{6}\right)$ •
- Quark Loop and non-Goldstone boson exchange: $\mathcal{O}\left(N_{c}\right)$ and $\mathcal{O}\left(p^{8}\right)$ •
- Goldstone bosons Loop: $\mathcal{O}(1)$ in $1 / N_{c}$ and $\mathcal{O}\left(p^{4}\right)$ -


## Based on this counting:

- Two full calculations
J. Bijnens, E. Pallante, J.P. (BPP)
M. Hayakawa, T. Kinoshita, A. Sanda (HKS)
- Dominant pseudo-scalar exchange: Extensive analytic analysis •
M. Knecht, A. Nyffeler (KN)

Found sign mistake $\sqrt{ }$
M. Knecht, A. Nyffeler, M. Perrottet, E. de Rafael
$\star$ New four-point form factor short-distance constraint:
K. Melnikov, A. Vainshtein
(see also M. Knecht, S. Peris, M. Perrottet, E. de Rafael)

Model: Full light-by-light saturated by pseudo-scalar and pseudo-vector pole exchanges •

## Dominant contribution $\Rightarrow$ pseudo-scalar exchange $\bullet$



Here, I discuss work in J. Bijnens, E. Pallante, J.P. •

We used a variety of $\pi^{0} \gamma^{*} \gamma^{*}$ form factors

$$
\mathcal{F}^{\mu \nu}\left(p_{1}, p_{2}\right)=\frac{N_{c}}{6 \pi} \frac{\alpha}{f_{\pi}} i \varepsilon^{\mu \nu \alpha \beta} p_{1 \alpha} p_{2 \beta} \underline{\mathcal{F}\left(p_{1}^{2}, p_{2}^{2}\right)}
$$

fulfilling as many as possible QCD constraints • (Short-distance, data, $\mathrm{U}_{A}(1)$ normalization and slope at the origin). In particular,

$$
\begin{aligned}
\mathcal{F}\left(Q^{2}, Q^{2}\right) & \rightarrow \frac{A}{Q^{2}} \\
\mathcal{F}\left(Q^{2}, 0\right) & \rightarrow \frac{B}{Q^{2}}
\end{aligned}
$$

for $Q^{2}$ Euclidean and very large

All form factors we used converge for $\mu \sim(2-4) \mathrm{GeV}$ and the numerical difference between them is small $\sqrt{ }$

Somewhat different $\pi^{0} \gamma^{*} \gamma^{*}$ form factors used in M. Hayakawa, T. Kinoshita, A. Sanda and M. Knecht, A. Nyffeler •

Results agree very well (after correcting a mistake in the sign of the phase space)

Adding $\pi^{0}, \eta$ and $\eta^{\prime}$ contributions

|  | $10^{10} \times a_{\mu}$ |
| :---: | :---: |
| BPP | $(8.5 \pm 1.3)$ |
| HKS | $(8.3 \pm 0.6)$ |
| KN | $(8.3 \pm 1.2)$ |

Need $a_{1}^{0} \gamma \gamma^{*}$ and $a_{1}^{0} \gamma^{*} \gamma^{*}$ form factors •
$\Rightarrow$ related to $\pi^{0} \gamma \gamma^{*}$ and $\pi^{0} \gamma^{*} \gamma^{*}$ by anomalous Ward identities $\sqrt{ }$

## Pseudo-vector exchange

|  | $10^{10} \times a_{\mu}$ |
| :---: | :---: |
| BPP | $(0.25 \pm 0.10)$ |
| HKS | $(0.17 \pm 0.10)$ |

Need $S^{0} \gamma \gamma^{*}$ and $S^{0} \gamma^{*} \gamma^{*}$ form factors •
They are constrained by CHPT at $\mathcal{O}\left(p^{4}\right)$ : $L_{i}$ 's reproduced
Within ENJL: Ward identities impose relations between Quark loop and Scalar exchange •

$$
a_{\mu}(\text { Scalar })=-(0.7 \pm 0.2) \cdot 10^{-10}
$$

Not included by M. Hayakawa, T. Kinoshita and A. Sanda nor by K. Melnikov and A. Vainshtein •


| $\Lambda[\mathrm{GeV}]$ | $10^{10} \times a_{\mu}$ |
| :---: | :---: |
| 0.7 | 2.2 |
| 1.0 | 2.0 |
| 2.0 | 1.9 |
| 4.0 | 2.0 |

- Low Energy ( 0 to $\Lambda$ ): ENJL model •
- High Energy $(\Lambda$ to $\infty)$ : Bare heavy quark loop with $m_{Q}=\Lambda$ -
- Numerical matching $\sqrt{ }$

Leading contribution in chiral counting, suppressed by $1 / N_{c}$


No $\gamma^{*} \gamma^{*} \rightarrow \pi \pi$ data available: Models needed!

| Model for $\pi \pi \gamma(\gamma)$ | $10^{10} \times a_{\mu}$ |
| ---: | :---: |
| BPP (Full VMD) | -1.8 |
| HKS (HGS) | -0.4 |

Kaon loop is much smaller: $-0.05 \times 10^{-10} \bullet$
K. Melnikov and A. Vainshtein

New short-distance constraint on four-point function form factor

$$
\langle 0| T\left[V^{\nu}\left(p_{1}\right) V^{\alpha}\left(p_{2}\right) V^{\rho}\left(-\left(p_{1}+p_{2}+p_{3}\right)\right)\right]\left|\gamma\left(p_{3} \rightarrow 0\right)\right\rangle
$$

using $\underline{\text { OPE with }}-p_{1}^{2} \simeq-p_{2}^{2} \gg-\left(p_{1}+p_{2}\right)^{2}$ Euclidean and large,

$$
T\left[V^{\nu}\left(p_{1}\right) V^{\alpha}\left(p_{2}\right)\right] \sim \frac{1}{\hat{p}^{2}} \varepsilon^{\nu \alpha \mu \beta} \hat{p}_{\mu}\left[\bar{q} \widehat{Q}^{2} \gamma_{\beta} \gamma_{5} q\right]\left(p_{1}+p_{2}\right)
$$

with $\hat{p}=\left(p_{1}-p_{2}\right) / 2 \simeq p_{1} \simeq-p_{2}$

## New OPE constraint saturated by pseudo-scalar exchange

 $\Rightarrow$ Model uses a point-like vertex when $p_{3} \rightarrow 0 \bullet$Not all OPE constraints satisfied: Negligible numerically •


## Axial-Vector exchange depends very much on the resonance mass mixing •

K. Melnikov and A. Vainsthein: Ideal mixing for $f_{1}(1285)$ and $f_{1}(1420)$

| Mass mixing | $10^{10} \times a_{\mu}$ |
| :---: | :---: |
| No New OPE (Nonet symmetry) | $0.3 \pm 0.1$ |
| $\mathrm{M}=1.3 \mathrm{GeV}$ (Nonet symmetry) | 0.7 |
| $\mathrm{M}=\mathrm{M}_{\rho}$ (Nonet symmetry) | 2.8 |
| Ideal mixing | $2.2 \pm 0.5$ |

Leading order in $N_{c}$ contributions:
Quark Loop + Pseudo-Scalar + Pseudo-Vector + Scalar Exchanges •

| Total at $\mathcal{O}\left(N_{c}\right)$ | $10^{10} \times a_{\mu}$ |
| :---: | :---: |
| BPP (Nonet symmetry) | $\frac{(10.9 \pm 1.9)}{-(0.7 \pm 0.1)}=(10.2 \pm 1.9)$ |
| HKS (Nonet symmetry) | $\underline{(9.4 \pm 1.6)+? ? S c a l a r ? ? ~}$ |

MV: Hadronic model saturated by pole exchanges:
Cannot compare individual contributions •

| Total at $\mathcal{O}\left(N_{c}\right)$ | $10^{10} \times a_{\mu}$ |
| :---: | :---: |
| MV (Nonet symmetry) | $\underline{(12.1 \pm 1.0)}+? ?$ Scalar?? |
| MV (Ideal mass mixing) | $\underline{(13.6 \pm 1.5)}+? ?$ Scalar?? |

Masses produce main difference in pseudo-vector exchange •

## Study of momenta regions contribution for $\pi^{0}$ exchange

$$
\begin{aligned}
& a_{\mu}^{\mathrm{lbl}}=\int \mathrm{d} P_{1} \mathrm{~d} P_{2} a_{\mu}^{P P}\left(P_{1}, P_{2}\right) \\
& \quad=\int \mathrm{d} l_{1} \mathrm{~d} l_{2} a_{\mu}^{L L}\left(l_{1}, l_{2}\right) \\
& \quad=\int \mathrm{d} l_{1} \mathrm{~d} l_{2} \mathrm{~d} q a_{\mu}^{P P Q}\left(l_{1}, l_{2}, q\right)
\end{aligned}
$$

with $l_{1} \equiv \ln \left(P_{1} / \mathrm{GeV}\right), l_{2} \equiv \ln \left(P_{2} / \mathrm{GeV}\right)$ and $q \equiv \ln (Q / \mathrm{GeV})$
$P_{1}^{2}=-p_{1}^{2}, \quad P_{2}^{2}=-p_{2}^{2}, \quad Q^{2}=-\left(p_{1}+p_{2}\right)^{2}$





Conclusions of this comparison
Cut-off in $Q=\varnothing 58 \%$ of numerical difference come from MV OPE violating regions $\bullet$

Fixing $P 1=P 2 \leftrightharpoons$ Numerical difference come from low values of $Q$ and moderate values of $P_{1}=P_{2} \bullet$

Important to control energy regions below 2 GeV •
Main ENJL quark-loop contribution is from that region •

Next to leading order in $1 / N_{c}$ contributions:

## Charged Pion and Kaon Loop •

| Model for $\pi \pi \gamma(\gamma)$ | $10^{10} \times a_{\mu}$ |
| ---: | :---: |
| BPP (Full VMD) | $-1.9 \pm 0.5$ |
| HKS (HGS) | $-0.45 \pm 0.8$ |

K. Melnikov and A. Vainshtein:

Full NLO in $1 / N_{c}$ estimate

$$
a_{\mu}=(0 \pm 1) \cdot 10^{-10}
$$

BPP vs HKS:

| Full | $10^{10} \times a_{\mu}$ |
| ---: | :--- |
| BPP | $8.3 \pm 3.2$ |
| HKS | $8.9 \pm 1.7$ |

No scalar exchange, different quark loop and different pion and kaon loops almost compensate •

BPP vs MV:

| Full | $10^{10} \times a_{\mu}$ |
| ---: | :---: |
| BPP | $8.3 \pm 3.2$ |
| MV | $13.6 \pm 2.5$ |

Several order $1.5 \cdot 10^{-10}$ differences, in addition to new OPE effects -
$-1.5 \cdot 10^{-10}$ (Different pseudo-vector mass mixing)
$-0.7 \cdot 10^{-10}$ (No scalar exchange)
$-1.9 \cdot 10^{-10}$ (No pion+kaon loop)
$=-4.1 \cdot 10^{-10}$
Final [BPP-MV] difference: $\underline{-5.3 \cdot 10^{-10}}$

Unsatisfactory situation: Needs new evaluation(s) of the full hadronic light-by-light contribution •
$\star$ At $\mathcal{O}\left(N_{c}\right)$ : Study full four-point function with large $N_{c}$ techniques • Granada-Lund-València

- Implement as many short-distance and low energy constraints as possible •
(possible problems J. Bijnens, E. Gámiz, E. Lipartia, J.P.)
$\star$ At NLO in $1 / N_{c} \rightleftharpoons$ Non-Goldstone bosons at one loop • Little is known (see recent work by A. Pich, I. Rosell, J. Sanz-Cillero) •

At present,
Large $N_{c}$ agree within $1 \sigma \Rightarrow$

$$
a_{\mu}^{\mathrm{lbl}}=(11.0 \pm 4.0) \times 10^{-10}
$$

More work needed to have a definite answer of hadronic light-by-light contribution to muon $g-2$ with reduced uncertainty

Goal: Control present model dependences •

