



# Topology-changing transitions with gravity

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# Talk outline

- moving around moduli space
- topological transitions and their consequences
- including gravity in cycle-collapses
- results
- conclusions

## Moduli space motion

Higher dimensional perspective:

$$ds^2 = \mathrm{d}s_{us}^2 + a^2(x)\mathrm{d}\theta^2 + b^2(x)\mathrm{d}\phi^2 + \dots$$

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Higher dimensional perspective:

$$ds^{2} = ds_{us}^{2} + a^{2}(x)d\theta^{2} + b^{2}(x)d\phi^{2} + \dots$$

Our perspective:

$$\mathcal{L} = R - \frac{1}{2} \partial_{\mu} a \partial^{\mu} a - \frac{1}{2} \partial_{\mu} b \partial^{\mu} b - \dots$$

## Moduli space

Motion with constant topology:

- Betti numbers unchanged
- intersection numbers unchanged



## Moduli space

Motion with changing topology:

Betti numbers changede.g. conifold transition





## **Conifold transition**

$$ds^{2} = dr^{2} + r^{2} \left( d\Omega_{(2)} + d\Omega_{(3)} \right)$$



## **Regulating the transition**



wrapped D-branes provide extra massless states which alter the effective theory

Strominger, Greene, Morrison, Vafa

## Effective action

Use knowledge of effective theory to include extra degrees of freedom for the wrapped branes

For the conifold in IIB we find 4D N=2 sugra, with one hypermultiplet for each degenerating cycle.

$$= \mathcal{R} \star 1 - g_{i\bar{j}} dz^{i} \wedge \star dz^{\bar{j}} - h_{uv} Dq^{u} \wedge Dq^{v} \\ + \frac{1}{2} Im(\mathcal{N})_{IJ} F^{I} \wedge \star F^{J} - \frac{1}{2} Re(\mathcal{N})_{IJ} F^{I} \wedge F^{J} - \mathcal{V}(q, z)$$

Brandle, Greene, Jarv, Lukas, Mohaupt, Morrison, Palti, Saffin, Saueressig, Vafa

L

### **Dynamical consequences**

#### Trapping at the transition point



U

W

q

40 time 60

20

0.6

0.4

0.2

0

IIB conifold: Lukas, Palti, Saffin

#### M-theory flop: Lukas, Brandle



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- the cycle continues to collapse, classical geometry requires the formation of a curvature singularity
  - does this create a naked singularity?
  - is a horizon formed?
  - what is the final state?

$$\begin{aligned} & \text{Cycle collapse - a simpler model} \\ & ds^2 = \alpha^{-1}(\rho)d\rho^2 + \frac{1}{4}\rho^2 \left[ (\sigma_1^2 + \sigma_2^2) + \alpha(\rho)\sigma_3^2 \right] \\ & \alpha(\rho) = 1 - \left(\frac{L}{\rho}\right)^4 \qquad \text{Eguchi, Hanson} \\ & \text{two-cycle:} \quad \frac{\rho = L + \frac{R^2}{L}}{ds^2 \simeq dR^2 + R^2\sigma_3^2 + \frac{L^2}{4} \left(\sigma_1^2 + \sigma_2^2\right)} \end{aligned}$$



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apparent horizon:

 $\left. \frac{d}{dt} (area) \right|_{null\ ray} = 0$ 

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Final state?

 $ds^{2} = -fdt^{2} + \frac{k^{2}}{f}dr^{2} + \frac{r^{2}}{4}\left[k(\sigma_{1}^{2} + \sigma_{2}^{2}) + \sigma_{3}^{2}\right]$  $f(r) = \frac{r^2 - r_+^2}{r^2}$  $k(r) = \frac{(r_{\infty}^2 - r_{+}^2)r_{\infty}^2}{(r_{\infty}^2 - r^2)^2}$ 

Gibbons, Maeda, Ishihara, Matsuno



### <u>conclusions</u>

- branes can regulate the certain singularities
- can derive effective actions for topology changes using the moduli space approximation
- the moduli space approximation breaks down as cycles collapse
- event horizons can form from cycles collapsing
- cycle collapses give a novel form of dynamical compactification

## Moduli space

- Betti numbers unchanged
- Intersection numbers changed
- e.g. flop transition





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