

SU(3) gauge theories with many massless fermions: methods and mysteries

David Schaich (University of Colorado)

Lattice Meets Experiment: Beyond the Standard Model
Boulder, 27 October 2012

[arXiv:1207.7162](#), [arXiv:1207.7164](#) and work in progress
with Anqi Cheng, Anna Hasenfratz and Gregory Petropoulos



Thank you!

For your participation

Contributions to the workshop, white paper and community planning

For your patience

As I say a few words about a small portion of our ongoing explorations

Thank you!

For your participation

Contributions to the workshop, white paper and community planning

For your patience

As I say a few words about a small portion of our ongoing explorations

$N_F = 8$ widely believed to be *well below* conformal window

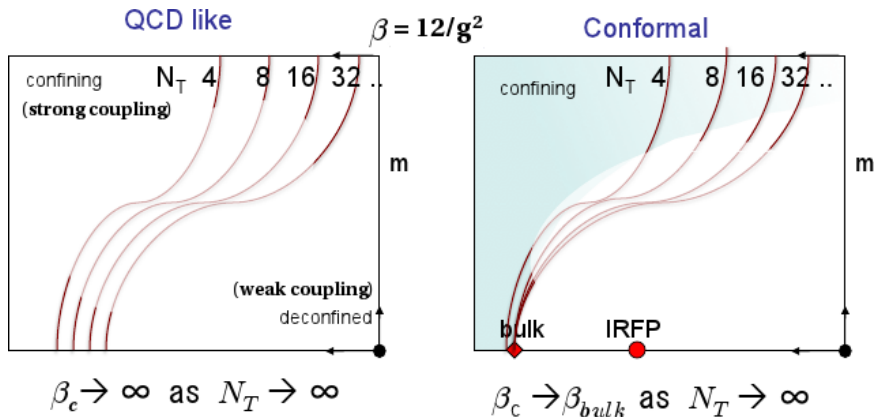
We are currently investigating $N_f = 8$ systems **directly at zero mass**

On this critical surface:

- 1 Finite-temperature phase transitions don't exhibit QCD-like scaling
- 2 Eigenvalues provide access to scale-dependence of $\gamma_m(\mu)$
which also shows clearly non-QCD behavior for $N_F = 8$

Qualitative expectations for the lattice phase diagram

Fermion mass vs. gauge coupling – critical surface is $m \rightarrow 0$ chiral limit

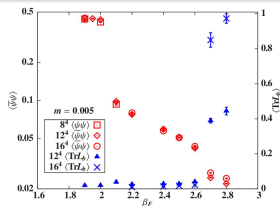
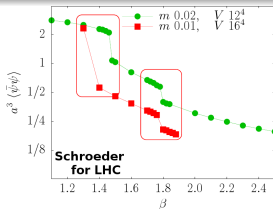
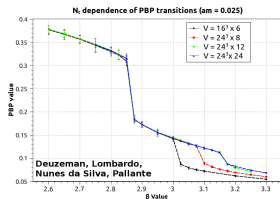


Hope for clear distinction between QCD-like and conformal cases from scaling $\Delta\beta$ of finite-temperature transitions as N_T increases

Of course, it's not that simple

Several groups find novel **intermediate** phase

(for $N_F = 12$)

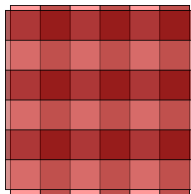
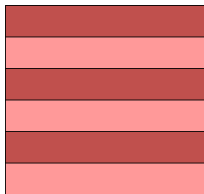


Suspect intermediate phase has no continuum limit

Below I refer to our observation

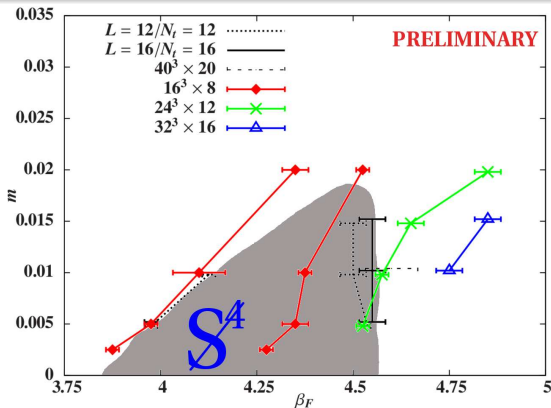
(for $N_F = 8$ and 12)

that it exhibits spontaneous **single-site shift symmetry breaking** (" \mathcal{S}^4 ")



$N_F = 8$ at non-zero mass (as of June)

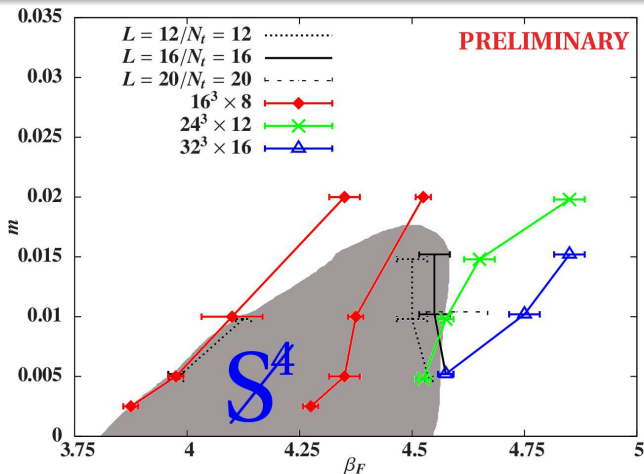
$T > 0$ transitions pass through bulk transitions surrounding \mathcal{S}^4 phase



Observe chiral symmetry breaking in systems
 between \mathcal{S}^4 phase and deconfinement at weak coupling (large β_F)
 $\Delta\beta_F$ agrees with two-loop prediction $\Delta\beta_F \approx 0.25$ for $N_T = 12 \rightarrow 16$

$N_F = 8$ at non-zero mass (as of July)

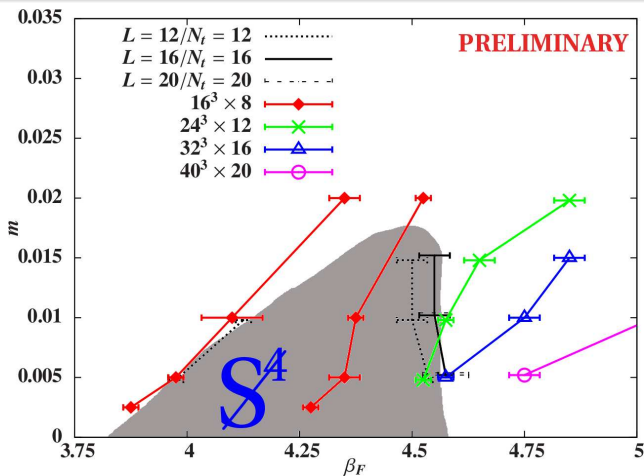
$N_T = 16$ transitions run into the S^4 phase at $m = 0.005$



We lose scaling with $N_T = 12 \rightarrow 16$ as we approach the chiral limit

$N_F = 8$ at non-zero mass

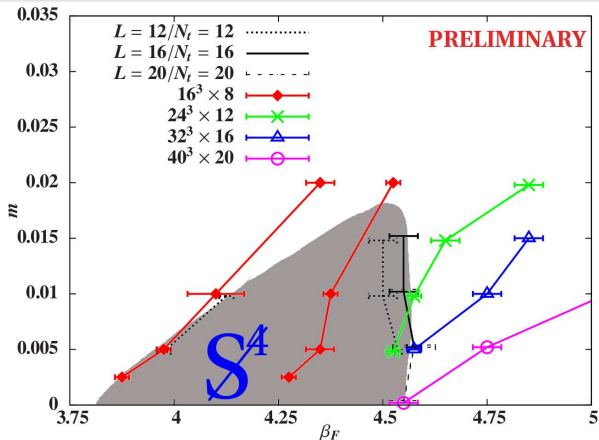
At $m \gtrsim 0.005$, still have scaling with $N_T = 16 \rightarrow 20$



Two-loop prediction is now $\Delta\beta_F \approx 0.2$ for $N_T = 16 \rightarrow 20$

$N_F = 8$ at zero mass

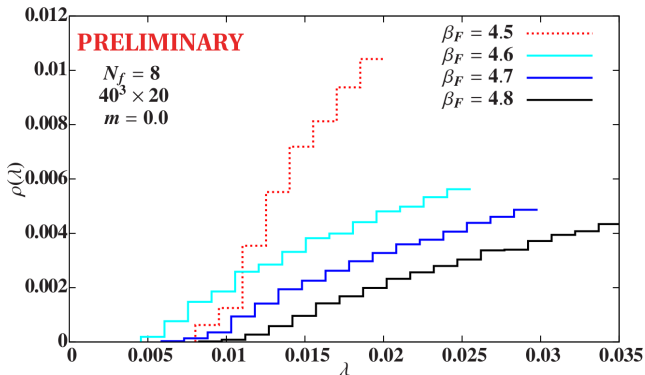
We can study $40^3 \times 20$ (and $24^3 \times 48$) **directly** at $m = 0$
on both sides of transition into S^4 phase



Again lose scaling with N_T , on the $m = 0$ critical surface

Closer look at $40^3 \times 20$ transition with $m = 0$

Eigenvalue densities (histogram) $\rho(\lambda)$ of **massless** Dirac operator
Strong couplings produce smaller λ , until we hit the \mathcal{S}^4 phase
None of these systems are chirally broken : $\langle \bar{\psi}\psi \rangle \propto \rho(0) = 0$



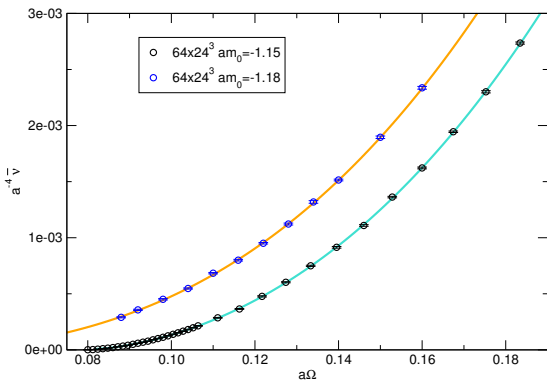
In addition to $\langle \bar{\psi}\psi \rangle$, $\rho(\lambda)$ also related to anomalous dimension γ_m

γ_m from eigenvalue mode number $\nu(\lambda)$

Del Debbio & Zwicky, arXiv:1005.2371

In the chiral limit $\rho(\lambda) \sim \lambda^\alpha \implies \nu(\lambda) = V \int_{-\lambda}^{\lambda} \rho(\omega) d\omega \sim V\lambda^{1+\alpha}$

RG invariance of **mode number** $\nu(\lambda) \implies 1 + \gamma_m = \frac{4}{1 + \alpha}$



Patella, arXiv:1204.4432

SU(2), $N_F = 2$ adjoint
believed IR-conformal

$\gamma_\star = 0.371(20)$
for fit range [0.091, 0.18]

Inspired us to look at $\nu(\lambda)$

Expectations for eigenvalue mode number analysis

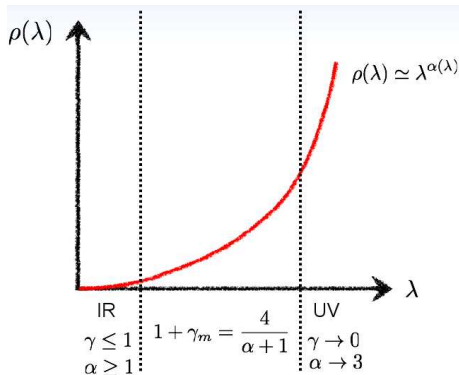
λ defines an energy scale;

fitting $\nu(\lambda) \propto \lambda^{1+\alpha(\lambda)}$ accesses $1 + \gamma_m(\lambda) = \frac{4}{1+\alpha(\lambda)}$ at that scale

For IR-conformal systems:

UV: Asymp. freedom $\Rightarrow \gamma_m(\lambda) \rightarrow 0$
corresponding to $\alpha(\lambda) \rightarrow 3$

IR: Fixed point $\Rightarrow \gamma_m(\lambda) \rightarrow \gamma_*$
 γ_* scheme-independent,
expect $\gamma_* \lesssim 1$



Form of $\rho(\lambda)$ changes from $\rho(\lambda) \propto \lambda^3$ in the UV to $\rho(\lambda) \propto \lambda^{\alpha_*}$ in the IR

Expectations for eigenvalue mode number analysis

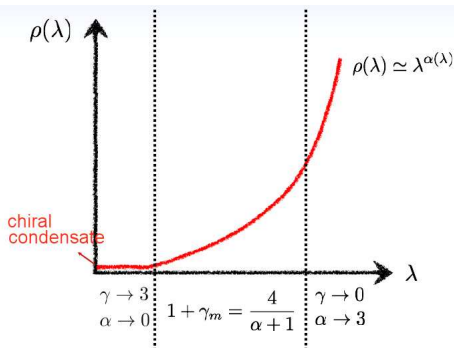
λ defines an energy scale;

fitting $\nu(\lambda) \propto \lambda^{1+\alpha(\lambda)}$ accesses $1 + \gamma_m(\lambda) = \frac{4}{1+\alpha(\lambda)}$ at that scale

For chirally broken systems:

UV: Asymp. freedom $\Rightarrow \gamma_m(\lambda) \rightarrow 0$
corresponding to $\alpha(\lambda) \rightarrow 3$

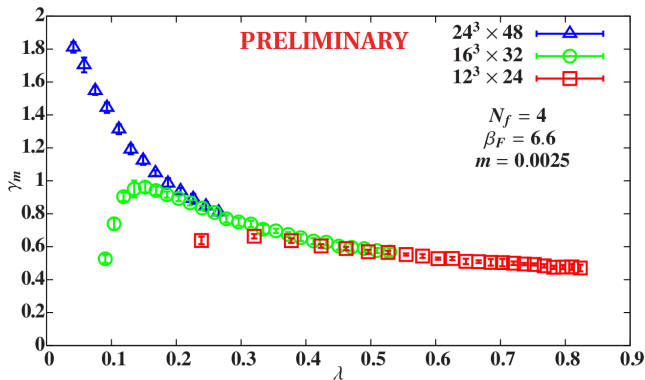
IR: $\langle \bar{\psi}\psi \rangle \propto \rho(0) > 0 \Rightarrow \alpha(\lambda) \rightarrow 0$
would produce “ $\gamma_m(\lambda) \rightarrow 3$ ”
but $\rho(\lambda)$ no longer $\sim \lambda^\alpha$



On the lattice we proceed by fitting $\nu(\lambda) \propto \lambda^{1+\alpha}$ in a limited range of λ

$N_F = 4$ runs rapidly : combine several volumes

Fit $\nu(\lambda) \propto \lambda^{1+\alpha}$ in a limited range of λ to find $1 + \gamma_m(\lambda) = \frac{4}{1 + \alpha(\lambda)}$



1000 eigenvalues
on each volume

Nearby points
use overlapping
fit ranges

$24^3 \times 48$ system
is chirally broken

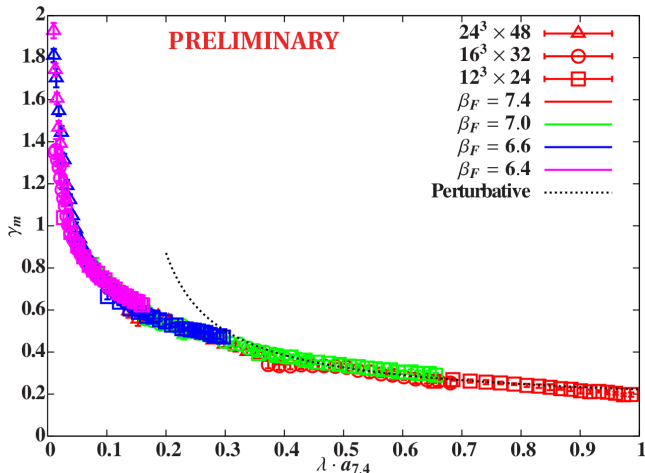
Focus on overlapping regions where different volumes agree
instead of studying finite-volume scaling behavior

Combine multiple couplings and volumes for $N_F = 4$

Rescale $\lambda \rightarrow \left(\frac{a_{7.4}}{a}\right)^{1+\gamma_m} \lambda$ to plot

in terms of single (smallest) lattice spacing

Match to one-loop perturbation theory at large $\lambda \cdot a_{7.4}$



Relative
lattice spacings
estimated from
Wilson flow &
MCRG matching:

$$a_{6.6} \approx 2a_{7.4}$$

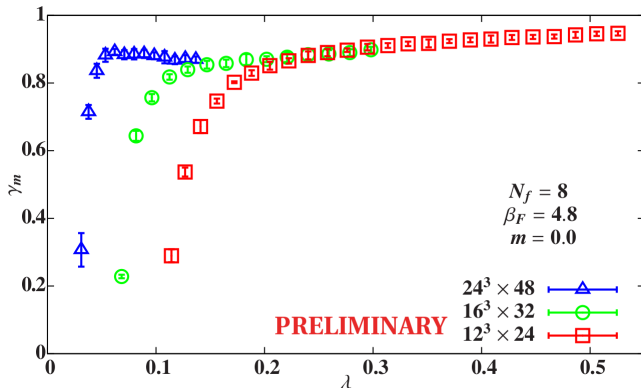
$$a_{6.4} \approx 2a_{7.0}$$

$$a_{6.4} \approx 1.3a_{6.6}$$

Finite-volume
“tails” omitted

$N_F = 8$ behaves very differently than $N_F = 4$

Fit $\nu(\lambda) \propto \lambda^{1+\alpha}$ in a limited range of λ to find $1 + \gamma_m(\lambda) = \frac{4}{1 + \alpha(\lambda)}$

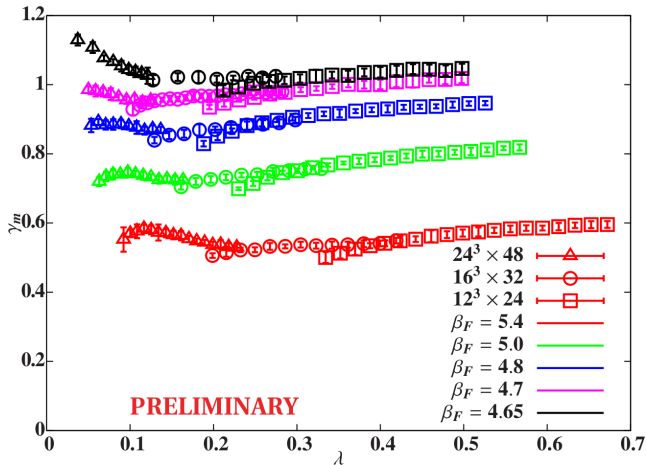


1000 eigenvalues
on each volume

$m = 0$,
all have $\rho(0) = 0$

In overlapping regions γ_m roughly independent of λ at fixed coupling β_F
No sign of asymptotic freedom – may be slight increase for larger λ

Cannot combine multiple couplings for $N_F = 8$



Relative
lattice spacings
not yet estimated
(not rescaling λ)

Finite-volume
“tails” omitted

γ_m remains roughly independent of λ in overlapping regions,
increases for stronger couplings until S^4 phase at $\beta \lesssim 4.65$
Clear contrast with $N_F = 4$

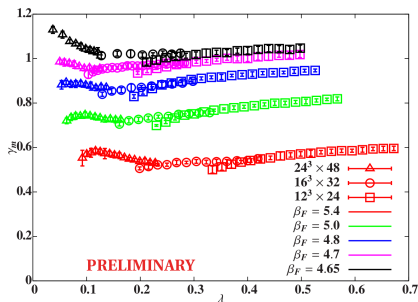
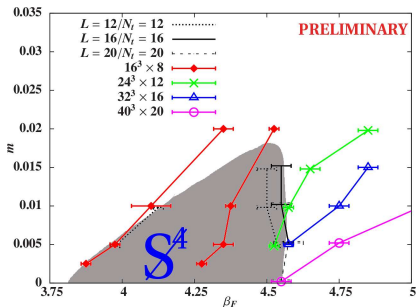
Recapitulation

$N_F = 8$ seems interesting, eigenvalues seem promising

Finite-temperature transitions on the $m = 0$ critical surface
show no scaling up to $N_T = 20$

$\gamma_m(\mu)$ accessible from eigenvalue mode number

Clear contrast between $N_F = 4$ and 8, latter more sensitive to β than λ



Thank you!

Thank you!

Collaborators

Anqi Cheng, Anna Hasenfratz, Gregory Petropoulos

Funding and computing resources

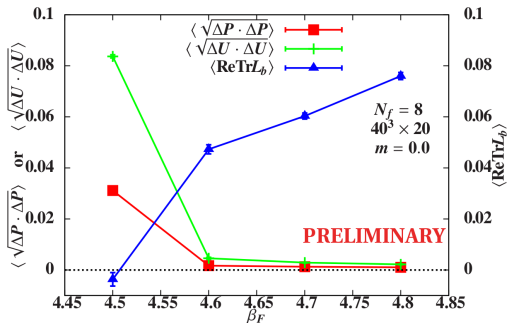
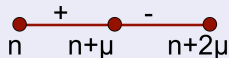


Backup: S^4 order parameters

Differences of plaquettes \square or links $\bar{\chi}U\chi$

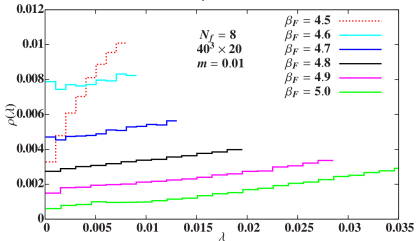
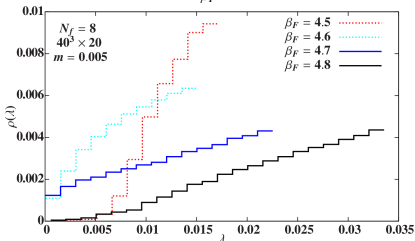
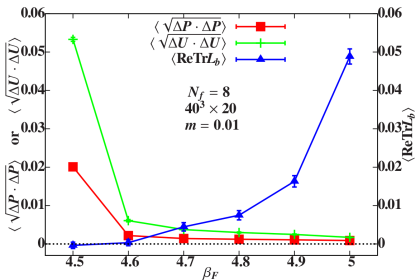
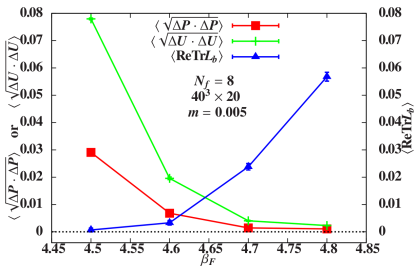
$$\Delta P_\mu = \langle \text{ReTr } \square_{n,\mu} - \text{ReTr } \square_{n+\mu,\mu} \rangle_{n_\mu \text{ even}}$$

$$\Delta U_\mu = \langle \alpha_{\mu,n} \bar{\chi}_n U_{\mu,n} \chi_{n+\mu} - \alpha_{\mu,n+\mu} \bar{\chi}_{n+\mu} U_{\mu,n+\mu} \chi_{n+2\mu} \rangle_{n_\mu \text{ even}}$$



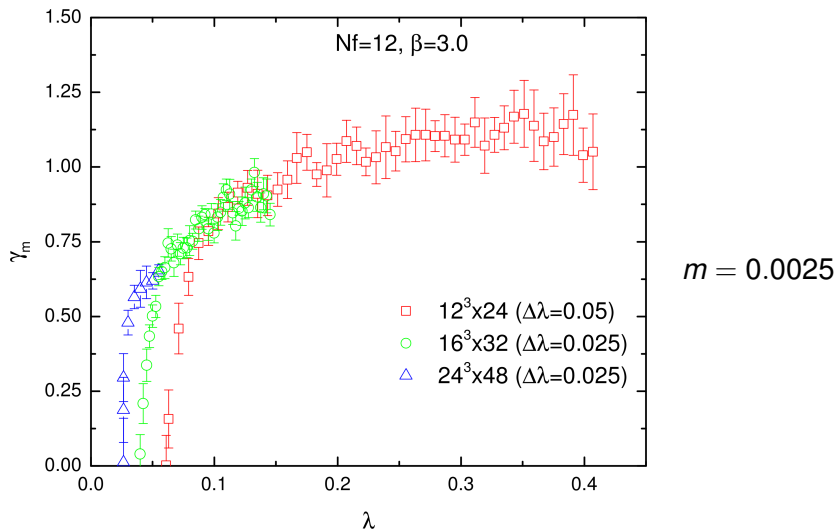
$N_F = 8$ results with $m = 0$ confirm signals in eigenvalue densities

Backup: $N_F = 8$ transitions for $40^3 \times 20$ with $m > 0$



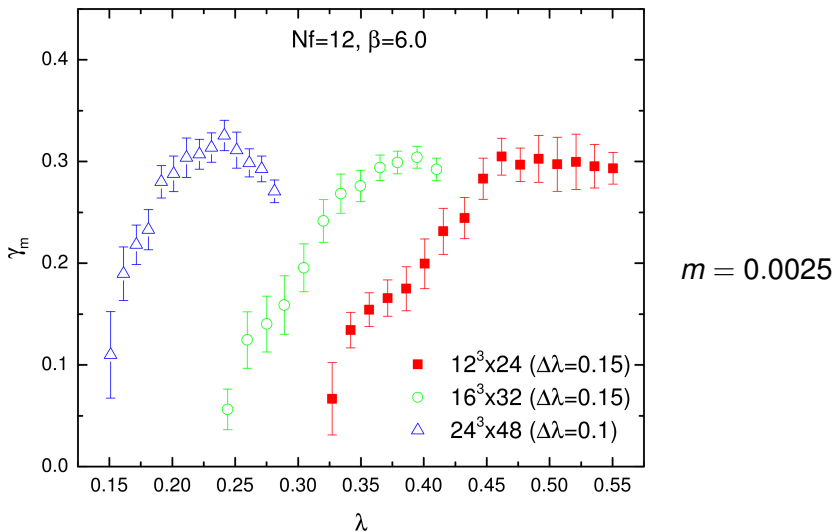
Between transitions, chirally broken systems with $\rho(0) > 0$

Backup: γ_m from eigenvalues for $N_F = 12$



At strong coupling (near S^4 phase) γ_m clearly increases with λ

Backup: γ_m from eigenvalues for $N_F = 12$



At weaker coupling, $\gamma_m \approx 0.3$ – better overlap needed