two-doublet lattice Higgs model

motivation
lattice action & symmetries
phases
symmetry breaking

(This is a preliminary report of work in progress with Richard Woloshyn.)

Recall that the standard SU(2)-Higgs model, with one scalar doublet, has two regions of parameter space *that are analytically connected*. (No gauge-invariant vev, no broken symmetry.)



adapted from Straumann, astro-ph/0409042 as adapted from textbook by Montvay & Münster Recall that the standard SU(2)-Higgs model, with one scalar doublet, has two regions of parameter space *that are analytically connected*. (No gauge-invariant vev, no broken symmetry.)

This phase diagram is easily mapped out by monitoring a quantity such as the gauge-invariant link $\equiv \left\langle \Phi^{\dagger}(x)U_{\mu}(x)\Phi(x+\mu) + h.c. \right\rangle$. analytic connection 0 0.25 ∞ $\kappa = \infty$ 0.8 Higgs region $\lambda = \infty$ 0.6 Σ ¥ $\partial \Sigma_1$ 0.4 $\lambda = 0$ K confinement region 0.2 $\lambda = \infty$ 0.2 0.4 0.6 0.8 0 $tanh(\beta/4)$ $\beta = \infty$ ß

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Wurtz, Lewis, Steele, Phys. Rev. D79, 074501 (2009).



Notice the effect of additional scalar doublets.

(These are identical doublets interacting only through gauge couplings.)



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Two-doublet models have been studied extensively by many authors using perturbative methods (continuum).

We want to understand the symmetries and symmetry breakings from a nonperturbative method (lattice). We consider an action that maintains a discrete symmetry for each scalar doublet:

$$S = \frac{\beta}{2} \sum_{x} \sum_{\mu=1}^{4} \sum_{\nu=1}^{4} \left[1 - \frac{1}{2} \operatorname{ReTr} U_{\mu}(x) U_{\nu}(x+\mu) U_{\mu}^{\dagger}(x+\nu) U_{\nu}^{\dagger}(x) \right]$$

$$- \sum_{x} \sum_{n=1}^{2} \sum_{\mu=1}^{4} \kappa_{n} \left[\Phi_{n}^{\dagger}(x) U_{\mu}(x) \Phi_{n}(x+\mu) + \operatorname{h.c.} \right]$$

$$+ \sum_{x} \sum_{n=1}^{2} \left[\Phi_{n}^{\dagger}(x) \Phi_{n}(x) + \lambda_{n} \left(\Phi_{n}^{\dagger}(x) \Phi_{n}(x) - 1 \right)^{2} \right]$$

$$+ \sum_{x} \left[\xi \Phi_{1}^{\dagger}(x) \Phi_{1}(x) \Phi_{2}^{\dagger}(x) \Phi_{2}(x) + \zeta \left| \Phi_{1}^{\dagger}(x) \Phi_{2}(x) \right|^{2} \right]$$

It is sometimes useful to change from doublet notation, $\Phi_n(x)$, to matrix notation,

$$\varphi_n(x) = \left(i\tau_2\Phi_n^*(x) \ \Phi_n(x)\right)$$

Heatbath algorithms, with additional accept-reject steps where necessary, are used for gauge updates and for scalar updates.

SU(2) gauge invariance:

 $\Phi_1(x) \to R(x)\Phi_1(x), \quad \Phi_2(x) \to R(x)\Phi_2(x), \quad \text{and} \quad U_\mu(x) \to R(x)U_\mu(x)R^{\dagger}(x+\mu)$

SU(2) global symmetry for scalar doublet #1:

$$\varphi_1(x) \to \varphi_1(x) R_{(1)}$$

SU(2) global symmetry for scalar doublet #2:

$$\varphi_2(x) \to \varphi_2(x) R_{(2)}$$

In the special case where the action and all constraints are functions of $\Phi_1^{\dagger}\Phi_1 + \Phi_2^{\dagger}\Phi_2$ rather than being functions of $\Phi_1^{\dagger}\Phi_1$ and $\Phi_2^{\dagger}\Phi_2$ separately, we also have a global U(2) symmetry,

$$\begin{pmatrix} \Phi_1(x) \\ \Phi_2(x) \end{pmatrix} \to R_{\text{flavor}} \begin{pmatrix} \Phi_1(x) \\ \Phi_2(x) \end{pmatrix}$$

but this is partially redundant. Only a $U(1) \times U(1)$ global symmetry is new.





0.8

0.9

1

0

0

0.2

0.1

0.3

0.4

0.5

κ₁

0.6

0.7





1. motivation 2. lattice action & symmetries 3. phases 4. symmetry breaking

To test for spontaneous symmetry breaking of SU(2)×SU(2), add a small explicit breaking term: $\delta S = \frac{\eta}{2} \text{Tr} \left[\varphi_1^{\dagger}(x) \varphi_2(x) \right]$



$$12^4$$
, $\beta=8$, $\lambda_1=\lambda_2=1$, $\xi=0$, $\zeta=0$



An order parameter after extrapolation to $\eta = 0$.



summary

We have mapped out the phase diagram of a two-doublet Higgs model through lattice simulations.

The simulations provide evidence of spontaneously broken global symmetries (in contrast to the one-doublet Higgs model).

The lattice formulation is gauge invariant, and the gauge symmetry is not broken in this formulation (which is also true in the one-doublet lattice Higgs model).