

THE IMPORTANCE OF WALKING

**Lattice Gauge Theory for the LHC
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References:

1. TC: S. Weinberg, Phys. Rev. **D19**, 1277 (1979);
L. Susskind, Phys. Rev. **D20**, 2619 (1979).
2. ETC and FCNC: E. Eichten and K. Lane, Phys. Lett. **B90**, 125 (1980).
3. WTC: B. Holdom, Phys. Rev. **D24**, 1441 (1981);
Phys. Lett. **B150**, 301 (1985);
especially: T. Appelquist, D. Karabali, and L. C. R. Wijewardhana, Phys. Rev. Lett. **57**, 957 (1986);
T. Appelquist and L. C. R. Wijewardhana,
Phys. Rev. **D36**, 568 (1987);
K. Yamawaki, M. Bando, and K. Matumoto,
Phys. Rev. Lett. **56**, 1335 (1986);
T. Akiba and T. Yanagida, Phys. Lett. **B169**, 432 (1986).

4. S,T,U: T. Appelquist and C. Bernard,
Phys. Rev. **D22**, 200 (1980);
A. Longhitano, Phys. Rev. **D22**, 1166 (1980);
Nucl. Phys. **B188**, 118 (1981);
B. W. Lynn, M. E. Peskin, and R. G. Stuart, in
Trieste Electroweak 1985, 213 (1985);
D. Kennedy and B. Lynn, Nucl. Phys. **B322**, 1 (1989);
M. E. Peskin and T. Takeuchi, Phys. Rev. Lett. **65**, 964
(1990); Phys. Rev. **D46**, 381 (1992);
M. Golden and L. Randall, Nucl. Phys. **B361**, 3 (1990);
B. Holdom and J. Terning, Phys. Lett. **B247**, 88 (1990);
G. Altarelli and R. Barbieri, Phys. Lett. **B253**, 161
(1990).
5. WTC for FCNC and S: K. Lane, *Two Lectures on Technicolor*, GIF-2001, Annecy-le-Vieux, hep-ph/0202255.

I. TECHNICOLOR and EXTENDED TECHNICOLOR

- ELECTROWEAK SYMM. BREAKING:

What **dynamics** are responsible for

$$SU(2)_{EW} \otimes U(1)_{EW} \longrightarrow U(1)_{EM} \quad ??$$

→ Technicolor (modeled on QCD and its chiral symmetry breaking)

- FLAVOR and FLAVOR SYMM. BRKNG:

The **origin** of quark and lepton flavors?

Why three identical generations?

What is the **dynamical origin** of nontrivial quark and lepton masses and mixings?

→ Extended Technicolor (modeled on the the Higgs VEV as a T-fermion condensate)

TECHNICOLOR — NATURAL EWSB

TC = asymptotically free gauge theory
of MASSLESS techni-fermions T
with $\Lambda_{TC} \simeq 0.1 - 1.0 \text{ TeV}$.

Assume (e.g.) $T_{L,Ri} = (U_i, D_i)_{L,R}$ form
 N_D LH $SU(2)_{EW}$ -doublets and $2N_D$ RH-singlets.

$\Rightarrow$ chiral symmetry **and** χ SB:

$$(SU(2) \otimes U(1))_{EW} \subset SU(2N_D)_L \otimes SU(2N_D)_R \\ \longrightarrow SU(2N_D)_V$$

$$\langle \bar{U}_{Li} U_{Rj} \rangle = \langle \bar{D}_{Li} D_{Rj} \rangle = -\frac{1}{2} \delta_{ij} \Delta_T$$

$$\left(\Delta_T \simeq 4\pi F_T^3 \text{ where } F_T = 246 \text{ GeV} / \sqrt{N_D} \right)$$

$$\Rightarrow (SU(2) \otimes U(1))_{EW} \longrightarrow U(1)_{EM}$$

$$\text{with } M_W^2 = M_Z^2 \cos^2 \theta_W = \frac{1}{2} g^2 N_D F_T^2$$

NATURALNESS: *TC's GUIDING PRINCIPLE!*

$$\Lambda_{EW} \equiv \Lambda_{\chi SB} \sim F_T \sim \Lambda_{TC} \sim M_{\text{technihadron}} \\ \text{at } \alpha_{TC}(\Lambda_{TC}) = \mathcal{O}(1)$$

EXTENDED TECHNICOLOR (ETC) SCENARIO FOR FLAVOR PHYSICS

- *WHY EXTENDED TECHNICOLOR?*

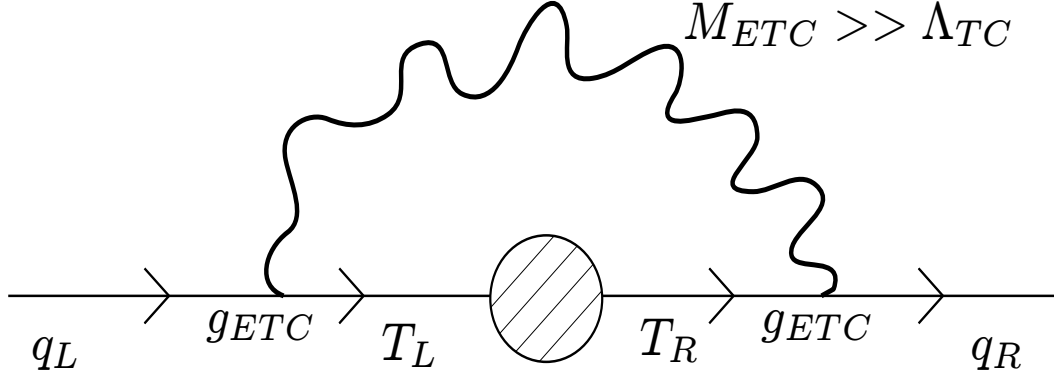
- Technicolor (and color) by itself leaves too much chiral symmetry unbroken.
- Explicit breaking of q, ℓ chiral symmetries required to give “hard” q, ℓ masses and avoid massless Goldstone bosons — $\pi, K, \eta, \dots$
- Explicit breaking of technifermion chiral symmetries required to give hard masses to T_i and avoid massless technipions π_T — “AXIONS”!

- *THE STRUCTURE OF ETC:*

ETC = gauge interaction of massless
 T_i, q_a, ℓ_a — all together in a few
ETC representations:

$$\underline{G_{ETC} \supset SU(N_{TC}) \otimes SU(3)_C \otimes \text{Flavor}}$$

$G_{ETC} \longrightarrow SU(N_{TC}) \otimes SU(3)_C$ at
 $M_{ETC} \gg \Lambda_{TC} \lesssim 1 \text{ TeV},$
*explicitly breaking **ALL** global flavor*
symmetries.



$$m_{q,\ell}(M_{ETC}) \simeq \frac{g_{ETC}^2}{M_{ETC}^2} \langle \bar{T} T \rangle_{ETC}$$

$$F_T^2 M_{\pi_T}^2 \simeq \frac{g_{ETC}^2}{2M_{ETC}^2} \langle \bar{T} T \bar{T} T \rangle_{ETC}$$

$$\langle \bar{T} T \rangle_{ETC} = \langle \bar{T} T \rangle_{TC} \exp \left(\int_{\Lambda_{TC}}^{M_{ETC}} \frac{d\mu}{\mu} \gamma_m(\mu) \right)$$

$$\begin{array}{c} \uparrow \\ \sim \underline{\underline{4\pi F_T^3}} \end{array}$$

$$\gamma_m(\mu) = \frac{3C_2(R) \alpha_{TC}(\mu)}{2\pi} + \dots$$

II. TC and ETC SHOW-STOPPERS

FLAVOR-CHANGING NEUTRAL CURRENTS

Generic ETC models with *realistic* m_q -matrices have FCNC involving *mass eigenstate* q, ℓ . e.g.:

$$\begin{aligned} [q^\dagger T, T^\dagger q'] &= q^\dagger q' \\ \Rightarrow \mathcal{H}'_{\Delta|S|=2} &= \frac{g_{ETC}^2 V_{ds}^2}{M_{ETC}^2} \bar{d} \Gamma^\mu_s \bar{d}'_{\mu} s + \text{h.c.} \end{aligned}$$

$$|\Delta S| = 2 \text{ FCNC} \Rightarrow \frac{M_{ETC}}{g_{ETC} \sqrt{\text{Re}(V_{ds}^2)}} \gtrsim 1300 \text{ TeV}$$

Scaling $\langle \bar{T} T \rangle_{ETC}$ from QCD (for $\gamma_m \ll 1$) $\Rightarrow$

$$m_{q,\ell}(M_{ETC}) \simeq \frac{g_{ETC}^2}{M_{ETC}^2} \langle \bar{T} T \rangle_{ETC} \lesssim \frac{0.1 \text{ MeV}}{|V_{ds}|^2 N_D^{3/2}}$$

$$M_{\pi_T} \simeq \frac{g_{ETC}^2}{\sqrt{2} M_{ETC}^2} \frac{\langle \bar{T} T \rangle_{ETC}}{F_T} \lesssim \frac{0.4 \text{ GeV}}{|V_{ds}| N_D}$$

PRECISION ELECTROWEAK TESTS

ASSUMING all new physics scales such as Λ_{TC} are $\gg M_{W,Z}$, “Oblique” correction factor S is:

$$\begin{aligned} S &= 16\pi \frac{d}{dq^2} \left[\Pi_{33}(q^2) - \Pi_{3Q}(q^2) \right]_{q^2=0} \\ &= 4\pi \int \frac{dm^2}{m^4} \left[\rho_V^3(m^2) - \rho_A^3(m^2) \right] \\ &= -0.04 \pm 0.09 \text{ (PDG)} \end{aligned}$$

ESTIMATES OF TC CONTRIBUTION TO S

ASSUME THAT TC is QCD-like:

e.g., Peskin & Takeuchi: QCD as analog computer — spectral-function sum rules; VMD; scale from ρ , a_1 masses $\Rightarrow$

$$S_{TC} = 4\pi \frac{N_D F_T^2}{M_{\rho_T}^2} \left(1 + \frac{M_{\rho_T}^2}{M_{a_T}^2} \right) \simeq 0.25 N_D \frac{N_{TC}}{3}$$

III. SAVING THE SHOW!

The KEY Equations of ETC — Again:

$$m_{q,\ell}(M_{ETC}) \simeq \frac{g_{ETC}^2}{M_{ETC}^2} \langle \bar{T}T \rangle_{ETC}$$

$$F_T^2 M_{\pi_T}^2 \simeq \frac{g_{ETC}^2}{2M_{ETC}^2} \langle \bar{T}T\bar{T}T \rangle_{ETC}$$

$$\langle \bar{T}T \rangle_{ETC} = \langle \bar{T}T \rangle_{TC} \exp \left(\int_{\Lambda_{TC}}^{M_{ETC}} \frac{d\mu}{\mu} \gamma_m(\mu) \right)$$

CAN γ_m BE LARGE for a LARGE μ -range —
ENHANCING $\langle \bar{T}T \rangle_{ETC} / \langle \bar{T}T \rangle_{TC}$?

YES!

(Well...**YOU** need to demonstrate this!!)

HOW to keep γ_m LARGE?

Suppose $\alpha_{TC}(\mu) \simeq \text{constant}$, i.e.,

$$\beta(\alpha_{TC}) = \frac{d\alpha_{TC}(\mu)}{d\log(\mu)} \simeq 0$$

Solution to the (approximate) equation for the anomalous dimension γ_m is

$$\implies \gamma_m(\mu) = 1 - \sqrt{1 - \alpha_{TC}(\mu)/\alpha_{\chi SB}}$$

where $\alpha_{\chi SB} = \pi/3C_2(R)$ (in ladder approx'n).

This indicates that χSB occurs *iff*

$$\alpha_{TC} \geq \alpha_{\chi SB}.$$

And, at the χSB scale Λ_{TC} , $\gamma_m(\Lambda_{TC}) = 1$!

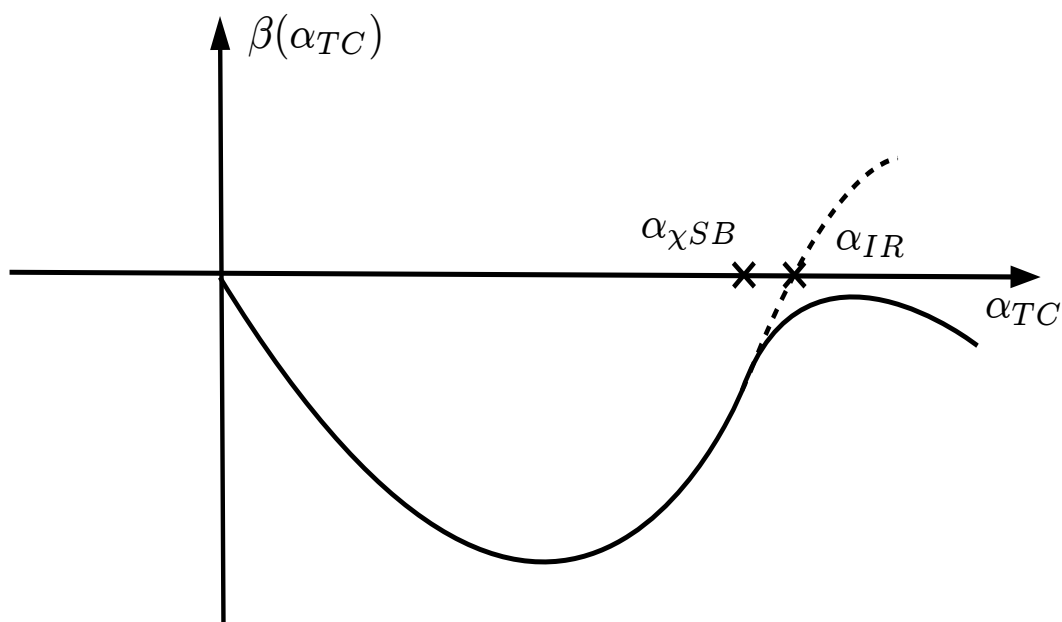
i.e., $\gamma_m = 1$ is the signal for χSB !!

Again: to keep γ_m **large** for a **large** range of μ

$$\beta(\alpha_{TC}(\mu)) \simeq 0 \text{ for } \mu \simeq \Lambda_{TC} \longrightarrow \sim M_{ETC}$$

$$\implies \alpha_{TC}(\mu) \simeq \alpha_{\chi SB} \implies \gamma_m(\mu) \simeq 1$$

HOW? $\alpha_{\chi SB}$ lies just below an IR-fixed point:



i.e., TC must be a WALKING gauge theory!

CURING FCNC

Must have $M_{ETC}/(g_{ETC}|V_{ds}|) \sim 10^3\text{--}10^4 \text{ TeV}$
for the 1st two generations! The WTC
enhancement of $\langle \bar{T}T \rangle_{ETC}$ compensates $1/M_{ETC}^2$.

In the *extreme* walking limit, e.g.:

$$\langle \bar{T}T \rangle_{ETC} \simeq \frac{M_{ETC}}{\Lambda_{TC}} \times \langle \bar{T}T \rangle_{TC}$$

$$m_{q,\ell}(M_{ETC}) \simeq \frac{4\pi\alpha_{ETC}\langle \bar{T}T \rangle_{TC}}{M_{ETC}\Lambda_{TC}} \simeq \frac{16\pi^2\alpha_{ETC}F_\pi^2F_T}{N_D M_{ETC}\Lambda_{TC}}$$

$$\Rightarrow M_{ETC} = (140 \rightarrow 700,000 \text{ TeV}) \left(\frac{10}{N_D} \frac{\alpha_{ETC}}{0.75} \frac{F_T}{\Lambda_{TC}} \right)$$

for $m_q = 5 \text{ GeV} \rightarrow 1 \text{ MeV}$

Furthermore...

$$M_{\pi_T} \simeq \underbrace{\sqrt{1.5\pi} 4\pi F_\pi}_{7(!!) \text{ TeV}} \left(\frac{F_T}{\Lambda_{TC}} \right) \sqrt{\frac{\alpha_{ETC}}{0.75 N_D}}$$

... π_T may NOT be pseudo-Goldstone bosons!

In any case, M_{π_T} are significantly enhanced!

This has ENORMOUS implications for WTC phenomenology.

See Tulika's talk.

S in walking technicolor

WTC $\implies$ spectral integrals

$$\int \frac{dm^2}{m^n} [\rho_V^3(m^2) - \rho_A^3(m^2)]$$

used in the S -calculation converge much slower than in QCD $\implies$ integrals **are not** dominated by just the lowest-lying resonances.

- Spectrum of ρ_T , a_T **very different** from QCD.
- Cannot justify scaling arguments from QCD.
- Large- N_{TC} arguments may fail for large N_D .

Expect a *tower* of ρ_T , a_T ??

$M_{a_T} \simeq M_{\rho_T}$ and $g_{a_T} \simeq g_{\rho_T}$ decrease $S \rightarrow 0$!

This talk needs a WOW FINISH!!

Either you'll provide it...or the LHC will.