

Supersymmetric Yang-Mills with domain wall fermions

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Lattice gauge theory for LHC physics
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Supersymmetry (SUSY) on the lattice

- Much effort devoted toward formulating supersymmetric lattice theories
 - fascinating theoretical challenge
 - potential role of SUSY in beyond the standard model physics
 - desire to better understand various nonperturbative aspects of SUSY theories via numerical simulation

SUSY on the lattice

- Naive lattice discretization explicitly breaks SUSY:

$$(\text{supercharges}) \quad \quad \quad (\text{generator of translations})$$

$\downarrow \quad \quad \quad \downarrow$

$$\{\bar{Q}, Q\} \sim P$$

- Undesirable SUSY violating operators may arise through radiative corrections
 - Fine-tuning of operators is required to reach SUSY limit

SUSY on the lattice

$N=1$ SYM is the only four-dimensional SUSY theory that can be studied using conventional lattice discretizations and without fine tuning

- Only relevant SUSY violating operator allowed by gauge and hypercubic symmetries is a gluino mass term
- Chiral symmetry realized on lattice implies SUSY restoration in continuum limit (e.g., overlap or domain wall fermions)

$N=1$ super Yang-Mills (SYM)

$$L = \frac{1}{g^2} \left[\bar{\lambda} \gamma_\mu D_\mu \lambda + \frac{1}{4} v_{\mu\nu} v_{\mu\nu} \right] , \quad \bar{\lambda} = \lambda^T C$$

- One of the simplest of SUSY gauge theories in terms of field content
 - 1 vector field and 1 adjoint Majorana fermion (gluino)
 - a single input parameter, the gauge coupling (g)
- Anomalous $U(1)_R$ (chiral) symmetry: $\lambda \rightarrow e^{-i\alpha\gamma_5} \lambda \Rightarrow \theta \rightarrow \theta - 2N_c\alpha$
- Z_{2N_c} subgroup of $U(1)_R$ survives at quantum level
 - partition function invariant for: $\alpha = \frac{\pi k}{N_c}$, $k = 0, \dots, 2N_c - 1$

$$L = \frac{1}{g^2} \left[\bar{\lambda} \gamma_\mu D_\mu \lambda + \frac{1}{4} v_{\mu\nu} v_{\mu\nu} \right] , \quad \bar{\lambda} = \lambda^T C$$

- Gluino condensation: $\langle \bar{\lambda} \lambda \rangle \neq 0$
 - discrete chiral symmetry breaking: $Z_{2N_c} \rightarrow Z_2$
- Confinement
 - colorless bound states (glue-glue, glue-gluino, gluino-gluino)
- No SUSY breaking (non-vanishing Witten index)

N=1 SYM

$$L = \frac{1}{g^2} \left[\bar{\lambda} \gamma_\mu D_\mu \lambda + \frac{1}{4} v_{\mu\nu} v_{\mu\nu} \right] , \quad \bar{\lambda} = \lambda^T C$$

- Why study N=1 SYM via numerical simulations?
 - natural starting point: simplest of SUSY theories, can simulate with well understood lattice actions
 - test exact theoretical predictions about SYM (e.g., discrete chiral symmetry breaking)
 - access to nonperturbative aspects of SYM, for which less is known (e.g., spectrum)

Numerical studies of N=1 SYM

- Wilson
 - DESY-Muenster-Roma Collaboration
(hep-lat/0112007 and references therein, arXiv:0811.1964)
- Domain wall fermions
 - G.T. Fleming, et al. (hep-lat/0008009)
 - M. G. Endres (arXiv:0810.0431, arXiv:0902.4267)
 - J. Giedt, et. al. (arXiv:0807.2032, arXiv:0810.5746)

Simulation details

- Wilson gauge action (fund. rep.) + DWF (adj. rep.)
- SU(2) gauge group
- Simulations performed:
 - using modified Columbia Physics System (CPS)
 - on QCDOC at Columbia University and New York Blue (BlueGene/L) at Brookhaven National Laboratory



Simulation parameters

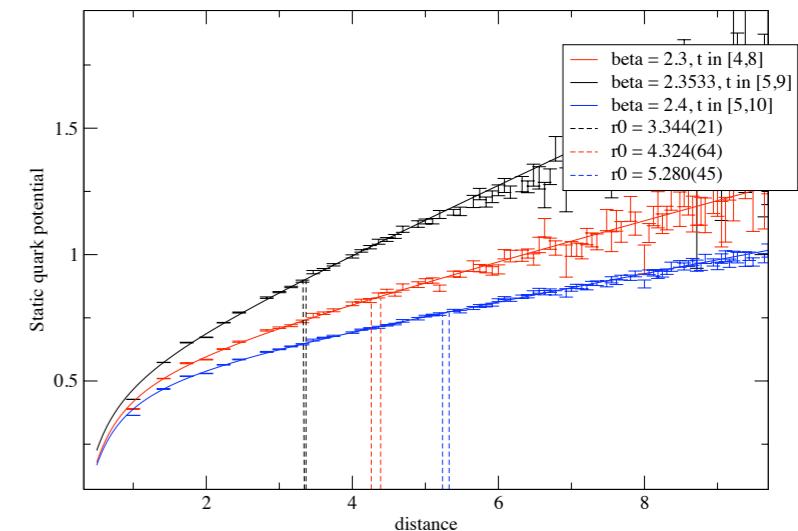
$L^3 \times T$	β	L_s	m_f
$16^3 \times 32$	2.3	16	0.01
			0.02
		20	0.04
			0.01
		24	0.02
			0.04
		28	0.01
			0.02
		32	0.04
			0.01
			0.02
		40	0.02
		48	0.02
2.353	28	0.02	
2.4	28	0.02	

$L^3 \times T$	β	L_s	m_f
$8^3 \times 8$	2.3	16	0.02
			20
		24	0.02

- Static potential
- Residual mass
- Gluino condensate
- Dirac spectrum
- Topological charge
- Spectrum

Summary of basic measurements

- Static potential:
 - $r_0/L \approx 0.18$ ($\beta=2.3$) to 0.30 ($\beta=2.4$)
- Residual mass:
 - $(m_f + m_{\text{res}}) r_0 \sim 0.35$ ($\beta=2.4$) - 0.5 ($\beta=2.3$) for $L_s=28$
 - m_{res} is 5-10 times larger than m_f
- Gluino condensate in the chiral limit:
 - 20-25% systematic errors due to large residual chiral symmetry breaking



Dirac spectrum

- Check if DWFs are working correctly for chosen parameters
 - are low modes bound to fifth dimension boundaries?
 - does Dirac spectrum reflect expected continuum properties?
 - do the matrix elements of the chirality operator make sense?
- Alternative approach for estimating residual mass
- Determine gluino condensate via the Banks-Casher relation

Topological charge

- Check topological ergodicity
- Check consistency of topological charge and near zero-modes of Dirac operator

Continuum (Euclidean) Dirac operator

Operator:	$\gamma_\mu D_\mu$	$D = \gamma_\mu D_\mu + m_g$	$D_H = \gamma_5 D$
Eigenvalue:	$\pm i\lambda$	$\pm i\lambda + m_g$	$\pm \lambda_H = \sqrt{\lambda^2 + m_g^2}$
Eigenstate:	$ \pm \lambda\rangle$	$ \pm \lambda\rangle$	$ \pm \lambda_H\rangle \propto \sqrt{-i\lambda + m_g} \lambda\rangle \pm \sqrt{i\lambda + m_g} -\lambda\rangle$

- Zero-modes unpaired and eigenstates of chirality operator
- Eigenvalues come in positive/negative pairs

$$\gamma_\mu^\dagger = \gamma_\mu \quad \{ \gamma_5, \gamma_\mu D_\mu \} = 0$$

- Eigenvalues are two-fold degenerate

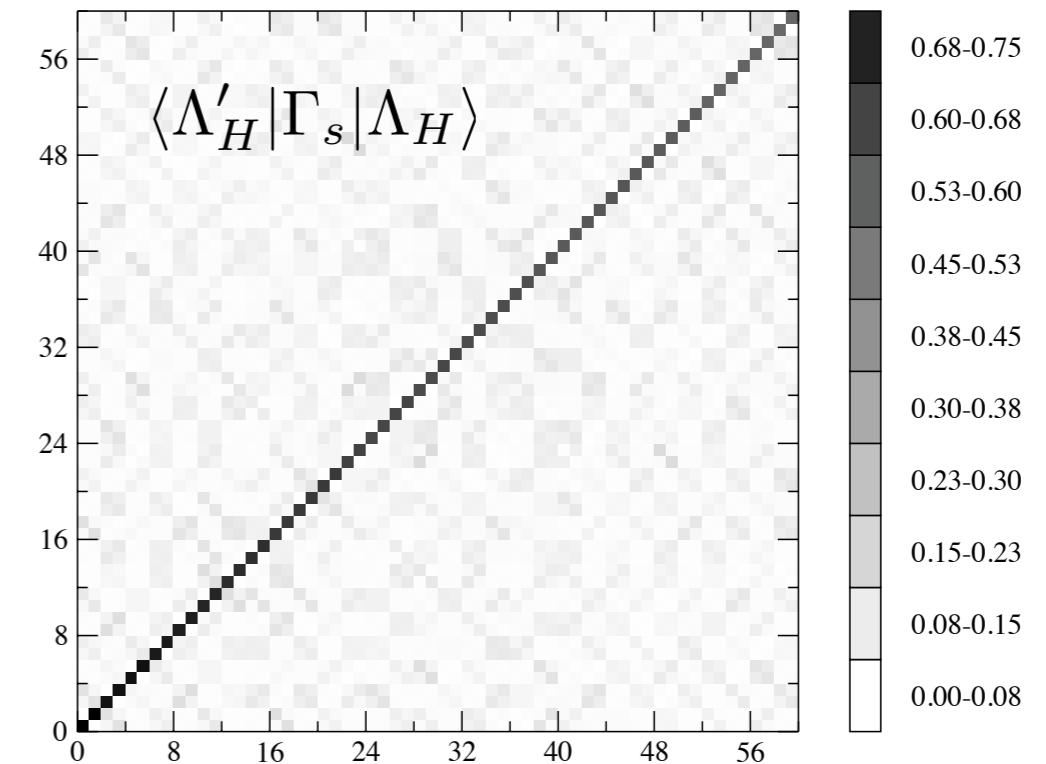
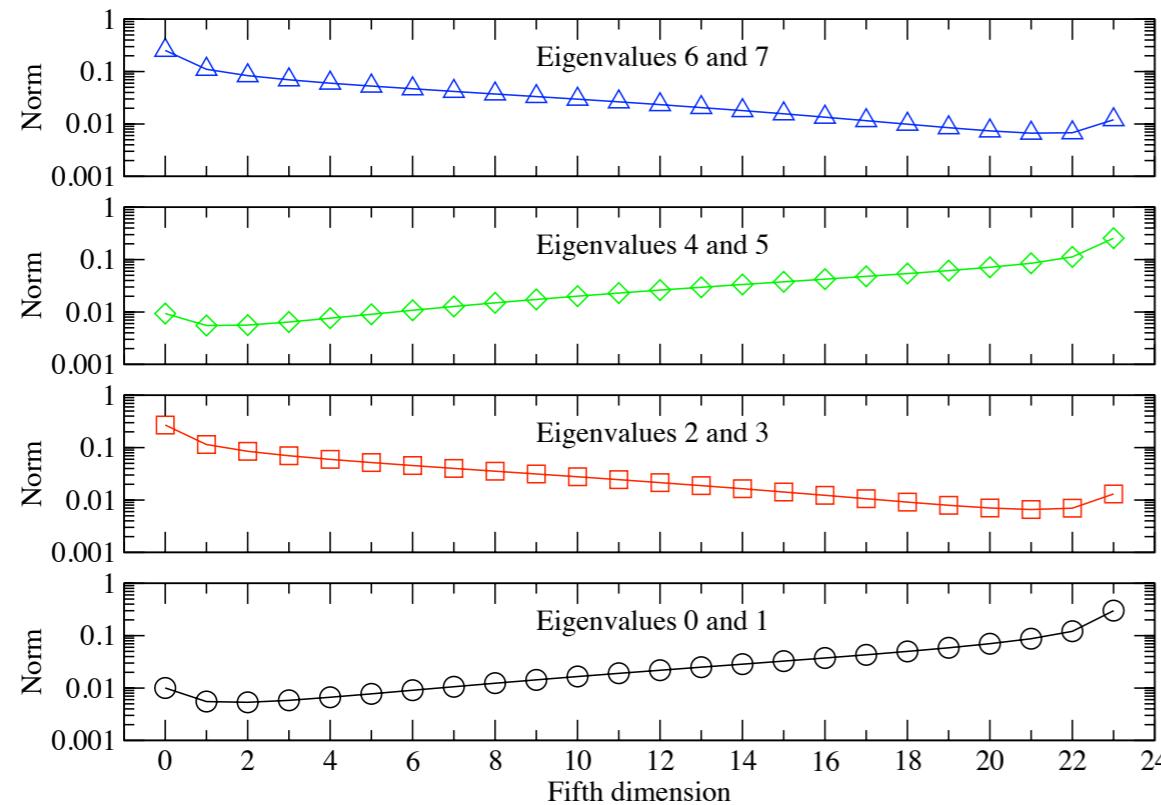
$$C^\dagger DC = D^*$$

Eigenvalue measurement details

- Performed eigenvalue studies of the Hermitian DWF Dirac operator (D_H)
 - eigenvalues are two-fold degenerate
 - \pm pairing of eigenvalues absent at finite lattice spacing
- Measurements performed on 8^4 lattices
- Obtained lowest 64 eigenvalues (Λ_H) on each of 150 gauge field configurations

Low modes of D_H

$\beta=2.3$, $m_f=0.2$ and $L_s=24$



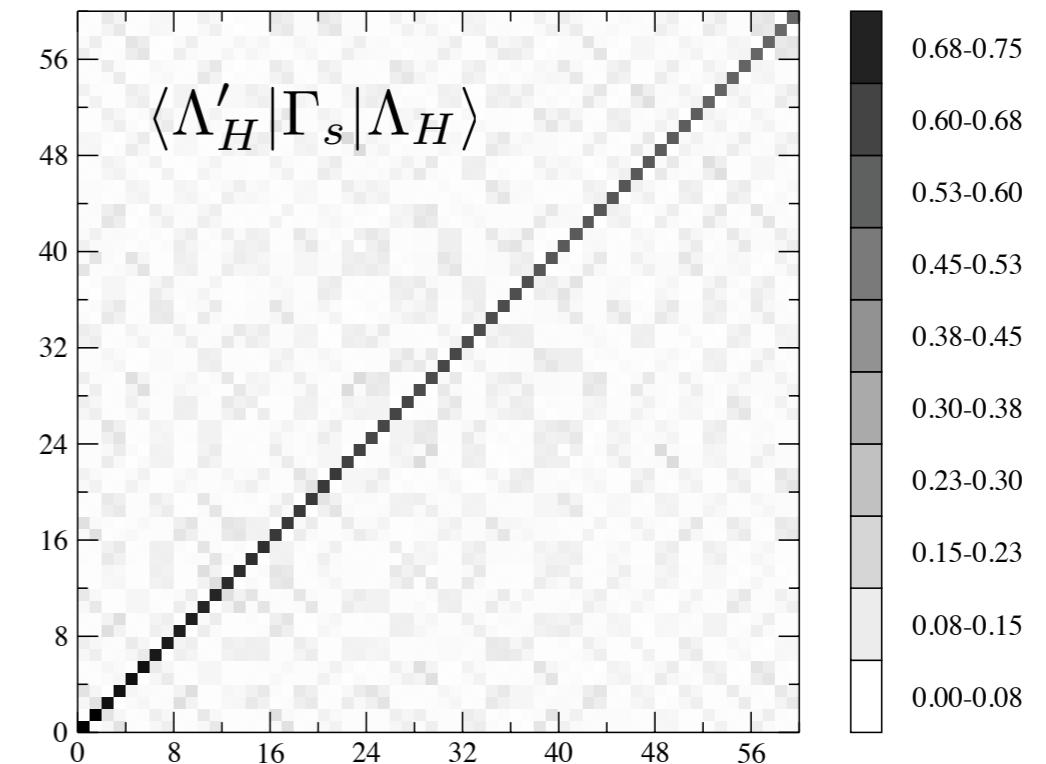
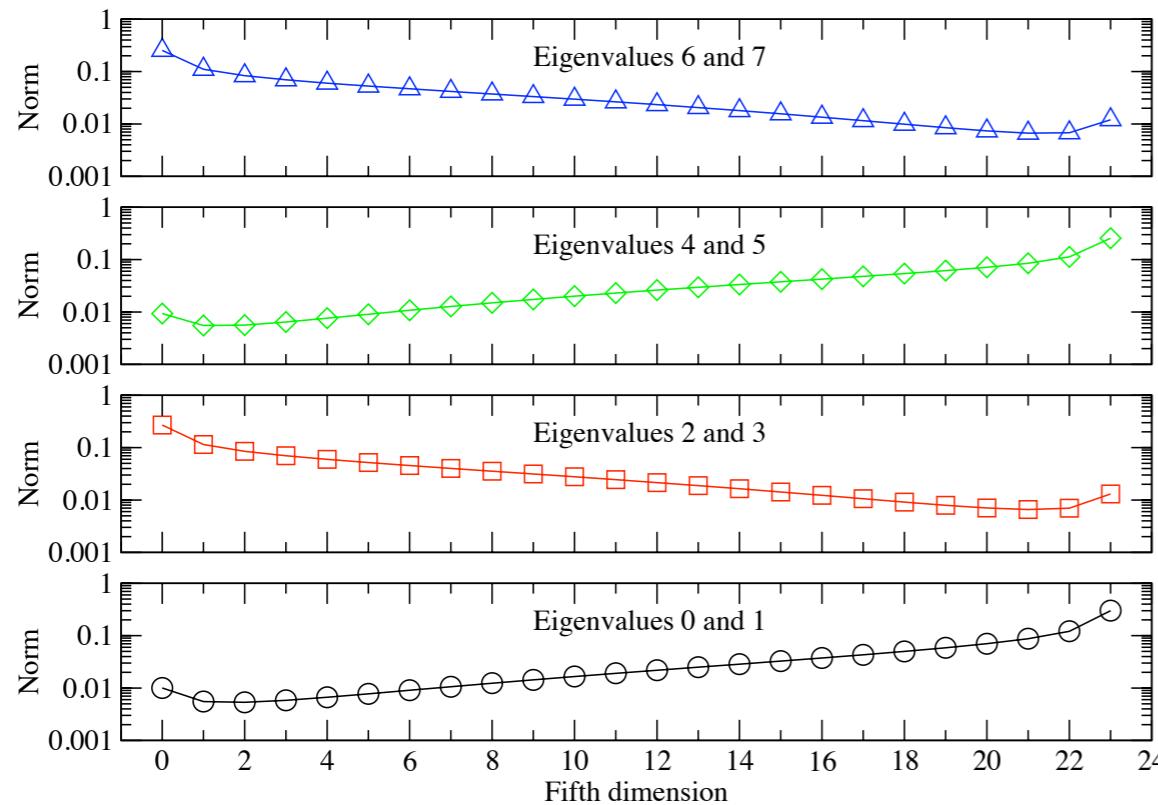
$$\mathcal{N}_{\Lambda_H}(s) = \sum_x \Psi_{\Lambda_H}^\dagger(x, s) \Psi_{\Lambda_H}(x, s)$$

$$D_H \Psi_{\Lambda_H} = \Lambda_H \Psi_{\Lambda_H}$$

$$\langle \Lambda'_H | \Gamma_s | \Lambda_H \rangle = \sum_{x,s} sgn \left(\frac{L_s - 1}{2} - s \right) \Psi_{\Lambda'_H}^\dagger(x, s) \Psi_{\Lambda_H}(x, s)$$

Low modes of D_H

$\beta=2.3, m_f=0.2$ and $L_s=24$



- Modes exponentially localized on left/right walls
- Modes have definite chirality, consistent with: $m_g \gg \lambda$
- in continuum theory: $\langle \lambda'_H | \gamma_5 | \lambda_H \rangle = \frac{1}{\lambda_H} [m_g \delta_{\lambda'_H, \lambda_H} + i|\lambda| \delta_{-\lambda'_H, \lambda_H}]$

Extracting information from Λ_H

- Fit valence mass dependence of eigenvalues using reparameterized Taylor expansion:

$$\Lambda_H^2(m_v) = n_5^2 [\lambda^2 + (m_v + \delta m)^2] + \mathcal{O}(m_v^3)$$

T. Blum, et al., hep-lat/0007038

- Can show:

$$\langle \bar{\psi} \psi \rangle = \frac{1}{12V} \sum_i \left\langle \frac{m_f + \delta m_i}{\lambda_i^2 + (m_f + \delta m_i)^2} \right\rangle + \mathcal{O}(m_v^3)$$

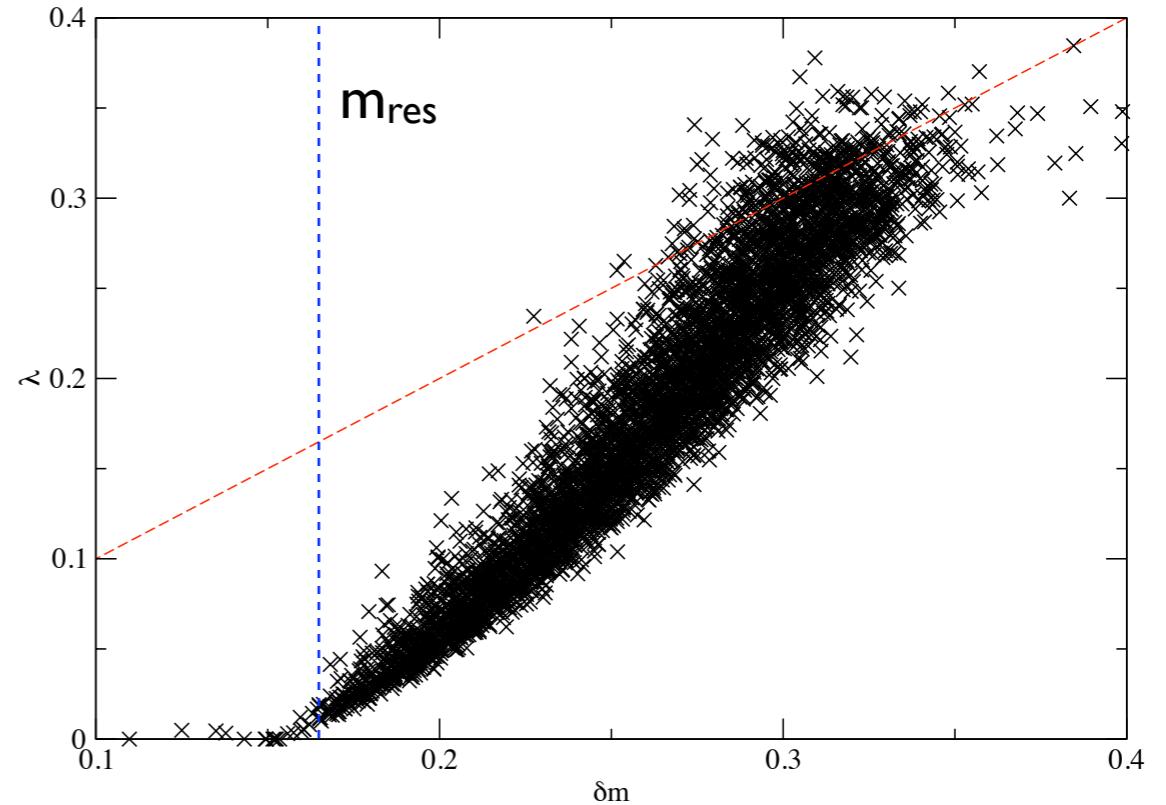
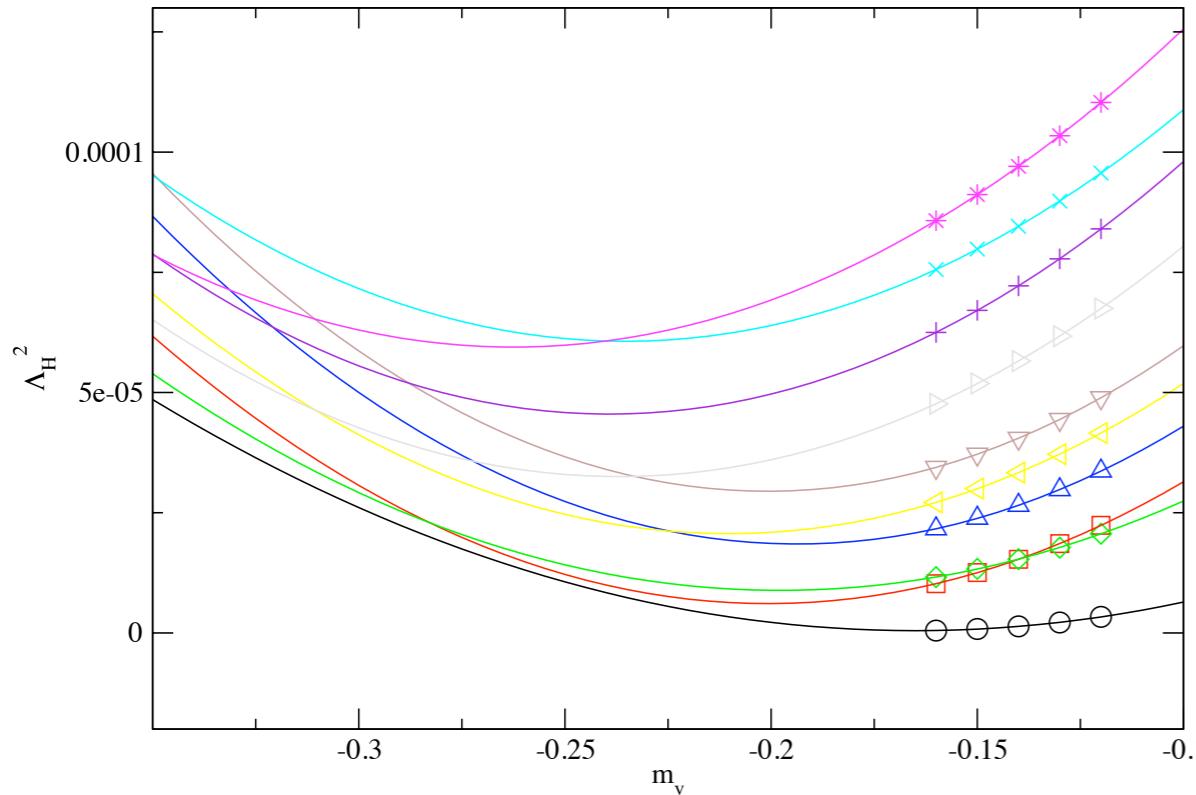
“effective 4D eigenvalue of D-slash”

reflects residual chiral symmetry
breaking due to finite L_s

- Use λ and δm distributions to extract condensate and m_{res}

Extracting information from Λ_H

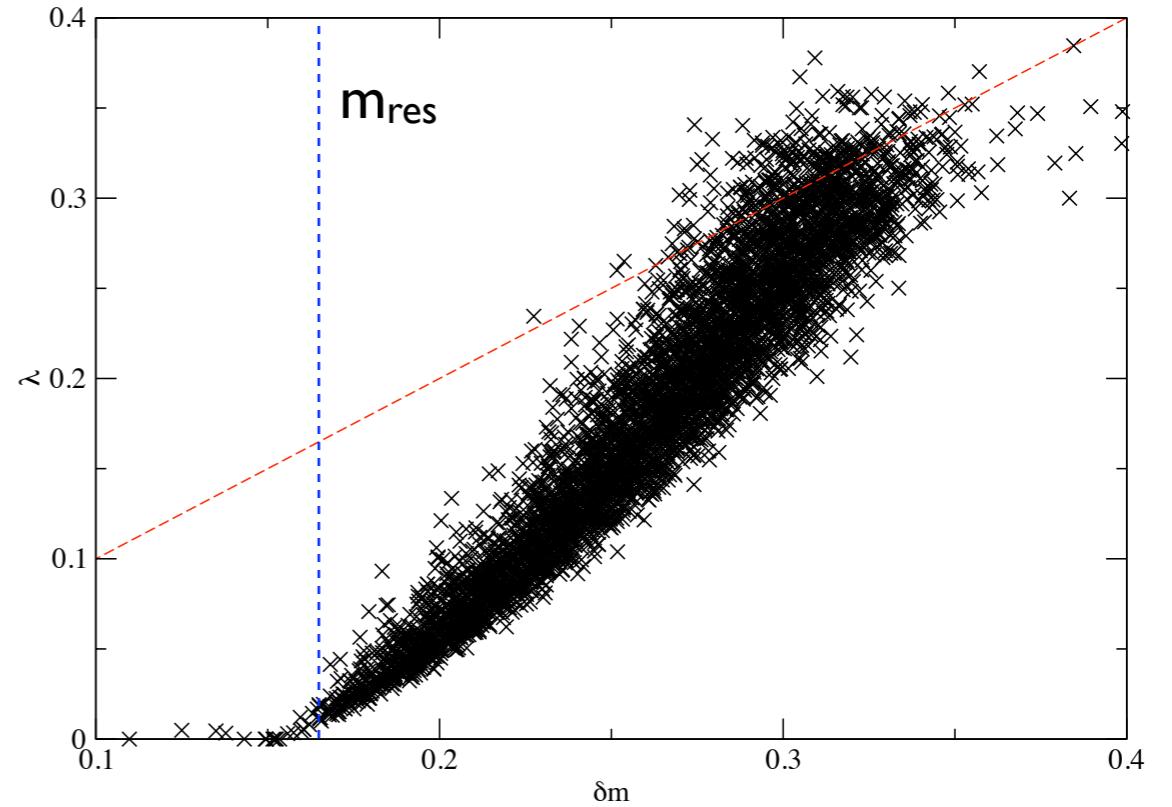
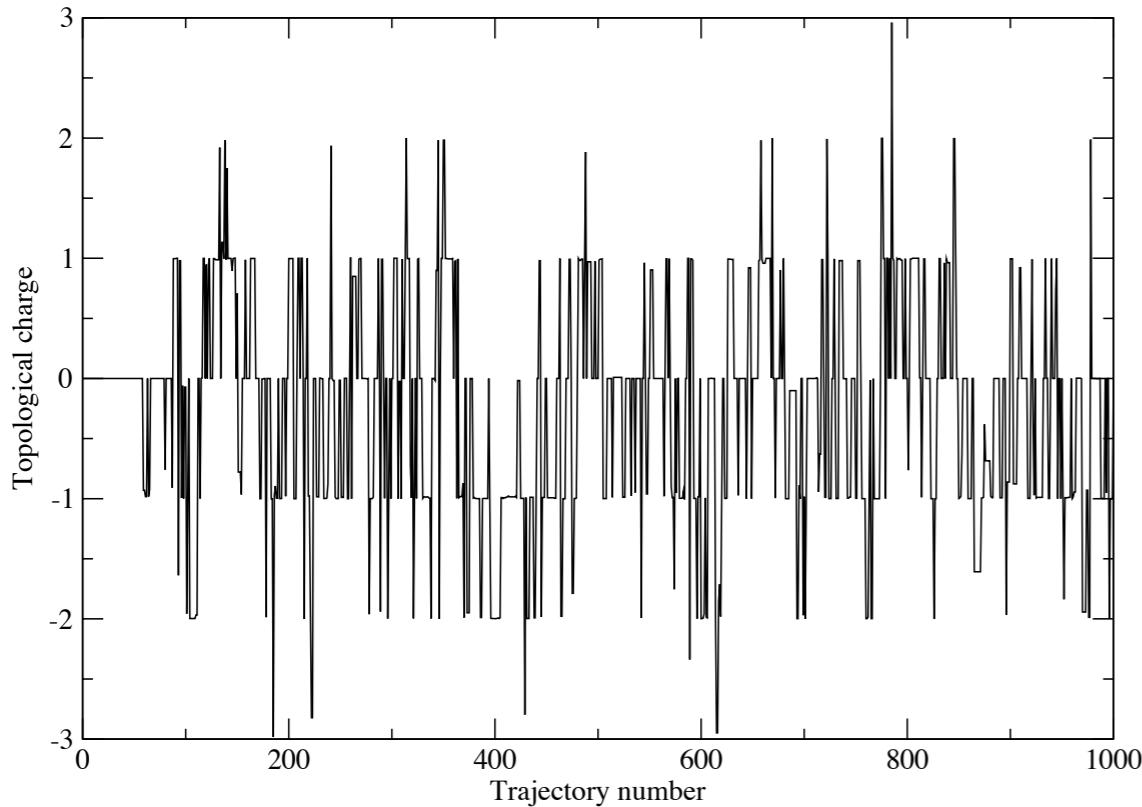
$\beta=2.3$, $m_f=0.2$ and $L_s=24$



- No evidence for $\pm \Lambda_H$ pairing suggests large lattice spacing
- $m_f + \delta m \gg \lambda$, consistent with “physical” γ_5 results
- Residual mass cannot be extracted from δm distribution
- Few near zero-modes

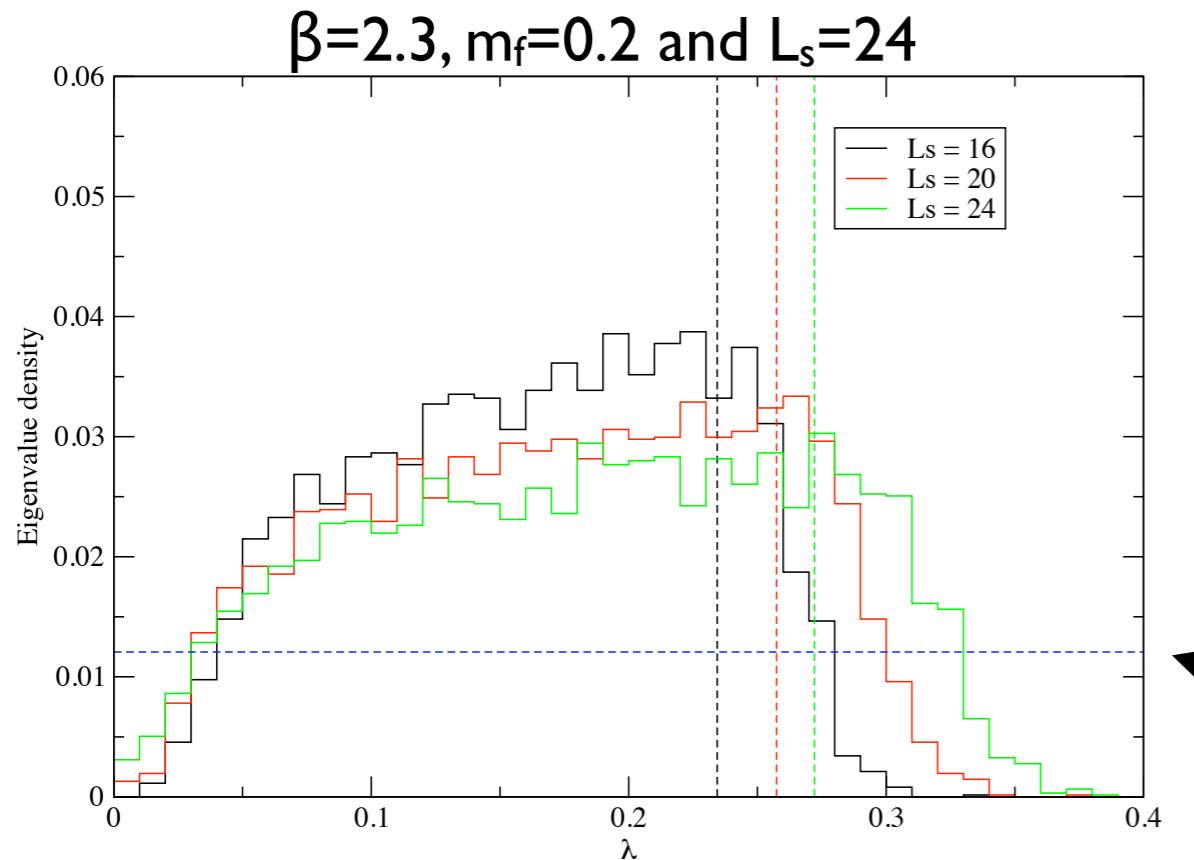
Near zero-modes and topological charge

$\beta=2.3$, $m_f=0.2$ and $L_s=24$



- Gluonic definition of topological charge for same ensemble suggest that there is topological ergodicity
- Nonzero topological charge is poorly reflected in analysis of eigenvalues

Banks-Casher relation for DWFs



$$\rho'(\lambda; m_g) = \left\langle \sum_i \delta(\lambda - \lambda_i) \right\rangle$$

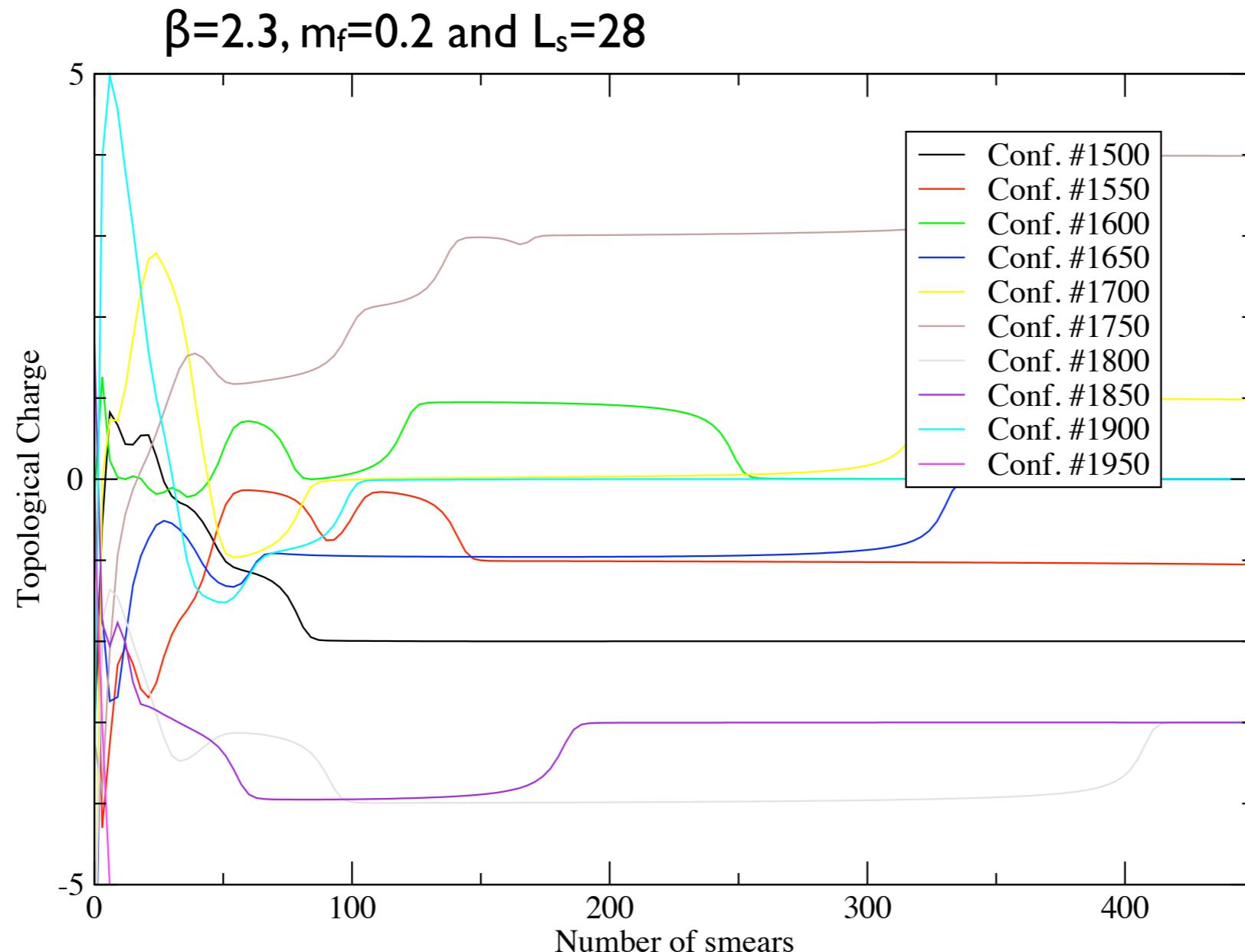
chiral limit extrap. of $\langle \bar{\Psi} \Psi \rangle$
(normalized appropriately)

- Large disagreement in density of eigenvalues and chiral limit value of the condensate
- Moderate L_s dependence of $\rho(\lambda)$ for $\lambda \sim 0.1 - 0.2$
- Spectrum too distorted by large residual chiral symmetry breaking effects

Topological charge

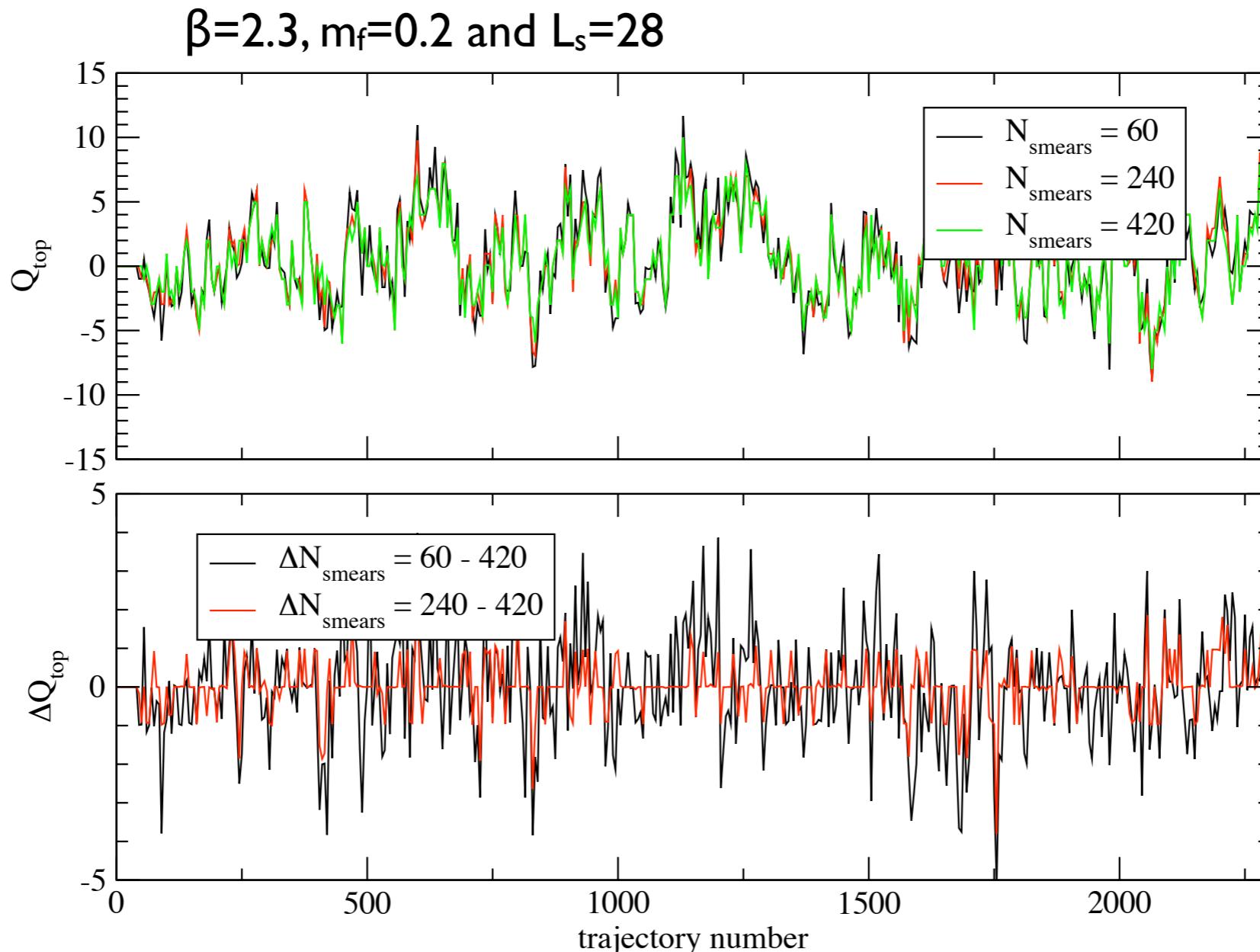
- Smoothed gauge fields using APE smearing (smearing coefficient equal to 0.45)
- Measured topological charge using two methods:
 - 5-loop improved (5li) definition of charge, $O(a^4)$ classically improved field strength tensor Ph. de Forcrand, M. Garcia Perez and I.-O. Stamatescu
 - fermionic definition based on stochastic estimate of $\langle \psi \gamma_5 \psi \rangle$

Topological charge



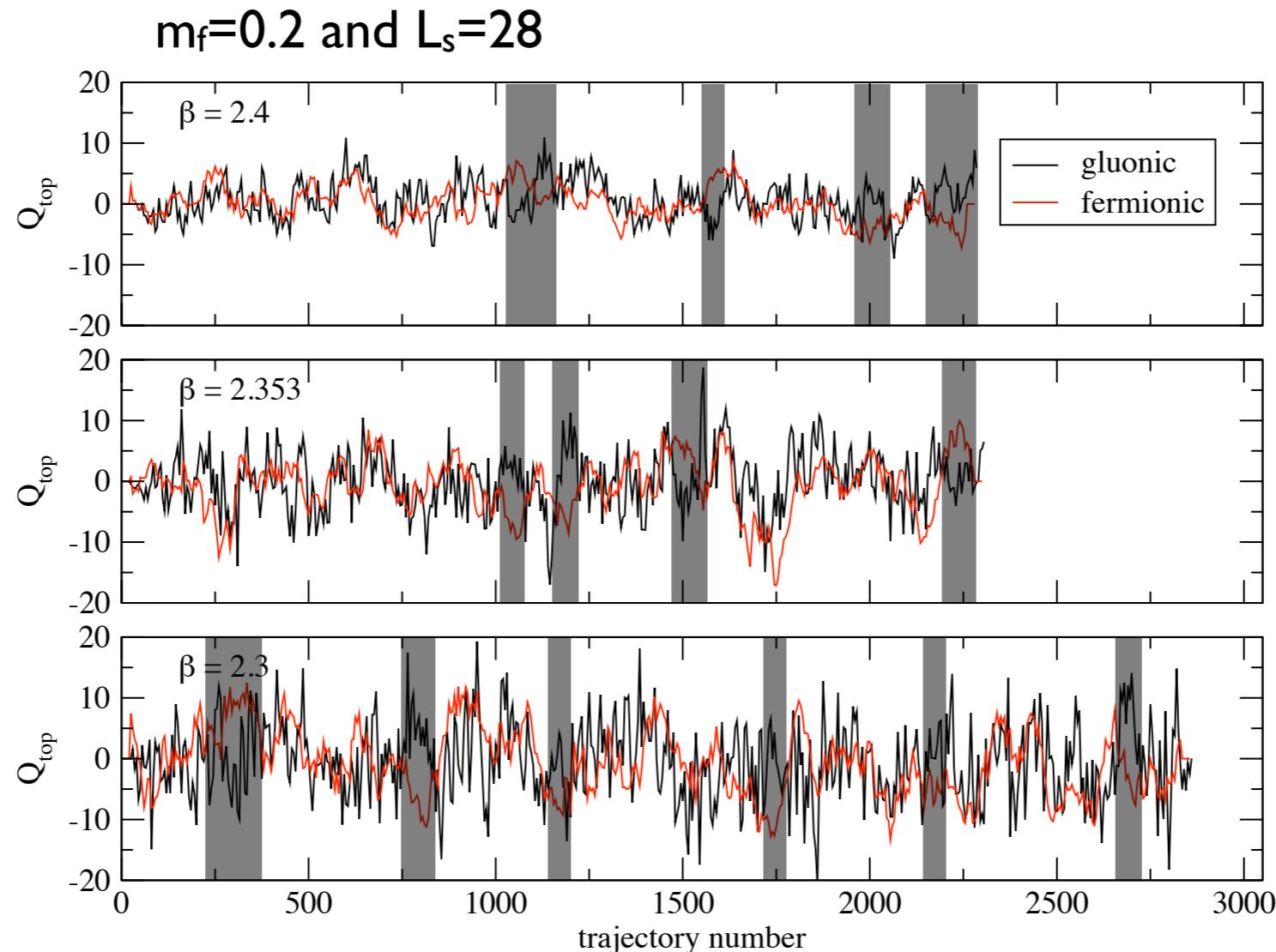
- Q_{top} fluctuation are a result of smearing away of small instantons
- Adequate for qualitative results, but should probably consider cooled lattices in future studies

Topological charge



- Qualitative picture of topological charge is independent of N_{smear}
- Variation in topological charge with $N_{\text{smear}} \sim 1-2$ on average

Topological charge



- Gluonic and fermionic definitions seem track reasonably well with each other at weakest couplings
- hints of an index theorem at weaker coupling?

Summary of results

- Current DWF SYM results are generally disappointing due to large residual chiral symmetry breaking
 - large systematic errors in chiral limit extrapolations
 - evident from Dirac spectrum studies and matrix elements of the “physical” chirality operator
- Substantial reduction in residual chiral symmetry breaking is crucial for reliable study of $N=1$ SYM
- Nonetheless, these simulations are important for understanding the relationship between the input parameters of the lattice action and physical quantities, which will guide future studies.

Future tasks and directions

- Combine improved action, larger L_s and weaker coupling to achieve smaller residual chiral symmetry breaking
 - Iwasaki gauge action
 - AuxDet DWFs recently implemented in Columbia Physics System (D. Renfrew) and modified for use with adjoint representation fermions (MGE)
- AuxDet simulations will be performed on New York Blue and QCDOC at Brookhaven National Laboratory in the near future
 - parameter tuning in progress... stay tuned.