# Introduction to lattice supersymmetry and its applications

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 $\begin{array}{l} \mbox{Outline}\\ \mbox{Goals, Methods}\\ \mbox{DWF for }\mathcal{N}=1\\ \mbox{Super QCD}\\ \mbox{Exact lattice SUSY}\\ \mathcal{N}=4 \mbox{SYM and AdS/CFT} \end{array}$ 

Goals, Methods

DWF for  $\mathcal{N}=1$ 

Super QCD

Exact lattice SUSY

 $\mathcal{N}=4$  SYM and AdS/CFT

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SUSY plays a key role in many theories of BSM physics

- Understanding soft SUSY breaking in MSSM in terms of strongly coupled high scale SQCD
- Extra dims models
- SUSY technicolor
- AdS/CFT duals, strings, black holes
- Many aspects are inherently non-perturbative eg dynamical SUSY breaking, gaugino condensation – lattice natural tool.

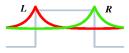
 $\begin{tabular}{c} \hline & & Outline \\ \hline & & Goals, Methods \\ DWF for $\mathcal{N}=1$ \\ Super QCD \\ Exact lattice SUSY \\ $\mathcal{N}=4$ SYM and AdS/CFT \end{tabular}$ 

#### Problems, solutions

- SUSY broken. Large amount fine tuning in general to take *a* → 0.
- But, some cases evaded/reduced:
  - ► 4D N = 1 SYM: chiral symmetry prohibits gaugino mass. Hence DWF/overlap good approach.
  - ► 4D N = 4 SYM: supersymmetric lattice action exists. (new ideas twisting, orbifolding, Kähler-Dirac fermions ).
  - Stepping stones to other theories: super QCD, deformations of *N* = 4 with eg mass terms, AdS/CFT duals

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# Why DWF good for $\mathcal{N}=1$



- ▶ 5D domain wall: 4d chiral, massless mode localized to wall.
- Can be gauged and subject to Majorana condition real, positive definite Pfaffian (c.f Wilson)
- Only relevant SUSY breaking op. is gaugino mass (Veneziano et al)
- Good chiral properties of DWF SUSY broken but no fine tuning a → 0 ...

But more expensive ... See M. Endres talk < -> < -> < -> < -> < ->

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### History

- ▶ Fleming, Kogut, Vranas in 2000:  $8^4 \times L_s$  lattices,  $\beta = 2.3$ ,  $L_s = 12 24$
- Recently 2008-9:
  - S.C,Brower,Fleming,Giedt,Vranas arXiv:0807.2032, arXiv:0810.5746 (16<sup>3</sup> × 32 × L<sub>s</sub> with L<sub>s</sub> = 16 − 64, β = 2.3, 2.4), (1Tflopyear)
  - ▶ M. Endres arXiv:0902.4267,  $(8^4/16^3 \times 32 L_s = 16 28 \beta = 2.3, 2.35, 2.4)$
- ▶ Both groups use hacked CPS code for SU(2). Run on BG/L

### Results so far

Broad agreement between 2 recent calcs. Show:

- Confining as expected static potential, string tension.
- Finite volume effects under control.
- Estimate cut-off effects now.
- Much better extrapolations to chiral limit measure residual mass - but large ...

Strong evidence for nonzero 
$$< ar{\lambda} \lambda >$$
 as  $a 
ightarrow 0.$ 

#### Future

- ► Large m<sub>res</sub> = \frac{\rho(0)}{L\_s} + ... Need larger L<sub>s</sub>. Or better DWFs (möbius ?,..)
- Non-perturbative renormalization to give physical condensate.
- Spectrum ? Interpolating operators known. But 1 Majorana fermion – disconnected diagrams. Hard. cf η' in QCD.
- Eg. 32<sup>3</sup> × 64 with L<sub>s</sub> = 64 − 128 at 3 β's would need at least 100 Tflopyear.

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# Super QCD - I

- MSSM weakly coupled. But soft SUSY breaking ops typically generated by strongly coupled SQCD like sector at high scales.
- Ingredients: add chiral multiplets (scalar+fermion in fundamental)
- DWF eliminates need to tune fermion masses (Giedt et al. arXiv:0806.0013)
- ▶ Fine tuning Yukawas swap for scalar kinetic terms. Flavor symmetries imply single  $m_{\phi}^2$
- ► Quartic ops still an issue n = (2 + few) such ops. depending on N<sub>f</sub>, N<sub>c</sub>.

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# Super QCD - II

- Find SUSY pt by requiring < ∂<sub>µ</sub>S<sub>µ</sub>(x)Θ<sup>i</sup>(0) >= 0 for set of set ops Θ<sup>i</sup> with i ≥ n number of ops.
- If we want to use multicanonical reweighting techniques need say 10 simulation pts for each fine tuned coupling.
- ► Implies need 10<sup>2+few</sup> more CPU than current N = 1 DWF simulations.
- Note: use for N = 4 SYM independent of supersymmetric lattice formulations

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# SUSY lattices

- ► Recent work: possible to discretize theories with Q ≥ 2<sup>D</sup> supercharges while preserving one or more SUSYs exactly (S.C, David Kaplan, Mithat Ünsal, Phys. Rep, arXiv:0903.4881)
- Simplest way to understand: discretizations of topologically twisted versions of target SYM theory
- In flat space: these twisted theories completely equivalent to usual theory.
- ▶ In particular lead to unique lattice formulation of N = 4 SYM with one exact SUSY.

## Exact SUSY for $\mathcal{N} = 4$ : pros and cons

- Exact SUSY ensures boson/fermion spectrum paired and *E*<sub>vac</sub> = 0.
- ▶ Reduces number of potential counterterms needed to take a → 0.
- Action is local, free of doublers, gauge invariant and supersymmetric.
- Price (mostly technical):
  - Most natural lattice is not hypercubic A<sup>\*</sup><sub>4</sub>
  - 16 Fermion fields distributed over links of lattice fermion action non-standard (Kähler-Dirac).
  - Bosons 5 Complex gauge links  $U_a = e^{(A_a + iB_a)}$
  - Sign problem ?

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#### Lattice $\mathcal{N} = 4$

#### Action:

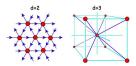
$$S = \sum_{\mathbf{x}} \operatorname{Tr} \left( \mathcal{F}_{mn}^{\dagger} \mathcal{F}_{mn} + \frac{1}{2} \left( \overline{\mathcal{D}}_{m}^{(-)} \mathcal{U}_{m} \right)^{2} - \left( \eta \overline{\mathcal{D}}_{m}^{(-)} \psi_{m} + \chi_{mn} \mathcal{D}_{[m}^{(+)} \psi_{n]} \right) \right)$$
  
+  $S_{\text{closed}}$ 

Fields defined on  $A_4^*$  with

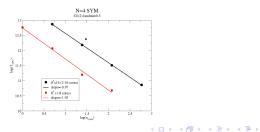
$$\mathcal{F}_{mn} = \mathcal{U}_m(x)\mathcal{U}_n(x+m) - \mathcal{U}_n(x)\mathcal{U}_m(x+n)$$
$$S_{\text{closed}} = -\frac{1}{8}\sum_{\mathbf{x}} \operatorname{Tr} \epsilon_{mnpqr} \chi_{qr}(\mathbf{x} + \mu_m + \mu_n + \mu_p)\overline{\mathcal{D}}_p^{(-)} \chi_{mn}(\mathbf{x} + \mu_p)$$

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#### Lattices, code



 $8^3 \times 16$  2-16 cores Jpsi (MPI comm using MDP in FermiQCD)



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### Renormalization

Lattice symmetries:

- Gauge invariance
- Q-symmetry.
- Point group symmetry eg.  $S^5$  for  $A_4^*$
- Exact fermionic shift symmetry  $\eta \rightarrow \eta + \epsilon I$

Conclusion:

 Only (marginally) relevant ops that can arise correspond to renormalizations of ops. in bare theory except for SUSY mass term Q(U<sub>a</sub>U<sup>†</sup><sub>a</sub>η)

Examine flows at 1-loop - in progress (with J. Giedt). Currently appears that no mass term induced!

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#### Lattice perturbation theory

Ingredients:

- ► Bosons:  $\langle \overline{\mathcal{A}}_a(k)\mathcal{A}_b(-k) \rangle = \frac{1}{\hat{k}^2}\delta_{ab}$  with  $\hat{k}^2 = 4\sum_a \sin^2(k_a/2)$
- Fermions:  $M_{\text{KD}}^{-1}(k) = \frac{1}{\hat{k}^2} M_{\text{KD}}(k)$  with M(k) a 16 × 16 block matrix.
- Vertices:  $\psi\eta$ ,  $\psi\chi$  and  $\chi\chi$ . No lattice tadpoles.
- Small number of Feynman graphs determine wavefunction renormalization from 1-loop corrections to fermion propagators.

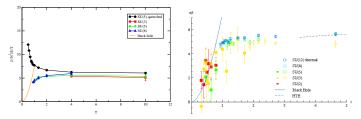
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# Gravity duals

- Large number of examples of conjectured dualities between (supersymmetric) YM and gravity/string theories.
- YM always strongly coupled –
   New SUSY lattice actions very useful in this respect.
- Simplest example studied so far:
   SYMQM at large N black hole type IIa SUGRA
   S.C and Wiseman arXiv:0803.4273, Nishimura et al. arXiv:0810.2884
- Current work:
  - $\mathcal{N} = 4$  in D = 4 AdS/CFT, (SUSY)Wilson loops
  - Dimensional reduction to 2D Is Gregory-LaFlamme black string-black hole transition dual to thermal PT in YM ?

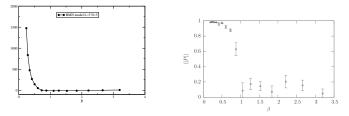
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#### Black holes from YM



Energy vs temperature for SYMQM system. Black hole prediction shown (no fit) Care ... divergence in thermal partition function - needs care ...

#### Mass deformed SYMQM – BMN model



Thermal phase transition between SUSY confining phase at low T and deconfined broken phase at high T (Relation to Hawking-Page transition in  $\mathcal{N} = 4$ ?)

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### Summary

Several promising approaches to SUSY on lattice

- Use DWF+reweighting  $\mathcal{N} = 1$  SYM and SQCD. Ultimate goal to understand soft terms in MSSM.
  - Break SUSY but ensure restored as  $a \rightarrow 0$  with minimum fine tuning
  - Price is (some) reweighting scalar sector and need DWF.
- Formulations with exact SUSY.
  - Reduced fine tuning due to lattice symmetries.
  - Lots of applications to AdS/CFT-like duals, understand quantum black holes, ...
  - Potential sign problem ?