### Peskin Takeuchi S-parameter on 2+1 QCD DW lattices<sup>†</sup>

Rich Brower

OEstimates for "technicolor" are base on scaling fromQCD.

•(Modest) goal: Evaluate S for QCD (or minimal technicolor).

Establish lattice methods

•Test phenomenological rationale for S parameter models.

• Other uses of Euclidean correlators:  $\Pi(Q^2) = FT < J_{\mu}(x) J_{\mu}(0) >$ 

Hadronic contribution to g-2

\*Determine "warping" for AdS/CFT models:  $Q \simeq 1/z$ 

• See Ami Katz (Inverse problems unique at  $N_c = \infty$ ?)

OComment on Application of Disconnected project to SUSY dark matter

•Largest uncertainty is Nucleon strange sigma term:  $m_s <N|s s|N>/M_N$ 

#### <sup>†</sup> LSD collaboration using QCDOC lattices



$$i \int dx e^{iqx} \langle J^R_\mu(x) J^L_\nu(0) \rangle = (q_\mu q_\nu - q^2 g_{\mu\nu}) \Pi(q^2)$$

S parameter<sup>†</sup> = 
$$S_{strong} - S_{SM}$$

$$S(m_{H,ref}) = 4 \int_{\epsilon}^{\infty} \frac{dq^2}{q^2} \left[ Im \Pi_{VV}(q^2) - Im \Pi_{AA}(q^2) \right] - \frac{1}{12\pi} \int_{\epsilon}^{\infty} \frac{dq^2}{q^2} \left[ (1 - (1 - m_{H,ref}/q^2)^3 \theta(q^2 - m_{H,ref}^2)) \right]$$

Using Dispersion Rel:  $q^2 \Pi(q^2) = q^2 \int \frac{ds}{\pi} \frac{Im \Pi_{VV}(s) - Im \Pi_{AA}(s)}{s - q^2} - f_{\pi}^2$ 

 $\begin{array}{ll} \textit{Naive Scaling Models.} \\ \rho/\text{A}_1 \left( \text{Weinberg sum rules} \right) ==> & S \simeq 0.25 \frac{N_f}{2} \frac{N_c}{3} \\ \chi \ \text{PT} \left( \text{leading logs} \right) & ==> & S \simeq \frac{1}{12\pi} [\frac{N_f^2}{4} - 1] \ln(\frac{m_\rho^2}{m_{PGB}^2}) \end{array}$ 

<sup>+</sup> "Estimation of oblique electroweak corrections" Peskin and Takeuchi, PR D46 (1992) 381

But "larger"  $n_f$  dependence of  $S(N_c, n_f)$  must be non-trivial.

(An IR fixed point at the two-loop  $\beta$  function. Gross and Wilczek, Banks and Zaks, ...)

•Conformal:  $n_f > n_f^*$ 

•Walking:  $n_f < n_{f}^*$ 



Pioneering work on  $a^{had}_{\mu} = (g-2)/2$ 



$$a^{(had)}{}_{\mu} = 460(78) \ge 10^{-10}$$

 $a^{(had)}_{\mu} = 713(15) \times 10^{-10}$ 

JLQCD<sup>†</sup> calculation of EM  $\Delta m_{\pi}^2$  and  $Q^2 \Pi_{VA}$  for S (Overlap Fermions using non-conserved currents with  $Z_A = Z_V = 1.38$ )



 $L_{10}(\mu) = -0.00474(23) \text{ vs exp } -0.00509(47)$ 200 config. 16<sup>3</sup> x 32 lattice.,  $m_{\pi}^2 = 0.27 - 0.082 \text{ Gev}^2$ 

<sup>†</sup> Shintani et al PoS (Lattice 2007) 137 using Dass-Guralnik-Mathur-Low-Young sum rule.

# Domain Wall Currents



#### Need for Pauli-Villars Subtraction

• Overlap Derivation: 
$$\langle J_{\mu}(y)J_{\mu}(x)\rangle \equiv \frac{\delta}{\delta A_{\mu}(y)} \frac{\delta}{\delta A_{\mu}(x)} \log[Zov(Ue^{iA})]$$
 at  $A_{\mu} = 0$   
• Overlap ==> DomainWall:  $\log[Z_{0v}] = Tr \log[D_{ov}(m)] \equiv \log[Z_{DW}] = Tr \log[D_{DW}(m)] - Tr \log[D_{DW}(1)]$   
• Equivalent DW expression:  
 $\langle J_{\mu}(y)J_{\mu}(x)\rangle \equiv \sum_{s',s} \langle j_{\mu}^{DW}(y,s')j_{\mu}^{DW}(x,s) - j_{\mu}^{PV}(y,s')j_{\mu}^{PV}(x,s)\rangle$ 

•With  $\Gamma_s = \text{Sign}(L_s/2 + 1 - s)$ 

$$\langle J^{5}_{\mu}(y)J^{5}_{\mu}(x)\rangle \equiv \sum_{s',s} \Gamma_{s}\Gamma_{s'}\langle j^{DW}_{\mu}(y,s')j^{DW}_{\mu}(x,s) - j^{PV}_{\mu}(y,s')j^{PV}_{\mu}(x,s)\rangle$$

•BUT cheaper to use local Non-conserved currents on boundary with  $Z_V = Z_A + O(m^2_{res})$ 

# SUSY Dark Matter



ONeed to compute disconnected diagram for strange loop

 Very expensive but new techniques<sup>†</sup> (dilution, multi-grid variance reduction, finite volume extrapolation) are working.

"Determination of Dark Matter Properties at High Energy Colliders" Balitz, Battaglia, Peskin and Wizansky hep-ph./0602187v4

- "Hadronic Uncertainty in the Elastic Scattering of Supersymmetric Dark Matter" Ellis, Olive and Savage hep-ph/0801.3656v2 "
- *†* "Strange quark contribution to nucleon form factors", Ronald Babich, Richard Brower, Michael Clark, George Fleming, James Osborn, Claudio Rebbi hep-lat/ 0710.5536

# Strange Quark Disconnected Diagram Project



OStrange sigma term in Nucleon  $R_s = \langle N(t) Tr[D^{-1}(0)] N(-t) \rangle \langle N(t) N(-t) \rangle$ 

• In the numerator fit to two forward states transition term. In the denominator fit two forward states (blue with  $c_B = 0$ ) plus one backward state (red).  $R_s = \frac{c_1^2 j_1^2 e^{-2m_1 t} + c_1 c_2 j_{12} e^{-(m_1 + m_2)t} + c_2^2 j_2^2 e^{-2m_2 t}}{c_1^2 e^{-2m_1 t} + c_2 e^{-2m_2 t} + c_R^2 e^{-2m_B (L/2-t)}}$