

The Conformal Window in QCD-like Theories

Ethan T. Neil, George T. Fleming, Thomas Appelquist
ethan.neil@yale.edu

Department of Physics
Yale University

May 2, 2008

Outline

- 1 Introduction
 - Motivation
 - Flavor dependence
- 2 Setup and Methods
 - Program of study
 - Lattice methods
 - Schrödinger Functional
 - Continuum extrapolation
- 3 Results
 - Results
 - Conclusion

Motivation and Introduction

- Many of the proposed theories for electroweak symmetry breaking that might be observable at the LHC involve new **strong dynamics**: technicolor, topcolor, composite Higgs...all involve additional strongly coupled $SU(N)$ Yang-Mills sectors.

Motivation and Introduction

- Many of the proposed theories for electroweak symmetry breaking that might be observable at the LHC involve new **strong dynamics**: technicolor, topcolor, composite Higgs...all involve additional strongly coupled $SU(N)$ Yang-Mills sectors.
- Also, Yang-Mills theory with many flavors is approximately conformal at long distances; conformal or near-conformal infrared behavior appears in many other ideas for physics beyond the Standard Model (AdS/CFT, unparticles, walking technicolor...)

Motivation and Introduction

- Many of the proposed theories for electroweak symmetry breaking that might be observable at the LHC involve new **strong dynamics**: technicolor, topcolor, composite Higgs...all involve additional strongly coupled $SU(N)$ Yang-Mills sectors.
- Also, Yang-Mills theory with many flavors is approximately conformal at long distances; conformal or near-conformal infrared behavior appears in many other ideas for physics beyond the Standard Model (AdS/CFT, unparticles, walking technicolor...)
- We can get a lot out of studying Yang-Mills theory!

Motivation and Introduction

- Many of the proposed theories for electroweak symmetry breaking that might be observable at the LHC involve new **strong dynamics**: technicolor, topcolor, composite Higgs...all involve additional strongly coupled $SU(N)$ Yang-Mills sectors.
- Also, Yang-Mills theory with many flavors is approximately conformal at long distances; conformal or near-conformal infrared behavior appears in many other ideas for physics beyond the Standard Model (AdS/CFT, unparticles, walking technicolor...)
- We can get a lot out of studying Yang-Mills theory!
- Fix $N = 3$, N_f fermions in the fundamental rep.

Motivation and Introduction

- Many of the proposed theories for electroweak symmetry breaking that might be observable at the LHC involve new **strong dynamics**: technicolor, topcolor, composite Higgs...all involve additional strongly coupled $SU(N)$ Yang-Mills sectors.
- Also, Yang-Mills theory with many flavors is approximately conformal at long distances; conformal or near-conformal infrared behavior appears in many other ideas for physics beyond the Standard Model (AdS/CFT, unparticles, walking technicolor...)
- We can get a lot out of studying Yang-Mills theory!
- Fix $N = 3$, N_f fermions in the fundamental rep.
- “Lattice Study of the Conformal Window in QCD-like Theories” (Thomas Appelquist, George T. Fleming, EN.) **PRL 100, 171607 (2008)**.

Flavor dependence of $SU(3)$ Yang-Mills

- The running coupling g of QCD is characterized by two important features: **asymptotic freedom** ($g \rightarrow 0$ in the UV) and **confinement** ($g \rightarrow \infty$ in the IR.)

Flavor dependence of $SU(3)$ Yang-Mills

- The running coupling g of QCD is characterized by two important features: **asymptotic freedom** ($g \rightarrow 0$ in the UV) and **confinement** ($g \rightarrow \infty$ in the IR.)
- These properties are strongly dependent on the number of fermion flavors, N_f :

	Short-distance (UV)	Long-distance (IR)
$0 < N_f < N_f^c$	Free ($g \rightarrow 0$)	Confined ($g \rightarrow \infty$)
$N_f^c < N_f < 16.5$	Free ($g \rightarrow 0$)	Fixed point ($g \rightarrow g^*$)
$N_f > 16.5$	Divergent ($g \rightarrow \infty$)	Trivial ($g \rightarrow 0$)

Flavor dependence of $SU(3)$ Yang-Mills

- The running coupling g of QCD is characterized by two important features: **asymptotic freedom** ($g \rightarrow 0$ in the UV) and **confinement** ($g \rightarrow \infty$ in the IR.)
- These properties are strongly dependent on the number of fermion flavors, N_f :

	Short-distance (UV)	Long-distance (IR)
$0 < N_f < N_f^c$	Free ($g \rightarrow 0$)	Confined ($g \rightarrow \infty$)
$N_f^c < N_f < 16.5$	Free ($g \rightarrow 0$)	Fixed point ($g \rightarrow g^*$)
$N_f > 16.5$	Divergent ($g \rightarrow \infty$)	Trivial ($g \rightarrow 0$)

- Theories in the second row of the table lie in the **conformal window**. Perturbation theory gives some insight, but near N_f^c the fixed point coupling is too strong \Rightarrow non-perturbative study is required.

Estimates of N_f^C

- Continuum study based on counting degrees of freedom (Appelquist, Cohen, Schmaltz 1999) yields a bound:

$$N_f^C \leq 4N \left(1 - \frac{1}{18N^2} + \dots \right)$$

Estimates of N_f^C

- Continuum study based on counting degrees of freedom (Appelquist, Cohen, Schmaltz 1999) yields a bound:

$$N_f^C \leq 4N \left(1 - \frac{1}{18N^2} + \dots \right)$$

- Gap equation studies (Appelquist et al, PRD 58: 105017, 1998) suggest that this bound is saturated, i.e. for $N = 3$, $N_f^C \approx 12$.

Estimates of N_f^c

- Continuum study based on counting degrees of freedom (Appelquist, Cohen, Schmaltz 1999) yields a bound:

$$N_f^c \leq 4N \left(1 - \frac{1}{18N^2} + \dots \right)$$

- Gap equation studies (Appelquist et al, PRD 58: 105017, 1998) suggest that this bound is saturated, i.e. for $N = 3$, $N_f^c \approx 12$.
- In supersymmetric $SU(N)$ Yang-Mills, the ACS inequality yields $N_f^c \leq 3N/2$; Seiberg duality can be used to show the bound is saturated, $N_f^c = 3N/2$.

Estimates of N_f^c

- Continuum study based on counting degrees of freedom (Appelquist, Cohen, Schmaltz 1999) yields a bound:

$$N_f^c \leq 4N \left(1 - \frac{1}{18N^2} + \dots \right)$$

- Gap equation studies (Appelquist et al, PRD 58: 105017, 1998) suggest that this bound is saturated, i.e. for $N = 3$, $N_f^c \approx 12$.
- In supersymmetric $SU(N)$ Yang-Mills, the ACS inequality yields $N_f^c \leq 3N/2$; Seiberg duality can be used to show the bound is saturated, $N_f^c = 3N/2$.
- However, previous lattice investigation of the conformal window (Iwasaki et al, PRD 69, 014507, 2004) claims the result $6 < N_f^c < 7$.

Program of study

- Goal: check the results of Iwasaki et al by obtaining an independent bound on N_f^c through lattice simulation.

Program of study

- Goal: check the results of Iwasaki et al by obtaining an independent bound on N_f^c through lattice simulation.
- Method: measure the running coupling over a wide range of scales, and look for the existence of an IR fixed point.

Program of study

- Goal: check the results of Iwasaki et al by obtaining an independent bound on N_f^c through lattice simulation.
- Method: measure the running coupling over a wide range of scales, and look for the existence of an IR fixed point.
- Use staggered fermions for computational efficiency, which naturally come in multiples of 4 flavors.

Program of study

- Goal: check the results of Iwasaki et al by obtaining an independent bound on N_f^c through lattice simulation.
- Method: measure the running coupling over a wide range of scales, and look for the existence of an IR fixed point.
- Use staggered fermions for computational efficiency, which naturally come in multiples of 4 flavors.
 - $N_f = 4$: expected to be well into the broken phase
 - $N_f = 8$: presence of IRFP unknown
 - $N_f = 12$: should be in the conformal window
 - $N_f = 16$: very perturbative IR fixed point

Program of study

- Goal: check the results of Iwasaki et al by obtaining an independent bound on N_f^c through lattice simulation.
- Method: measure the running coupling over a wide range of scales, and look for the existence of an IR fixed point.
- Use staggered fermions for computational efficiency, which naturally come in multiples of 4 flavors.
 - $N_f = 4$: expected to be well into the broken phase
 - $N_f = 8$: presence of IRFP unknown
 - $N_f = 12$: should be in the conformal window
 - $N_f = 16$: very perturbative IR fixed point

Simulate here!

Measuring the running coupling

- When simulating on the lattice, must work at scales well-separated from the lattice spacing a and the box size L . Normally hard enough, but we want to measure over a huge range of scales!

Measuring the running coupling

- When simulating on the lattice, must work at scales well-separated from the lattice spacing a and the box size L . Normally hard enough, but we want to measure over a huge range of scales!
- To avoid box-size effects, we measure the **Schrödinger Functional** coupling $\bar{g}^2(L)$, which is defined directly at the scale L .

Measuring the running coupling

- When simulating on the lattice, must work at scales well-separated from the lattice spacing a and the box size L . Normally hard enough, but we want to measure over a huge range of scales!
- To avoid box-size effects, we measure the **Schrödinger Functional** coupling $\bar{g}^2(L)$, which is defined directly at the scale L .
- In order to measure the evolution of $\bar{g}^2(L)$, we use the **step scaling** procedure to link together results of simulations at many different lattice spacings a . Measure in discrete steps: $\bar{g}^2(L) \rightarrow \bar{g}^2(2L) \rightarrow \dots$

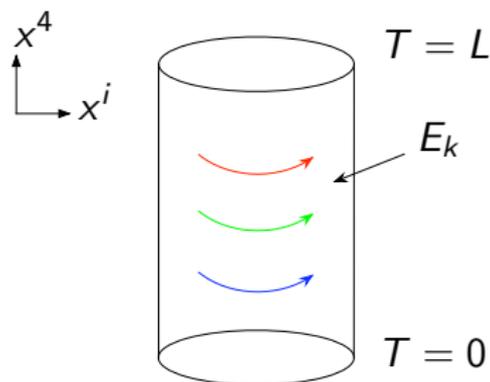
Measuring the running coupling

- When simulating on the lattice, must work at scales well-separated from the lattice spacing a and the box size L . Normally hard enough, but we want to measure over a huge range of scales!
- To avoid box-size effects, we measure the **Schrödinger Functional** coupling $\bar{g}^2(L)$, which is defined directly at the scale L .
- In order to measure the evolution of $\bar{g}^2(L)$, we use the **step scaling** procedure to link together results of simulations at many different lattice spacings a . Measure in discrete steps: $\bar{g}^2(L) \rightarrow \bar{g}^2(2L) \rightarrow \dots$
- Define the **step-scaling function**,

$$\Sigma(2, \bar{g}^2(L), a/L) \equiv \bar{g}^2(2L) + O(a/L)$$

The continuum limit $\sigma(2, u) \equiv \lim_{a \rightarrow 0} \Sigma(2, u, a/L)$ is basically a discretized version of the β -function.

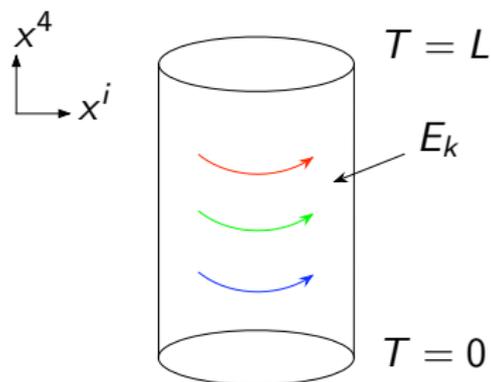
The Schrödinger Functional



Schrödinger Functional

simulations introduce Dirichlet boundaries in time; boundary gauge fields are chosen to give a constant chromoelectric background field.

The Schrödinger Functional



Schrödinger Functional

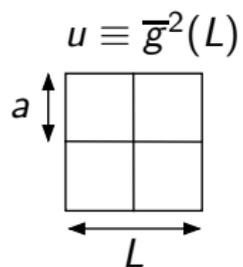
simulations introduce Dirichlet boundaries in time; boundary gauge fields are chosen to give a constant chromoelectric background field.

Running coupling

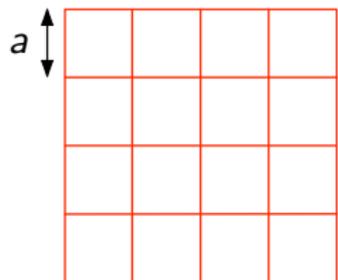
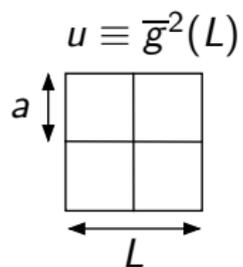
The SF running coupling $\bar{g}^2(L)$ is defined to vary inversely with the response of the action to the strength η of the background field,

$$\frac{dS}{d\eta} = \frac{k}{\bar{g}^2(L)} \Big|_{\eta=0}.$$

Step scaling

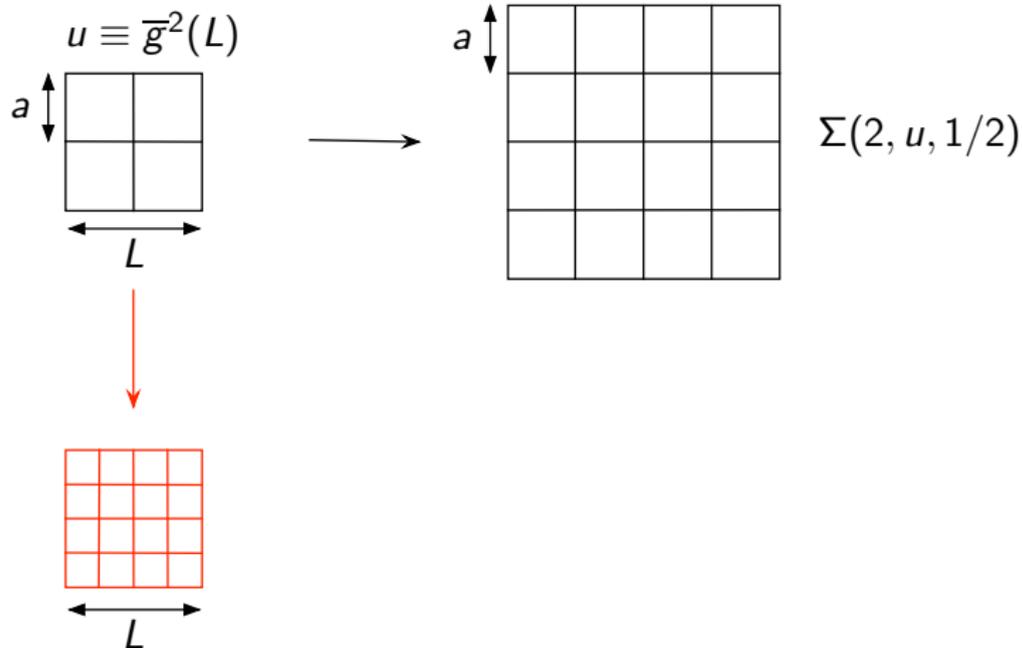


Step scaling

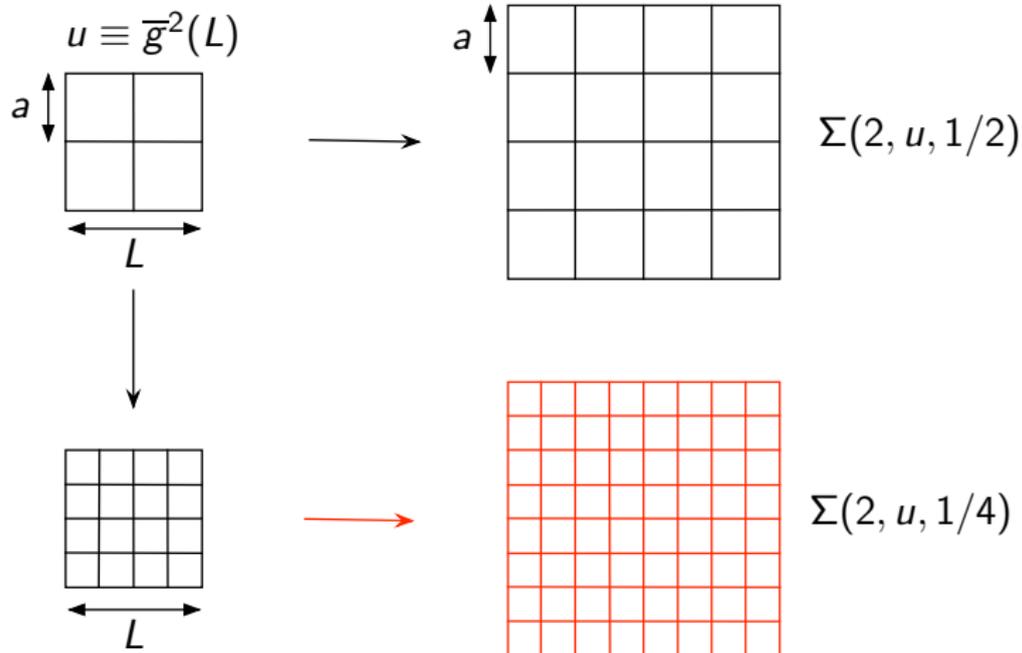


$$\Sigma(2, u, 1/2)$$

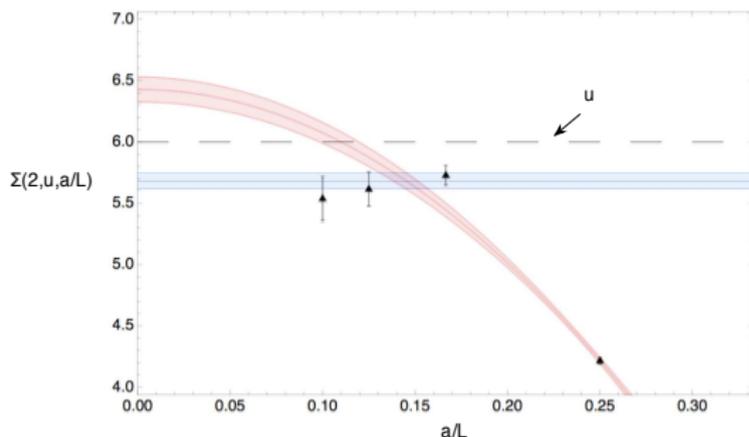
Step scaling



Step scaling

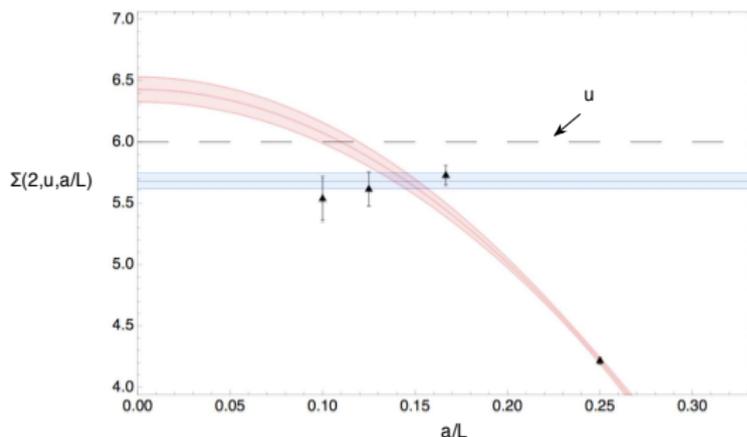


Continuum extrapolation



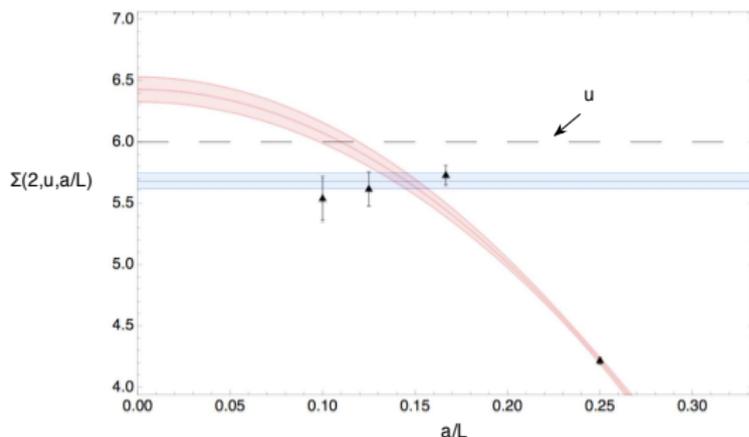
- With a limited number of data points (only even L), continuum extrapolation is poorly constrained.

Continuum extrapolation



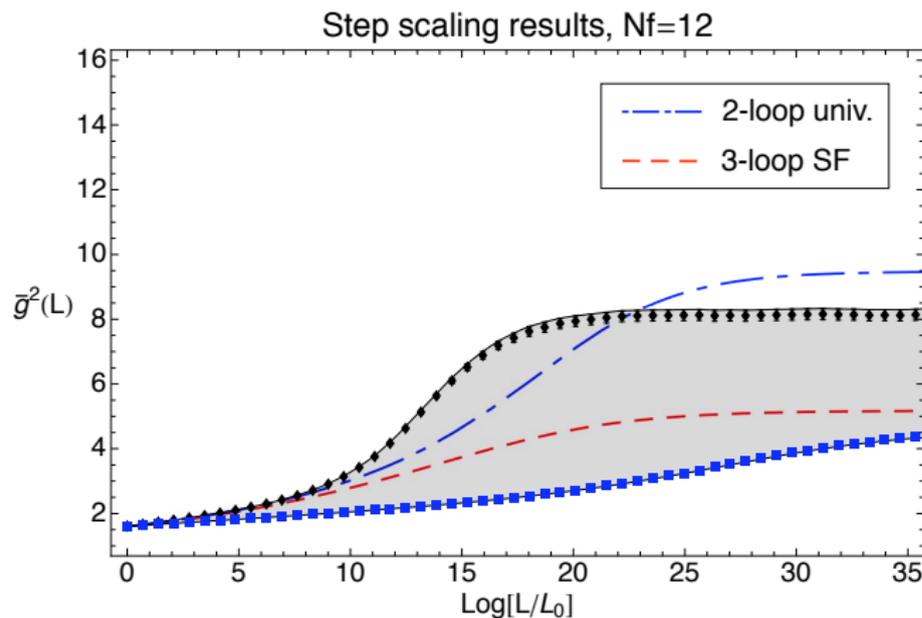
- With a limited number of data points (only even L), continuum extrapolation is poorly constrained.
- Also, some of the high- L data is sparse and not finished running \Rightarrow error bars are likely underestimated.

Continuum extrapolation

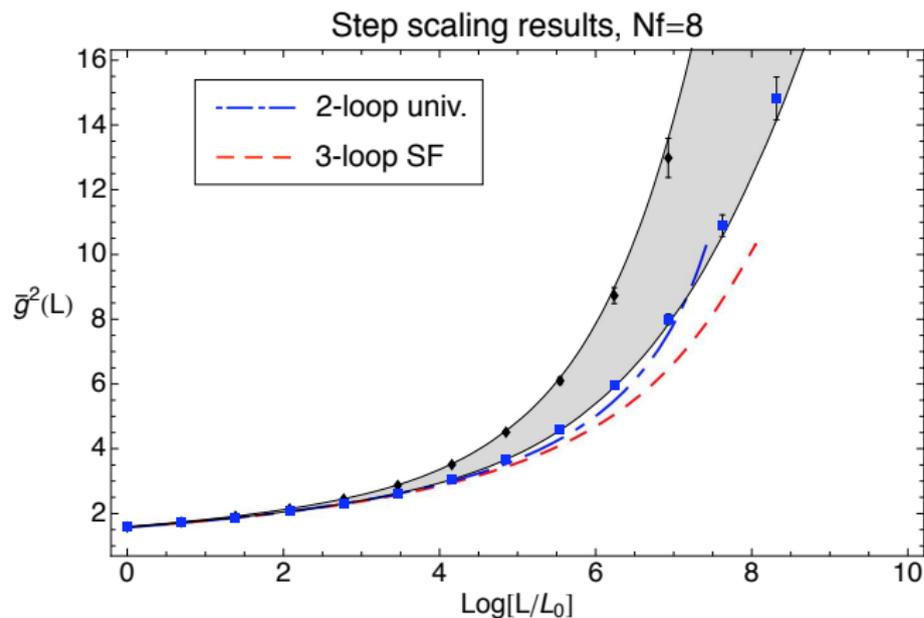


- With a limited number of data points (only even L), continuum extrapolation is poorly constrained.
- Also, some of the high- L data is sparse and not finished running \Rightarrow error bars are likely underestimated.

Any reasonable continuum extrapolation should be bounded by the two methods shown above, so we take them to define a systematic error band.

Results, $N_f = 8$ and 12

Apparent IR fixed point.

Results, $N_f = 8$ and 12

No evidence of a fixed point $\Rightarrow 8 < N_f^c < 12$.

Conclusions

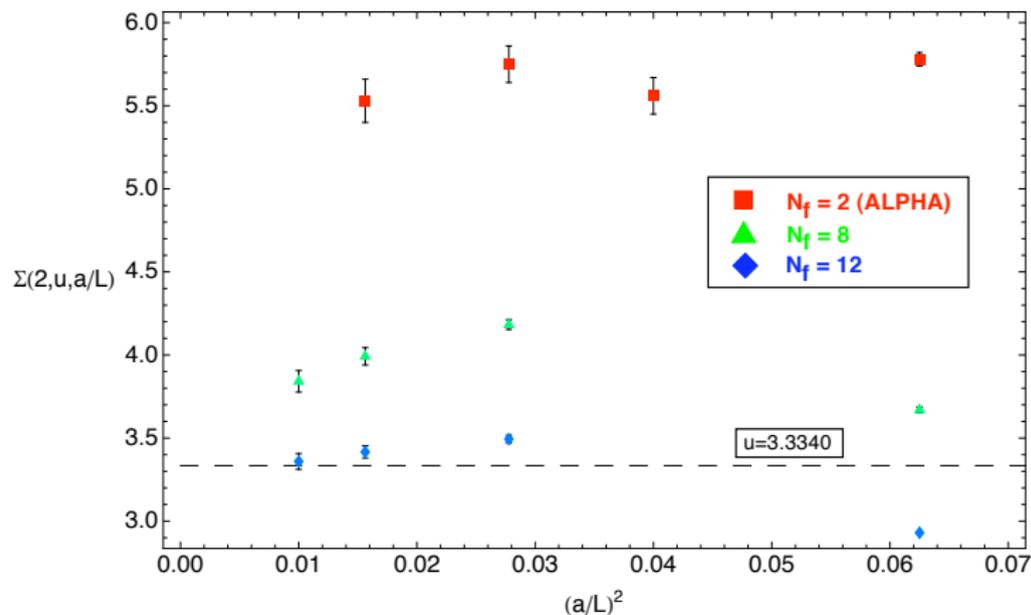
Summary

- We have constrained the lower boundary of the conformal window:
 $8 < N_f^c < 12$.

Future work

- Continued simulations at 8 and 12 flavors, to reduce systematics.
- Study of running coupling at $N_f = 10$ (allocation granted by USQCD)
- Dynamical simulation at $N_f = 8$, to verify the presence of chiral symmetry breaking.
- Simulation at other values of N_c .

Data comparison with ALPHA



(Ref: Della Morte et. al. (ALPHA), hep-lat/0411025, NPB 713 (2005) p.378.)