

MATH423 String Theory Solutions 10

1a.

The nontrivial Lorentz transformations for this boost are:

$$\begin{aligned}x'^0 &= \gamma(x^0 - \beta x^1) \\x'^1 &= \gamma(x^1 - \beta x^0)\end{aligned}$$

We then find

$$\begin{aligned}x'^{\pm} = \frac{x'^0 \pm x'^1}{\sqrt{2}} &= \frac{\gamma}{\sqrt{2}}(x^0 - \beta x^1 \pm x^1 \mp \beta x^0) \\&= \frac{\gamma(1 \mp \beta)}{\sqrt{2}}(x^0 \pm x^1) = \gamma(1 \mp \beta)x^{\pm}\end{aligned}$$

Therefore

$$x'^+ = \sqrt{\frac{1-\beta}{1+\beta}}x^+ \quad , \quad x'^- = \sqrt{\frac{1+\beta}{1-\beta}}x^- \quad , \quad x'^2 = x^2 \quad , \quad x'^3 = x^3 \quad (1)$$

The light cone coordinates x^+ and x^- do not mix under boosts along X^1 !

b.

The new rotated coordinates are given by

$$\begin{aligned}x'^0 &= x^0 \\x'^1 &= \cos \theta x^1 + \sin \theta x^2 \\x'^2 &= -\sin \theta x^1 + \cos \theta x^2 \\x'^3 &= x^3\end{aligned}$$

After some algebra we find

$$\begin{aligned}x'^+ &= \frac{1}{2}(1 + \cos \theta)x^+ + \frac{1}{2}(1 - \cos \theta)x^- + \frac{\sin \theta}{\sqrt{2}}x^2 \\x'^- &= \frac{1}{2}(1 - \cos \theta)x^+ + \frac{1}{2}(1 + \cos \theta)x^- - \frac{\sin \theta}{\sqrt{2}}x^2 \\x'^2 &= -\frac{\sin \theta}{\sqrt{2}}x^+ + \frac{\sin \theta}{\sqrt{2}}x^- + \cos \theta x^2\end{aligned}$$

c. For a boost with velocity parameter β along x^3 , the Lorentz transformations are:

$$\begin{aligned}x'^0 &= \gamma(x^0 - \beta x^3) \\x'^1 &= x^1 \\x'^2 &= x^2 \\x'^3 &= \gamma(x^3 - \beta x^0)\end{aligned}$$

Therefore,

$$\begin{aligned}x'^+ &= \frac{x'^0 + x'^1}{\sqrt{2}} = \frac{\gamma(x^0 - \beta x^3) + x^1}{\sqrt{2}} \\&= \frac{\gamma}{2}(x^+ + x^-) - \frac{\gamma\beta x^3}{\sqrt{2}} + \frac{1}{2}(x^+ - x^-).\end{aligned}$$

Similalry for x'^- , x'^2 and x'^3 we get

$$\begin{aligned}x'^+ &= \frac{1}{2}(\gamma + 1)x^+ + \frac{1}{2}(\gamma - 1)x^- - \gamma\beta\sqrt{2}x^3, \\x'^- &= \frac{1}{2}(\gamma - 1)x^+ + \frac{1}{2}(\gamma + 1)x^- - \gamma\beta\sqrt{2}x^3, \\x'^2 &= x^2 \\x'^3 &= \gamma x^3 - \frac{\gamma\beta}{\sqrt{2}}(x^+ + x^-).\end{aligned}$$

2a.

$$\begin{pmatrix} x \\ ct \end{pmatrix} \sim \begin{pmatrix} x \\ ct \end{pmatrix} + 2\pi \begin{pmatrix} R \\ -R \end{pmatrix}. \quad (2)$$

We use (2) to compute

$$\begin{aligned}x^+ &= \frac{ct + x}{\sqrt{2}} \sim \frac{1}{\sqrt{2}}(ct - 2\pi R + x + 2\pi R) = x^+ \\x^- &= \frac{ct - x}{\sqrt{2}} \sim \frac{1}{\sqrt{2}}(ct - 2\pi R - x - 2\pi R) = x^- - 2\pi(\sqrt{2}R)\end{aligned}$$

We therefore have the light –cone identifications

$$x^+ \sim x^+ , \quad x^- \sim x^- - 2\pi(\sqrt{2}R).$$

The compactification does not involve light–cone time.

b. With the Lorentz transformation $x' = \gamma(x - \beta ct)$ and $ct' = \gamma(ct - \beta x)$, we have

$$\begin{aligned} x' &\sim \gamma(x + 2\pi R - \beta[ct - 2\pi R]) = x' + 2\pi R\gamma(1 + \beta) \\ ct' &\sim \gamma(ct - 2\pi R - \beta[x + 2\pi R]) = ct' - 2\pi R\gamma(1 + \beta) \end{aligned}$$

$$\begin{pmatrix} x' \\ ct' \end{pmatrix} \sim \begin{pmatrix} x' \\ ct' \end{pmatrix} + 2\pi \sqrt{\frac{1+\beta}{1-\beta}} \begin{pmatrix} R \\ -R \end{pmatrix}.$$

Comparing with (2) we see that in the boosted frame the effective radius of compactification is increased.

Considering the compactification

$$\begin{pmatrix} x \\ ct \end{pmatrix} \sim \begin{pmatrix} x \\ ct \end{pmatrix} + 2\pi \begin{pmatrix} \sqrt{R^2 + R_s^2} \\ -R \end{pmatrix}. \quad (5)$$

c. We search for a frame (x', ct')

$$\begin{aligned} x' &= \gamma(x - \beta ct) \\ ct' &= \gamma(ct - \beta x) \end{aligned}$$

in which this compactification is standard. For this we need $ct' \sim ct'$. So,

$$\begin{aligned} ct' &\sim \gamma(ct - 2\pi R - \beta[x + 2\pi\sqrt{R^2 + R_s^2}]) \\ &= ct' - \gamma(2\pi)(R + \beta\sqrt{R^2 + R_s^2}) \end{aligned}$$

Thus, we need $R + \beta\sqrt{R^2 + R_s^2} = 0$, so

$$\beta = -\frac{R}{\sqrt{R^2 + R_s^2}}.$$

Now examine the x' identification to find the radius:

$$\begin{aligned}
x' &\sim \gamma(x + 2\pi\sqrt{R^2 + R_s^2} - \beta(ct - 2\pi R)) \\
x' &\sim x' + 2\pi\gamma(\sqrt{R^2 + R_s^2} + \beta R)
\end{aligned}$$

Similarly,

$$\begin{aligned}
\gamma(\sqrt{R^2 + R_s^2} + \beta R) &= \frac{1}{\sqrt{1 - \beta^2}}(\sqrt{R^2 + R_s^2} - \frac{R^2}{\sqrt{R^2 + R_s^2}}) \\
&= \frac{\sqrt{R^2 + R_s^2}}{R_s} \left(\frac{R^2 + R_s^2 - R^2}{\sqrt{R^2 + R_s^2}} \right) = R_s
\end{aligned}$$

Thus, $x'x' + 2\pi R_s$. The radius of compactification is R_s . It is interesting to note that $\gamma = \sqrt{R^2 + R_s^2}/R_s$. In the limit as R_s is small this is a very large Lorentz factor $\gamma \sim R/R_s$.