

MATH425 - Quantum Field Theory
Set Work: Sheet 8

1. The Dirac field may be written

$$\psi(x) = \frac{1}{(2\pi)^3} \sum_{s=1}^2 \int \frac{d^3\mathbf{p}}{2p^0} [\psi_p^{(s)}(x) a_s(\mathbf{p}) + \tilde{\psi}_p^{(s)}(x) b_s^\dagger(\mathbf{p})],$$

where

$$\begin{aligned}\psi_p^{(s)}(x) &= e^{-ip \cdot x} u_s(p), \\ \tilde{\psi}_p^{(s)}(x) &= e^{ip \cdot x} v_s(p).\end{aligned}$$

With the scalar product $\langle \psi_1, \psi_2 \rangle = \int \psi_1^\dagger \psi_2 d^3\mathbf{x}$, we have already shown

$$\langle \psi_p^{(r)}, \psi_{p'}^{(s)} \rangle = \delta_{rs} 2p^0 (2\pi)^3 \delta(\mathbf{p} - \mathbf{p}').$$

Show also using

$$\begin{aligned}u_r(p) &= \sqrt{p^0 + m} \begin{pmatrix} \chi_r \\ \frac{\boldsymbol{\sigma} \cdot \mathbf{p}}{p^0 + m} \chi_r \end{pmatrix} \\ v_r(p) &= \sqrt{p^0 + m} \begin{pmatrix} \frac{\boldsymbol{\sigma} \cdot \mathbf{p}}{p^0 + m} \chi_r \\ \chi_r \end{pmatrix}\end{aligned}$$

that

$$\langle \psi_p^{(r)}, \tilde{\psi}_{p'}^{(s)} \rangle = 0.$$

Use these results together with the basic anticommutation relations

$$\{\psi_\alpha(\mathbf{x}, t), \psi_\beta^\dagger(\mathbf{x}', t)\} = \delta(\mathbf{x} - \mathbf{x}') \delta_{\alpha\beta}$$

to show

$$\{b_r(\mathbf{p}), b_s^\dagger(\mathbf{p}')\} = (2\pi)^3 \delta_{rs} 2p^0 \delta(\mathbf{p} - \mathbf{p}').$$

(b) Defining the energy-momentum vector by

$$P^\mu = i \bar{\psi} \gamma^\mu \partial^0 \psi,$$

show that P^μ may be rewritten as

$$P^\mu = \sum_{s=1}^2 \int \frac{d^3\mathbf{p}}{2p^0} p^\mu [a_s^\dagger(\mathbf{p}) a_s(\mathbf{p}) - b_s(\mathbf{p}) b_s^\dagger(\mathbf{p})].$$