

1. See copy of separate handwritten sheet.

2.

$$UU^\dagger = \mathbb{1} \Rightarrow U = e^{iH}; \quad U^\dagger = e^{-iH^\dagger} = e^{-iH},$$

hence we must have $H = H^\dagger$ hermitian.

3. Unitary equivalence and the ‘ $SU(2)$ miracle’:

(a) *First step:*

$$i\sigma_2 = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}, \quad \text{and} \quad \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix},$$

so $i\sigma_2$ is unitary.

Second step: Prove by explicit matrix multiplication that $\sigma_2 \sigma_a^* \sigma_2 = -\sigma_a$ ($i = 1, 2, 3$). For $i = 1$ we have e.g.

$$\sigma_2 \sigma_1^* \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \begin{pmatrix} 0 & -1 \\ -1 & 0 \end{pmatrix} = -\sigma_1.$$

Third step: Choose $W = i\sigma_2$, and with $\sigma_2^* \sigma_2 = \mathbb{1} = \sigma_2 \sigma_2$ and the result of step two we write:

$$\begin{aligned} W^\dagger U^* W &= -i\sigma_2^\dagger U^* i\sigma_2 = \sigma_2 \exp\left(-\frac{i}{2}\theta_a \sigma_a^*\right) \sigma_2 \\ &= \sigma_2 \left(\mathbb{1} - \frac{i}{2}\theta_a \sigma_a^* - \frac{1}{4} \frac{1}{2!} \theta_a \theta_a - \frac{i}{8} \frac{1}{3!} \theta_a \theta_a \theta_b \sigma_b^* - \dots \right) \sigma_2 \\ &= \left(\mathbb{1} + \frac{i}{2}\theta_a \sigma_a - \frac{1}{4} \frac{1}{2!} \theta_a \theta_a + \frac{i}{8} \frac{1}{3!} \theta_a \theta_a \theta_b \sigma_b - \dots \right) = U. \end{aligned}$$

Taking the complex conjugate of

$$W^\dagger U^* W = U$$

we now also have

$$W^{\dagger*} U W^* = U^*$$

and as $W = i\sigma_2 = W^*$ we arrive at the desired relation

$$U^* = W^\dagger U W.$$

(b) *Will be discussed later.*

In the Standard Model, fermion and gauge boson masses are obtained in a gauge invariant way through electroweak symmetry breaking which is mediated by a Higgs potential. Because of the unitary equivalence between the fundamental and the complex conjugate representations of $SU(2)$, gauge invariant mass terms for both up- and down quarks (which are grouped together in $SU(2)$ doublets) can be constructed from only one complex Higgs doublet. In other words, it is due to this special property of $SU(2)$ that the Higgs sector in the Standard Model is the minimal one resulting in only one physical Higgs boson.