

Spinor–Vector Duality and Sterile Neutrinos in String Derived Models

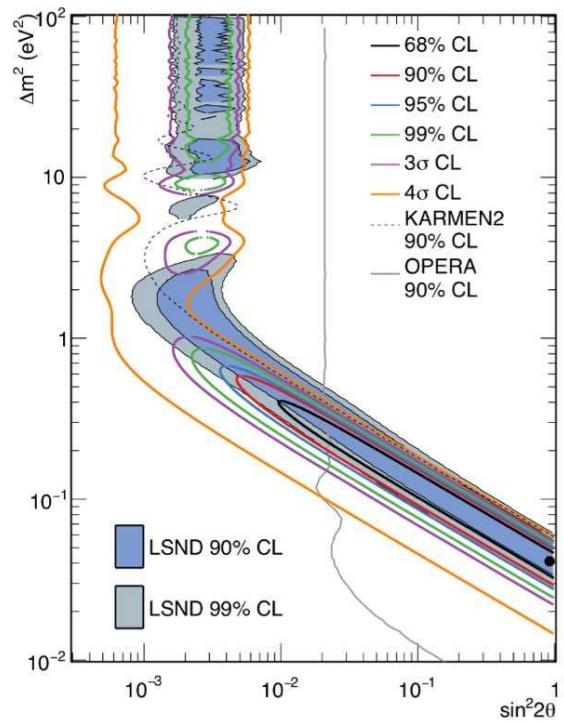


- Spinor–vector duality:

With: Costas Kounnas, John Rizos, Carlo Angelantonj, Mirian Tsulaia,
Ioannis Florakis, Thomas Mohaupt, Panos Athanasopoulos,
Doron Gepner

- Sterile neutrinos:

AEF, European Phys Journal C78 (2018) 867.



LSND 1993: Evidence for sterile neutrinos

MiniBooNE 2018: Further evidence for sterile neutrinos

Sterile neutrinos in string vacua?

String Phenomenology: An answer in search of a question

The Answer : The Standard Model and its BSM extensions

A question : What is the true string vacuum?

A question : Can we identify signatures of classes of string

compactifications in the experimental particle data?

In this talk:

Light sterile neutrinos in heterotic string vacua require a light Z'

A light Z' is hard to implement in heterotic string constructions

PHENOMENA

DATA → STANDARD MODEL

EWX → PERTUBATIVE

STANDARD MODEL → UNIFICATION

EVIDENCE: 16 of SO(10), Log running, proton stability, neutrino masses

+ GRAVITY < -- > STRINGS

PRIMARY GUIDES:

3 generations
SO(10) embedding

REALISTIC STRING MODELS :

heterotic 10D \rightarrow heterotic 4D

6D compactifications $(T^2 \times T^2 \times T^2)$

Orbifold – twists of flat 6D torus



FREE FERMIONIC MODELS –

$Z_2 \times Z_2$ Orbifold $\rightarrow U(1)_Y \in SO(10)$

$$\frac{6}{2} = 1+1+1$$

Realistic free fermionic models

'Phenomenology of the Standard Model and Unification'

- Minimal Superstring Standard Model NPB 335 (1990) 347
(with Nanopoulos & Yuan)
- Top quark mass $\sim 175\text{--}180\text{GeV}$ PLB 274 (1992) 47
- Generation mass hierarchy NPB 407 (1993) 57
- CKM mixing NPB 416 (1994) 63 (with Halyo)
- Stringy seesaw mechanism PLB 307 (1993) 311 (with Halyo)
- Gauge coupling unification NPB 457 (1995) 409 (with Dienes)
- Proton stability NPB 428 (1994) 111
- Squark degeneracy NPB 526 (1998) 21 (with Pati)
- Moduli fixing NPB 728 (2005) 83
- Classification & Exophobia 2003 – . . .
(with Nooij, Assel, Christodoulides, Kounnas, Rizos & Sonmez)

Other approaches

Geometrical

- Greene, Kirklin, Miron, Ross (1987)
Donagi, Ovrut, Pantev, Waldram (1999)
Blumenhagen, Moster, Reinbacher, Weigand (2006)
Heckman, Vafa (2008)
-

Orbifolds

- Ibanez, Nilles, Quevedo (1987)
Bailin, Love, Thomas (1987)
Kobayashi, Raby, Zhang (2004)
Lebedev, Nilles, Raby, Ramos-Sanchez, Ratz, Vaudrevange, Wingerter (2007)
Blaszczyk, Groot–Nibbelink, Ruehle, Trapletti, Vaudrevange (2010)
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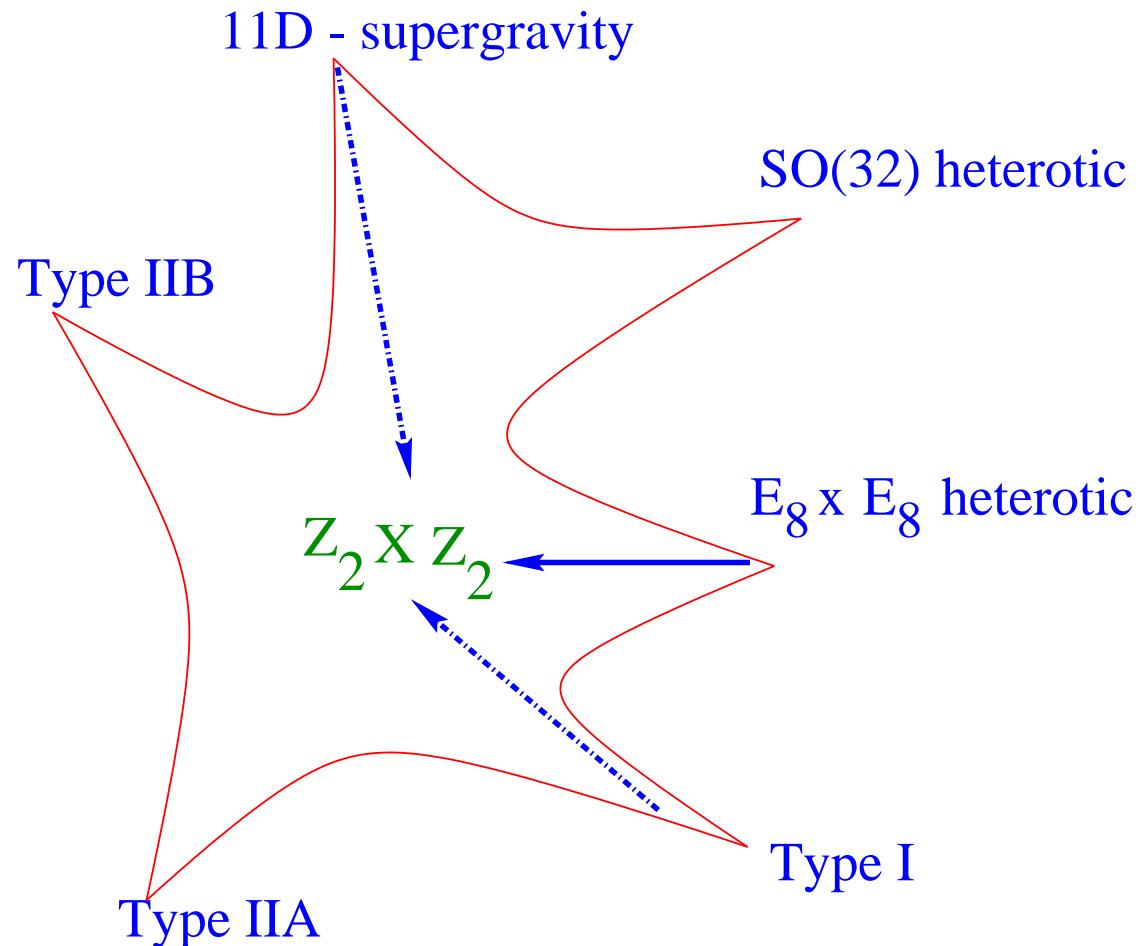
Other CFTs

- Gepner (1987)
Schellekens, Yankielowicz (1989)
Gato–Rivera, Schellekens (2009)
-

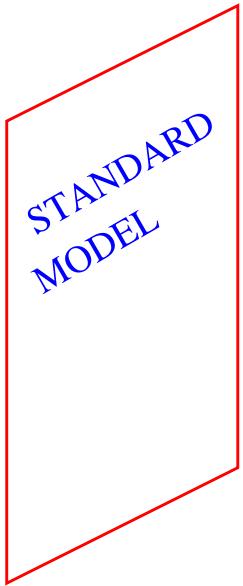
Orientifolds

- Cvetic, Shiu, Uranga (2001)
Ibanez, Marchesano, Rabadan (2001)
Kiritsis, Schellekens, Tsulaia (2008)
-

Point, String, Membrane



Sterile neutrinos in large volume scenarios



RHN
●

$$\nu_R^{(c)}(x, y) = \sum_n \frac{1}{\sqrt{2\pi r}} \nu_{Rn}^{(c)}(x) e^{iny/r}$$

$$S^{\text{int}} = \int d^4x \lambda l(x) h^*(x) \nu_R(x, y=0)$$

$$\lambda_{(4)} = \frac{\lambda}{\sqrt{r^n M_*^n}}, \quad \text{plus tower of KK modes} \longrightarrow \text{sterile neutrinos}$$

String GUT models in the Free Fermionic Construction:

Left-Movers: $\psi_{1,2}^\mu, \chi_i, y_i, \omega_i$ ($i = 1, \dots, 6$)

Right-Movers

$$\bar{\phi}_{A=1,\dots,44} = \begin{cases} \bar{y}_i, \bar{\omega}_i & i = 1, \dots, 6 \\ \bar{\eta}_i & U(1)_i \quad i = 1, 2, 3 \\ \bar{\psi}_{1,\dots,5} & SO(10) \\ \bar{\phi}_{1,\dots,8} & SO(16) \end{cases}$$

Models \longleftrightarrow Basis vectors + one-loop phases

Free Fermionic Models:

$Z_2 \times Z_2$ orbifolds
with discrete Wilson lines

Sterile neutrinos in free Fermionic models

Analyse mass terms contributing to the neutrino mass matrix

AEF & Edi Halyo, PLB 307 (1993) 311

AEF & Claudio Coriano, PLB 581 (2004) 99

The generic form:

$$\begin{pmatrix} \nu_i, N_i, \phi_i \end{pmatrix} \begin{pmatrix} 0 & (M_D)_{ij} & 0 \\ (M_D)_{ij} & 0 & \langle \bar{N} \rangle_{ij} \\ 0 & \langle \bar{N} \rangle_{ij} & \langle \phi \rangle_{ij} \end{pmatrix} \begin{pmatrix} \nu_j \\ N_j \\ \phi_j \end{pmatrix},$$

$$m_{\nu_j} \sim \left(\frac{k M_u^j}{\langle \bar{N} \rangle} \right)^2 \langle \phi \rangle \quad , \quad m_{N_j}, m_\phi \sim \langle \bar{N} \rangle .$$
$$\langle \bar{N} \rangle > 10^{14} \text{GeV}$$

No Light sterile neutrino

Specific example:

	L_3	L_2	L_1	N_3	N_2	N_1	H_{23}	H_{25}	Φ_{13}	Φ_{45}	Φ_2^+	$\bar{\Phi}_2^+$
L_3	0	0	0	0	0	r	0	0	0	0	0	0
L_2	0	0	0	r	0	r	0	0	0	0	0	0
L_1	0	0	0	r	0	v	0	0	0	0	0	0
N_3	0	r	r	0	0	0	0	x	0	0	0	0
N_2	0	0	0	0	0	0	z	0	u	u	z	u
N_1	r	r	v	0	0	0	0	w	0	0	0	0
H_{23}	0	0	0	0	z	0	p	0	x	x	x	x
H_{25}	0	0	0	x	0	w	0	0	0	0	0	0
Φ_{13}	0	0	0	0	u	0	x	0	q	y	0	y
Φ_{45}	0	0	0	0	z	0	p	0	x	x	x	x
Φ_2^+	0	0	0	0	z	0	x	0	0	q	0	x
$\bar{\Phi}_2^+$	0	0	0	0	0	0	x	0	y	q	x	q

mass eigenvalues $\{1.7 \times 10^{13}, \dots, 2.4^{-8}\}$ GeV

No Light sterile neutrinos \rightarrow non-chiral matter \rightarrow heavy

Sterile neutrinos \rightarrow Chiral under and extra $U(1)$ symmetry

Low scale Z' in free fermionic models:

- $\frac{3}{2}U(1)_{B-L} - 2U(1)_R \in SO(10) @ 1TeV$ MPL A6 (1991) 61
(with Nanopoulos)
- But $m_t = m_{\nu_\tau}$ & $1TeV Z' \Rightarrow m_{\nu_\tau} \approx 10MeV$ PLB 245 (1990) 435
 $E_6 \rightarrow SO(10) \times U(1)_A \implies U(1)_A$ is anomalous!
 $\implies U(1)_A \notin$ low scale $U(1)_{Z'}$
- 1996-2013, Pati, AEF, Guzzi, Mehta, Athanasopoulos, $U(1) \notin E_6$
- On the other hand(AEF, Viraf Mehta, PRD88 (2013) 025006)
 $\sin^2 \theta_W(M_Z), \alpha_s(M_Z) \implies U(1)_{Z'} \in E_6$
- Z' string derived model, (with Rizos) NPB 895 (2015) 233

Classification of fermionic $Z_2 \times Z_2$ orbifolds

(PS talk by Ben Percival)

Basis vectors:

$$1 = \{\psi^\mu, \chi^{1,\dots,6}, y^{1,\dots,6}, \omega^{1,\dots,6} \mid \bar{y}^{1,\dots,6}, \bar{\omega}^{1,\dots,6}, \bar{\eta}^{1,2,3}, \bar{\psi}^{1,\dots,5}, \bar{\phi}^{1,\dots,8}\}$$

$$S = \{\psi^\mu, \chi^{1,\dots,6}\},$$

$$z_1 = \{\bar{\phi}^{1,\dots,4}\},$$

$$z_2 = \{\bar{\phi}^{5,\dots,8}\},$$

$$e_i = \{y^i, \omega^i | \bar{y}^i, \bar{\omega}^i\}, \quad i = 1, \dots, 6, \quad N = 4 \text{ Vacua}$$

$$b_1 = \{\chi^{34}, \chi^{56}, y^{34}, y^{56} | \bar{y}^{34}, \bar{y}^{56}, \bar{\eta}^1, \bar{\psi}^{1,\dots,5}\}, \quad N = 4 \rightarrow N = 2$$

$$b_2 = \{\chi^{12}, \chi^{56}, y^{12}, y^{56} | \bar{y}^{12}, \bar{y}^{56}, \bar{\eta}^2, \bar{\psi}^{1,\dots,5}\}, \quad N = 2 \rightarrow N = 1$$

$$\alpha = \{\bar{\psi}^{4,5}, \bar{\phi}^{1,2}\} \quad \& \quad SO(10) \rightarrow SO(6) \times SO(4) \times \cdots$$

$$\beta = \{\bar{\psi}^{1,\dots,5} \equiv \frac{1}{2}, \dots\} \quad \& \quad SO(10) \rightarrow SU(5) \times U(1) \times \cdots$$

Independent phases $c_{v_j}^{[vi]} = \exp[i\pi(v_i|v_j)]$: upper block

	1	S	e_1	e_2	e_3	e_4	e_5	e_6	z_1	z_2	b_1	b_2	α
1	-1	-1	\pm										
S			-1	-1	-1	-1	-1	-1	-1	-1	1	1	-1
e_1				\pm									
e_2					\pm								
e_3						\pm							
e_4							\pm						
e_5								\pm	\pm	\pm	\pm	\pm	\pm
e_6									\pm	\pm	\pm	\pm	\pm
z_1										\pm	\pm	\pm	\pm
z_2										\pm	\pm	\pm	\pm
b_1											\pm	\pm	\pm
b_2												-1	\pm
α													

A priori 66 independent coefficients $\rightarrow 2^{66}$ distinct vacua

The twisted matter spectrum:

$$B_{\ell_3^1 \ell_4^1 \ell_5^1 \ell_6^1}^1 = S + b_1 + \ell_3^1 e_3 + \ell_4^1 e_4 + \ell_5^1 e_5 + \ell_6^1 e_6$$

$$B_{\ell_1^2 \ell_2^2 \ell_5^2 \ell_6^2}^2 = S + b_2 + \ell_1^2 e_1 + \ell_2^2 e_2 + \ell_5^2 e_5 + \ell_6^2 e_6$$

$$B_{\ell_1^3 \ell_2^3 \ell_3^3 \ell_4^3}^3 = S + b_3 + \ell_1^3 e_1 + \ell_2^3 e_2 + \ell_3^3 e_3 + \ell_4^3 e_4 \quad l_i^j = 0, 1$$

sectors $B_{pqrs}^i \rightarrow 16$ or $\overline{16}$ of $SO(10)$ with multiplicity $(1, 0, -1)$

$B_{pqrs}^i + x \rightarrow 10$ of $SO(10)$ with multiplicity $(1, 0)$

$x = \{\bar{\psi}^{1,\dots,5}, \bar{\eta}^1, \bar{\eta}^2, \bar{\eta}^3\}$ x – map \leftrightarrow spinor–vector map

Algebraic formulas for $S = \sum_{i=1}^3 S_+^{(i)} - S_-^{(i)}$ and $V = \sum_{i=1}^3 V^{(i)}$

Spinor–vector duality:

Duality under exchange of spinors and vectors.

First Plane			Second plane			Third Plane			
s	\bar{s}	v	s	\bar{s}	v	s	\bar{s}	v	# of models
2	0	0	0	0	0	0	0	0	1325963712
0	2	0	0	0	0	0	0	0	1340075584
1	1	0	0	0	0	0	0	0	3718991872
0	0	2	0	0	0	0	0	0	6385031168

of models with $\#(16 + \overline{16}) = \#$ of models with $\#(10)$

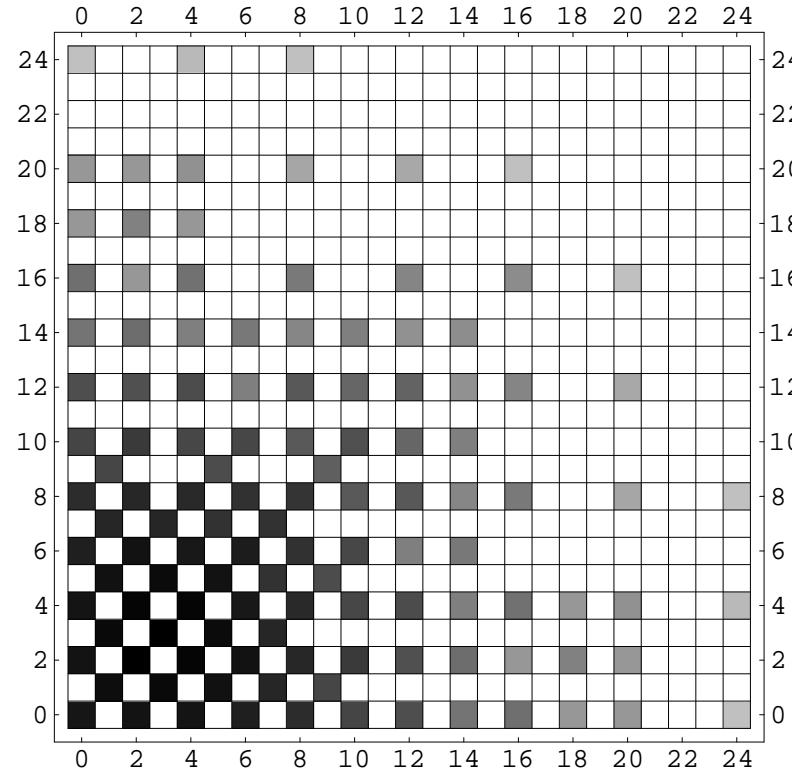
For every model with $\#(16 + \overline{16}) \& \#(10)$

There exist another model in which they are interchanged

Reflects discrete exchange of phases

Spinor–vector duality:

Invariance under exchange of $\#(16 + \overline{16}) < - > \#(10)$



Symmetric under exchange of rows and columns

$$E_6 : \quad 27 = 16 + 10 + 1 \quad \overline{27} = \overline{16} + 10 + 1$$

Self-dual: $\#(16 + \overline{16}) = \#(10)$ without E_6 symmetry

light Z' heterotic–string model $c[v_i]_{v_j} = \exp[i\pi(v_i|v_j)]$:

$$(v_i|v_j) = \begin{pmatrix} 1 & S & e_1 & e_2 & e_3 & e_4 & e_5 & e_6 & b_1 & b_2 & z_1 & z_2 & \alpha \\ 1 & \begin{pmatrix} 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 \end{pmatrix} \\ S & \begin{pmatrix} 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 \end{pmatrix} \\ e_1 & \begin{pmatrix} 1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \end{pmatrix} \\ e_2 & \begin{pmatrix} 1 & 1 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 1 & 0 \end{pmatrix} \\ e_3 & \begin{pmatrix} 1 & 1 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 1 & 1 \end{pmatrix} \\ e_4 & \begin{pmatrix} 1 & 1 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \end{pmatrix} \\ e_5 & \begin{pmatrix} 1 & 1 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 1 & 1 \end{pmatrix} \\ e_6 & \begin{pmatrix} 1 & 1 & 0 & 1 & 0 & 0 & 1 & 0 & 0 & 0 & 1 & 0 & 0 \end{pmatrix} \\ b_1 & \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 1 & 0 & 0 & 0 \end{pmatrix} \\ b_2 & \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 1 & 0 & 0 & 1 \end{pmatrix} \\ z_1 & \begin{pmatrix} 1 & 1 & 0 & 0 & 0 & 1 & 0 & 1 & 0 & 0 & 1 & 1 & 0 \end{pmatrix} \\ z_2 & \begin{pmatrix} 1 & 1 & 0 & 1 & 1 & 0 & 1 & 0 & 0 & 0 & 1 & 1 & 0 \end{pmatrix} \\ \alpha & \begin{pmatrix} 1 & 1 & 1 & 0 & 1 & 0 & 1 & 0 & 1 & 0 & 1 & 0 & 1 \end{pmatrix} \end{pmatrix}$$

Observable gauge group: $SO(6) \times SO(4) \times U(1)_{1,2,3}$

$U(1)_\zeta = U(1)_1 + U(1)_2 + U(1)_3$ is anomaly free

sector	field	$SU(4) \times SU(2)_L \times SU(2)_R$	$U(1)_1$	$U(1)_2$	$U(1)_3$	$U(1)_\zeta$
$S + b_1$	\bar{F}_{1R}	($\bar{\mathbf{4}}, \mathbf{1}, \mathbf{2}$)	1/2	0	0	1/2
$S + b_1 + e_3 + e_5$	F_{1R}	($\mathbf{4}, \mathbf{1}, \mathbf{2}$)	1/2	0	0	1/2
$S + b_2$	F_{1L}	($\mathbf{4}, \mathbf{2}, \mathbf{1}$)	0	1/2	0	1/2
$S + b_2 + e_1 + e_2 + e_5$	F_{2L}	($\mathbf{4}, \mathbf{2}, \mathbf{1}$)	0	1/2	0	1/2
$S + b_2 + e_1$	\bar{F}_{2R}	($\bar{\mathbf{4}}, \mathbf{1}, \mathbf{2}$)	0	1/2	0	1/2
$S + b_2 + e_2 + e_5$	\bar{F}_{3R}	($\bar{\mathbf{4}}, \mathbf{1}, \mathbf{2}$)	0	1/2	0	1/2
$S + b_3 + e_1 + e_2$	F_{3L}	($\mathbf{4}, \mathbf{2}, \mathbf{1}$)	0	0	1/2	1/2
$S + b_3 + e_2$	\bar{F}_{4R}	($\bar{\mathbf{4}}, \mathbf{1}, \mathbf{2}$)	0	0	1/2	1/2
$S + b_3 + x$	h_1	($\mathbf{1}, \mathbf{2}, \mathbf{2}$)	-1/2	-1/2	0	-1
$S + b_2 + x + e_5$	h_2	($\mathbf{1}, \mathbf{2}, \mathbf{2}$)	-1/2	0	-1/2	-1
$S + b_2 + x + e_1 + e_2$	h_3	($\mathbf{1}, \mathbf{2}, \mathbf{2}$)	-1/2	0	-1/2	-1
$S + b_3 + x + e_1$	D_4	($\mathbf{6}, \mathbf{1}, \mathbf{1}$)	-1/2	-1/2	0	-1
	χ_1^+	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	1/2	1/2	1	+2
	χ_1^-	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	1/2	1/2	-1	0
	$\zeta_a, a = 2, 3$	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	1/2	-1/2	0	0
	$\bar{\zeta}_a, a = 2, 3$	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	-1/2	1/2	0	0
$S + b_2 + x + e_1 + e_5$	D_5	($\mathbf{6}, \mathbf{1}, \mathbf{1}$)	-1/2	0	-1/2	-1
	χ_2^+	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	1/2	1	1/2	+2
	χ_2^-	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	1/2	-1	1/2	0
	$\zeta_a, a = 4, 5$	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	1/2	0	-1/2	0
	$\bar{\zeta}_a, a = 4, 5$	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	-1/2	0	1/2	0
$S + b_2 + x + e_2$	D_6	($\mathbf{6}, \mathbf{1}, \mathbf{1}$)	-1/2	0	-1/2	-1
	χ_3^+	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	1/2	1	1/2	+2
	χ_3^-	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	1/2	-1	1/2	0
	$\zeta_a, a = 6, 7$	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	1/2	0	-1/2	0
	$\bar{\zeta}_a, a = 6, 7$	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	-1/2	0	1/2	0
$S + b_1 + x + e_3$	\bar{D}_6	($\bar{\mathbf{6}}, \mathbf{1}, \mathbf{1}$)	0	1/2	1/2	+1
	$\bar{\chi}_4^+$	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	-1	-1/2	-1/2	-2
	$\bar{\chi}_4^-$	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	1	-1/2	-1/2	0
	$\zeta_a, a = 8, 9$	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	0	1/2	-1/2	0
	$\bar{\zeta}_a, a = 8, 9$	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	0	-1/2	1/2	0
$S + b_1 + x + e_5$	D_7	($\mathbf{6}, \mathbf{1}, \mathbf{1}$)	0	-1/2	-1/2	-1
	χ_5^+	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	1	1/2	1/2	+2
	χ_5^-	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	-1	1/2	1/2	0
	$\zeta_a, a = 10, 11$	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	0	1/2	-1/2	0
	$\bar{\zeta}_a, a = 10, 11$	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	0	-1/2	1/2	0
$S + b_3 + x + e_2 + e_3$	ζ_1	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	1/2	-1/2	0	0
	$\bar{\zeta}_1$	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	-1/2	1/2	0	0
$S + b_1 + x + e_3 + e_4 + e_6$	ϕ_1	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	0	1/2	1/2	+1
	$\bar{\phi}_1$	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	0	-1/2	-1/2	-1
$S + b_1 + x + e_4 + e_5 + e_6$	ϕ_2	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	0	1/2	1/2	+1
	$\bar{\phi}_2$	($\mathbf{1}, \mathbf{1}, \mathbf{1}$)	0	-1/2	-1/2	-1

Table 1: Observable twisted matter spectrum and $SU(4) \times SU(2)_L \times SU(2)_R \times U(1)^3$ quantum numbers.

The chiral spectrum is self–dual under spinor–vector duality

maintains the E_6 embedding $\Rightarrow U(1)_\zeta$ is anomaly free

Exophobic, three generations + heavy and light higgs + $\lambda_t \sim 1$

$$U(1)_{Z'} = \frac{1}{5} U(1)_C - \frac{1}{5} U(1)_L - U_\zeta$$

$Z' \Rightarrow$ Additional matter at the Z' breaking scale

Exotic $SO(10)$ singlets with non–standard $U(1)_\zeta$ charges

\Rightarrow Natural Wilsonian dark–matter candidates

Z' model at low scales Heavy Higgs $\langle \mathcal{N} \rangle \sim M_{\text{String}}$ \rightarrow high seesaw \rightarrow

Field	$SU(3)_C$	$\times SU(2)_L$	$U(1)_Y$	$U(1)_{Z'}$
Q_L^i	3	2	$+\frac{1}{6}$	$-\frac{2}{5}$
u_L^i	$\bar{3}$	1	$-\frac{2}{3}$	$-\frac{2}{5}$
d_L^i	$\bar{3}$	1	$+\frac{1}{3}$	$-\frac{4}{5}$
e_L^i	1	1	+1	$-\frac{2}{5}$
L_L^i	1	2	$-\frac{1}{2}$	$-\frac{4}{5}$
D^i	3	1	$-\frac{1}{3}$	$+\frac{4}{5}$
\bar{D}^i	$\bar{3}$	1	$+\frac{1}{3}$	$+\frac{6}{5}$
H^i	1	2	$-\frac{1}{2}$	$+\frac{6}{5}$
\bar{H}^i	1	2	$+\frac{1}{2}$	$+\frac{4}{5}$
S^i	1	1	0	-2
h	1	2	$-\frac{1}{2}$	$-\frac{4}{5}$
\bar{h}	1	2	$+\frac{1}{2}$	$+\frac{4}{5}$
ϕ	1	1	0	-1
$\bar{\phi}$	1	1	0	+1
ζ^i	1	1	0	0

Additional matter states at $U(1)_{Z'}$ breaking scale

Neutrino mass spectrum

a plausible scenario

$$\begin{pmatrix} L_i & S_i & H_i & \bar{H}_i & N_i \\ L_i & 0 & 0 & \lambda n & \lambda v \\ S_i & 0 & 0 & \lambda v_2 & \lambda v_3 \\ H_i & 0 & \lambda v_2 & 0 & z' \\ \bar{H}_i & \lambda n & \lambda v_3 & z' & 0 \\ N_i & \lambda v & 0 & 0 & \mathcal{N}^2/M \end{pmatrix},$$

$$\lambda v = 1 \text{GeV};$$

$$\lambda v_2 = 5 \times 10^{-4} \text{GeV} \approx m_e;$$

$$z' = 5 \times 10^4 \text{GeV} = 50 \text{TeV};$$

$$\lambda n = 5 \times 10^{-4} \text{GeV}; \approx m_e;$$

$$\lambda v_3 = 5 \times 10^{-4} \text{GeV}; \approx m_e;$$

$$\mathcal{N} = 5 \times 10^{14} \text{GeV},$$

$$m_{1,2} = \{10^{-2} \text{eV}, 10^{-3} \text{eV}\} \quad \text{mix of } \nu_L^i \text{ and } S_i, \text{ with } \sin \theta \approx 0.98.$$

$$m_{3,4,5} = \{50 \text{TeV}, 50 \text{TeV}, 2.5 \times 10^{11} \text{GeV}\}.$$

$$m_5 \leftrightarrow N_i \quad m_{3,4} \leftrightarrow \text{equal mix of } H_i \text{ and } \bar{H}_i$$

Spinor–Vector Duality from a Novel Basis

$$S = \{\psi^\mu, \chi^{1,\dots,6}\},$$

$$z_1 = \{\bar{\phi}^{1,\dots,4}\},$$

$$z_2 = \{\bar{\phi}^{5,\dots,8}\},$$

$$z_3 = \{\bar{\psi}^{1,\dots,4}\},$$

$$z_4 = \{\bar{\eta}^{0,\dots,3}\},$$

$$e_i = \{y^i, \omega^i | \bar{y}^i, \bar{\omega}^i\}, \quad i = 1, \dots, 6,$$

$$\bar{\eta}^0 \equiv \bar{\psi}^5$$

$N = 4$ Vacua

$$1 = S + \sum e_i + z_1 + z_2 + z_3 + z_4$$

$$b_1 = \{\chi^{34}, \chi^{56}, y^{34}, y^{56} | \bar{y}^{34}, \bar{y}^{56}, \bar{\eta}^1, \bar{\psi}^{1,\dots,5}\}, \quad N = 4 \rightarrow N = 2$$

Vector bosons: NS, $z_{1,2,3,4}$, $z_i + z_j$ NS $\rightarrow SO(8)^4$

$SO(12)$ -GUT \rightarrow from enhancement

Duality picture is facilitated

$$\text{Spinor} \longleftrightarrow \text{Vector map} \longrightarrow B \longleftrightarrow B + z_4$$

$$SO(12) \text{ enhancement} \longrightarrow B \longleftrightarrow B + z_3$$

A convenient basis to study dualities; modular properties

→ 2D → 24 dimensional lattices

GUT structure is obscured

Conclusions

- DATA → UNIFICATION
- STRINGS → GAUGE & GRAVITY UNIFICATION
- Free fermionic models → A Fertile Crescent $\longleftrightarrow Z_2 \times Z_2$ orbifolds
- Sterile neutrinos → chiral matter charged under low scale extra $U(1)$
- Low scale Z' → hard to implement in string GUT constructions
- Spinor–vector duality (alternatively: extended GUT embeddings)
- String Phenomenology → Physics of the third millennium
e.g. Aristarchus to Copernicus