#### Very high order perturbation theory: Lattice Wilson loops

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#### or

#### NNNNNNNNNNNNNNNNNNNNN QCD

#### Introduction

Talk about results from 20 loop perturbation theory. Progress: More loops, More quantities measured.

"Wilson loops in very high order lattice perturbation theory" E.-M. Ilgenfritz, Y. Nakamura, H. Perlt, P.E.L. Rakow, G. Schierholz, A. Schiller PoS(LAT2009)236 arXiv:0910.2795v1 [hep-lat]

Any sign that perturbation theory diverges?

### Introduction

- Introduction
- Divergence of Perturbation theory
- Stochastic Perturbation Theory
- Plaquette
- Wilson Loops
- Perturbation theory meets results

#### Conclusions

In perturbation theory we calculate physical quantities as series in a coupling constant.

Very succesful for QED, ( $\alpha_{EM} = 1/137$ ).

More difficult for QCD, ( $\alpha_{QCD} \sim 1/10$ ), most useful for high energy processes (asymptotic freedom). (One motive for lattice QCD.)

Expectation is that perturbation series diverge. (Nearest singularity is at  $g^2 = 0$ .)

SU(2) in two dimensions.

Gauge theories in two dimensions are simple, can often be solved exactly.

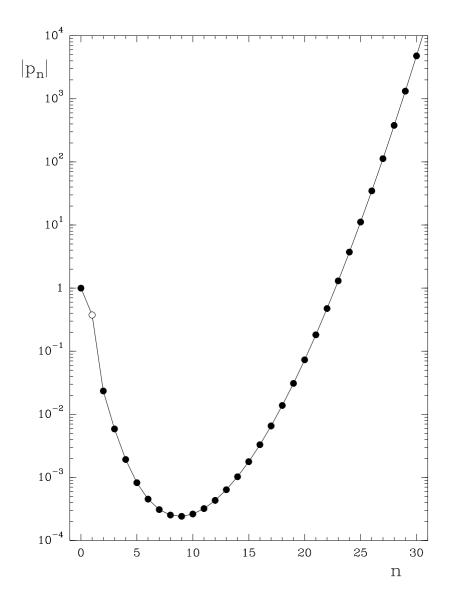
The "plaquette",  $\langle F_{\mu\nu}F^{\mu\nu}\rangle$ , for SU(2) is

 $P = \frac{I_2(4/(ga)^2)}{I_1(4/(ga)^2)}$ 

We know how the Bessel functions  $I_{\nu}$  behave at large argument, so we can expand this expression to get a series for the plaquette when  $(ga)^2$  is small.

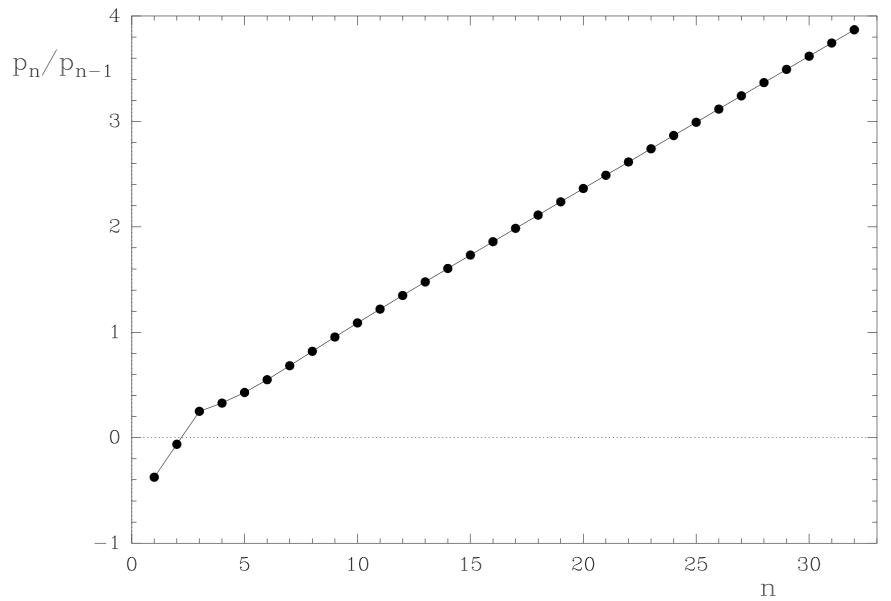
$$P = 1 - \frac{3}{8}(ga)^2 + \frac{3}{128}(ga)^4 + \frac{3}{512}(ga)^6 + \cdots$$
$$\equiv 1 + \sum_n p_n (ga)^{2n}$$

Series looks harmless:



Although the coefficients of the series start by decreasing rapidly, they soon turn around and grow faster than any exponential.

We can see how the coefficients are growing by looking at the ratio of succesive  $p_n$ . (Like the ratio test in Year 1 lectures)



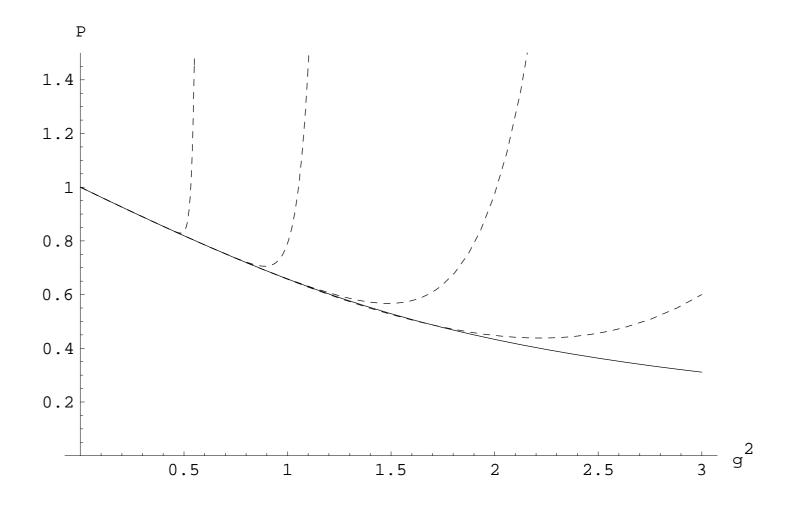
Very high order perturbation theory:Lattice Wilson loops - p.10/48

The ratio grows linearly with n, so the coefficients are growing factorially.

 $p_n \sim n!$ 

"Renormalon behaviour"

How does this effect the usefullness of the series?



Increasing the number of terms decreases the region of validity. (Works the opposite way to convergent series.) More and more ....

The interesting case - SU(3) in four dimensions Does the same sort of thing happen for QCD? How can we begin to look at this question?

Exact solution: Win \$ 1,000,000

### **Conventional Perturbation Theory**

$$P_{pert}(g^2) = 1 + \sum_n p_n g^{2n}$$

Lattice Feynman Rules are more complicated than continuum.

Nevertheless the first 3 coefficients in the series have been found by conventional techniques. (4 d integral, 8 dim integral, 12 dim integral) Alles, Campostrini, Feo, Panagopolous

Wilson gauge action Pure gauge theory

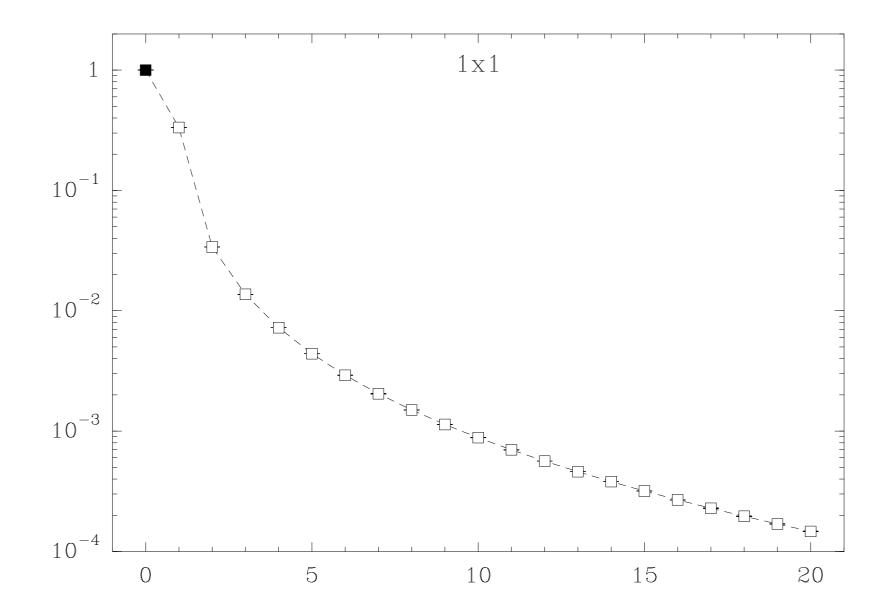
Need more terms, but difficulty grows very fast with increased loop number.

### **Stochastic Perturbation Theory**

A mechanised way to do lattice perturbation theory Burgio, Di Renzo, Onofri, Marchesini

Replace each field of lattice gauge theory by a series in g. Gives a Monte Carlo estimate for each  $p_n$ . Precision not as large as by doing integrals.

Effort grows as a polynomial of loop number. Can go much farther in n.



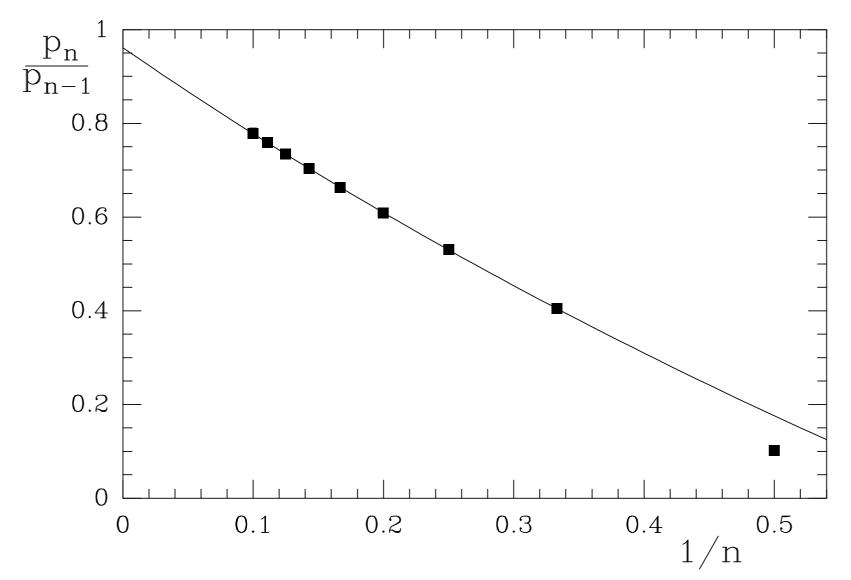
Doesn't look like the 2-d pattern. Terms dropping, curvature decreasing. Large *n* behaviour ? Domb-Sykes plot. Ratio of coefficients vs  $\frac{1}{n}$ .

$$f(g^2) \sim (1 - ug^2)^q$$

$$c_n \sim n^{-(1+q)}u^n$$

$$r_n \sim u \left[1 - \frac{(1+q)}{n}\right]$$

Straight line  $\Leftrightarrow$  Power law singularity. Convergence radius 1/u.



R Horsley, PEL Rakow, G Schierholz, hep-lat/0110210 (2001)

Fit to 10-loop data.

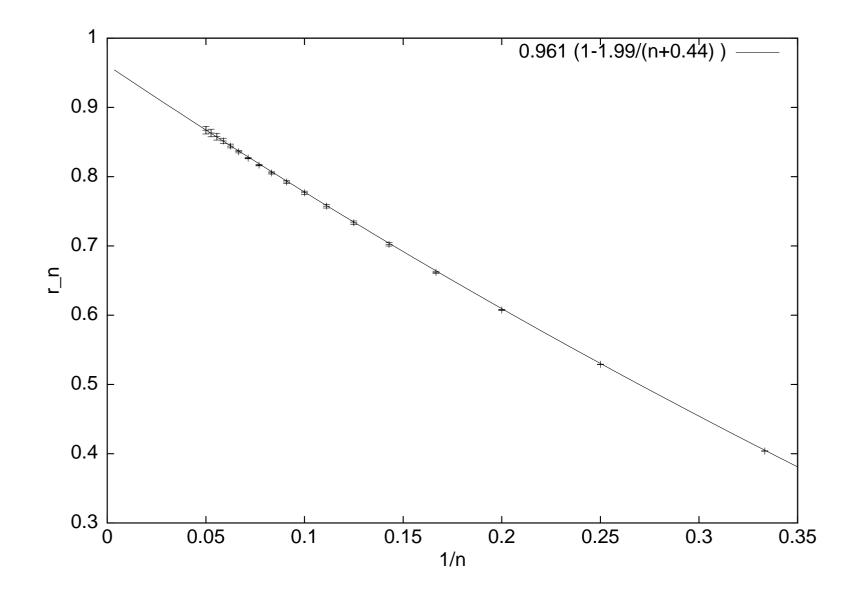
To allow some curvature we made a fit

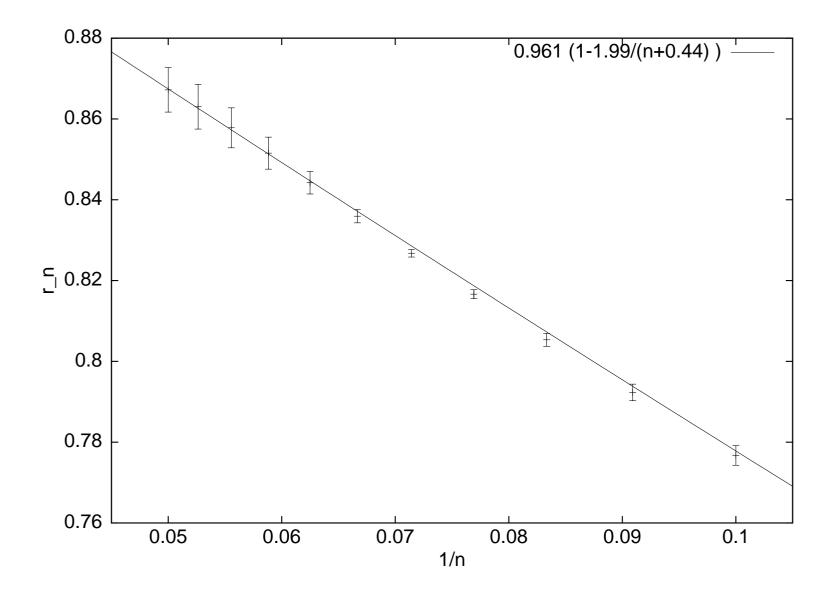
$$r_n = u\left(1 - \frac{(1+q)}{n+s}\right) = u\left(\frac{n+s-1-q}{n+s}\right)$$

and found

 $u = 0.961(9), \quad q = 0.99(7), \quad s = 0.44(10)$ .

Looks like a series which converges for  $g^2 < 1.04$ . How does this look now that we have 20 terms in the series (2009)?

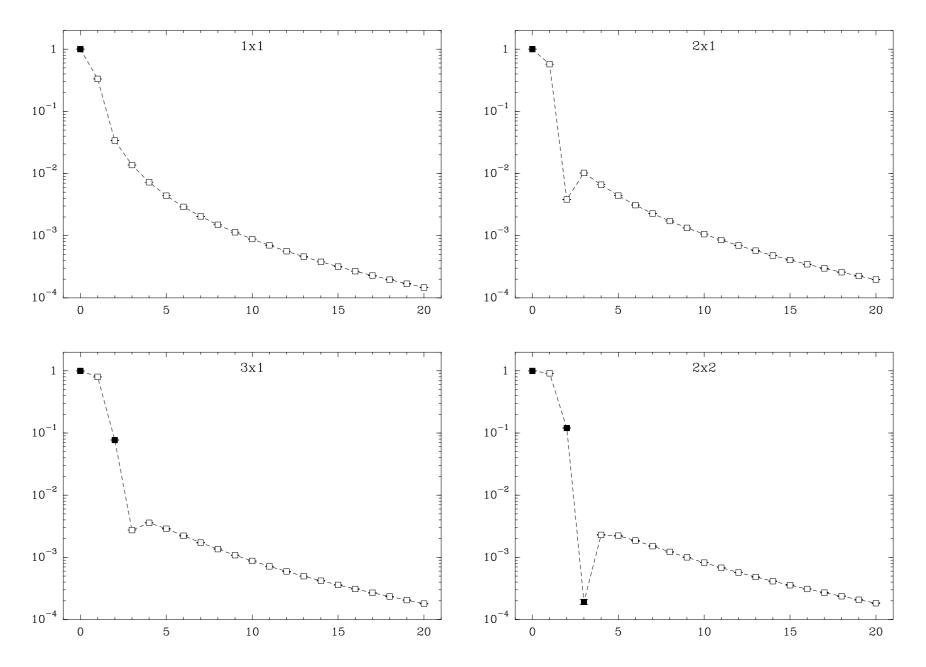


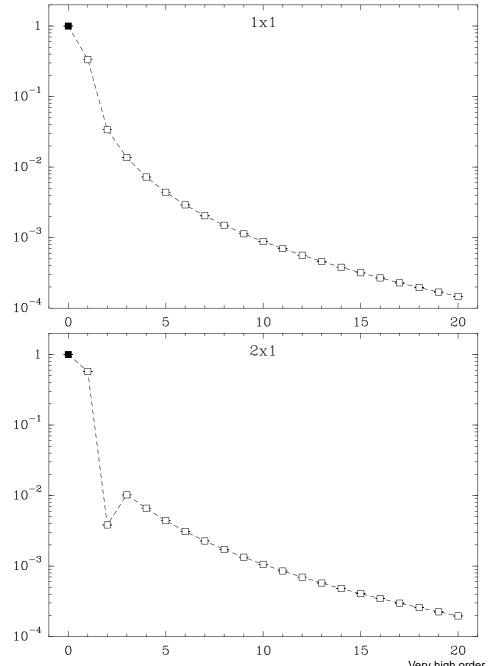


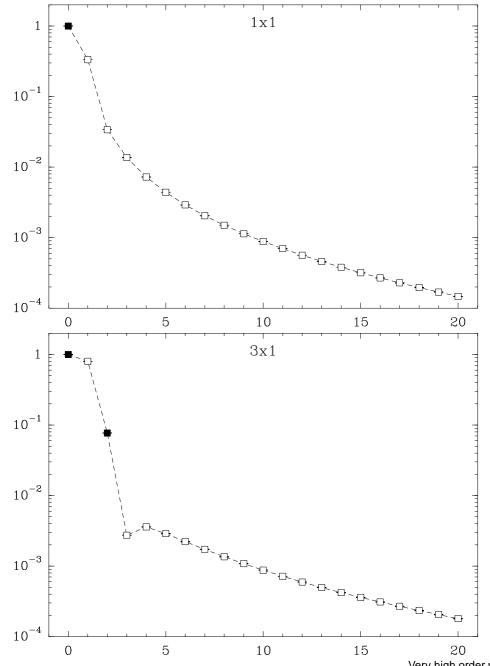
Until recently the plaquette was the only quantity for which we had a really long series.

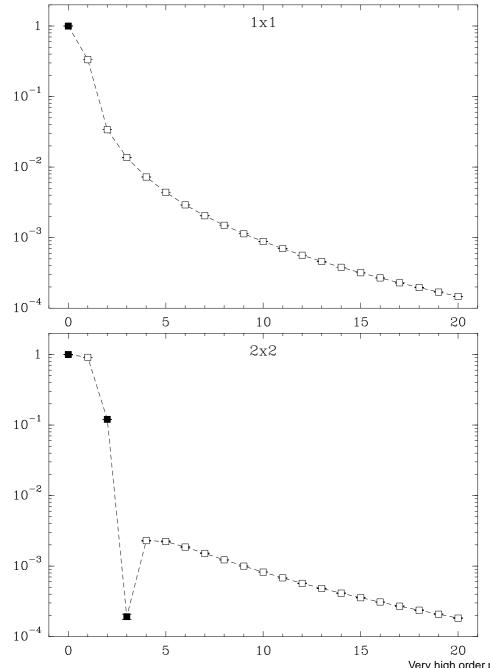
Now we have series for Wilson loops of various sizes.

 $\begin{array}{l} 1\times1, 2\times1, 3\times1, 2\times2\\ 4\times1, 5\times1, 6\times1, 3\times3, 4\times4\\ 1\times1 \text{ is the plaquette}\\ \end{array}$  Was the plaquette a freak? What do the other Wilson loops look like?

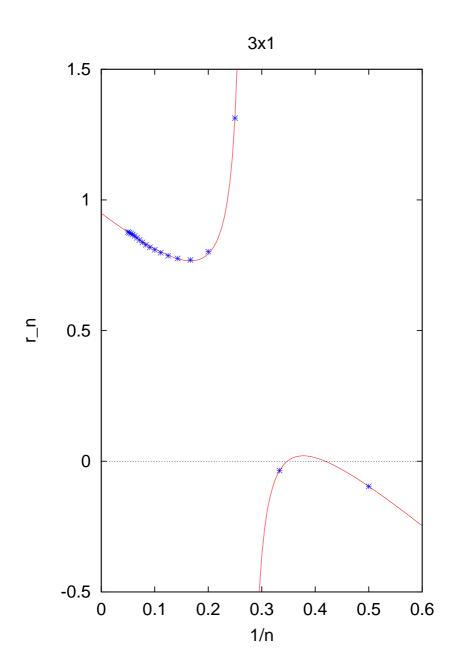


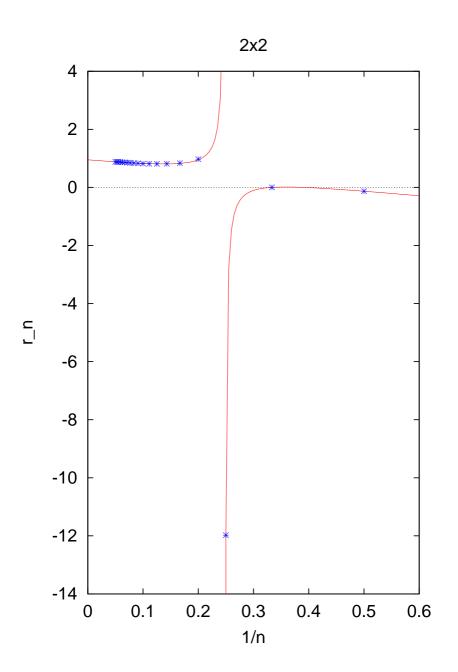






Other Wilson loops show more complex behaviour than plaquette at small n, but settle down to similar behaviour at large n.

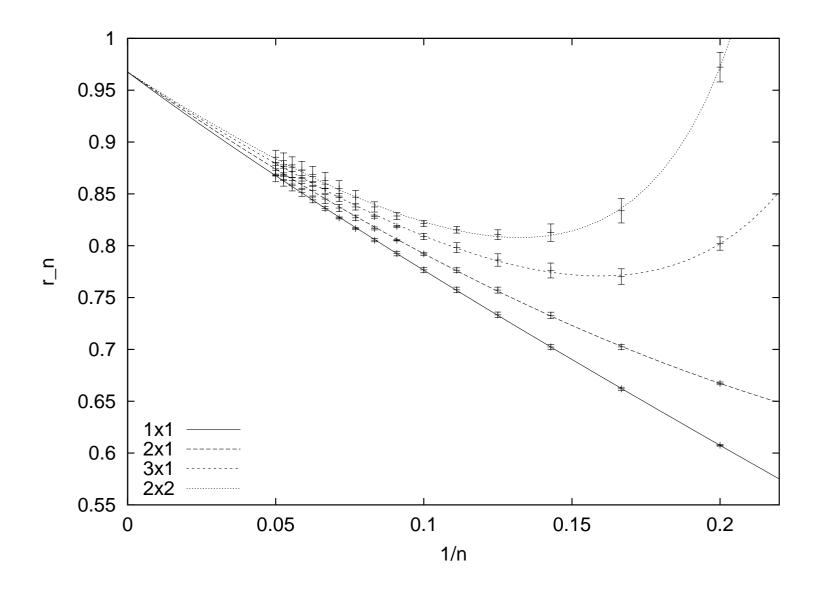


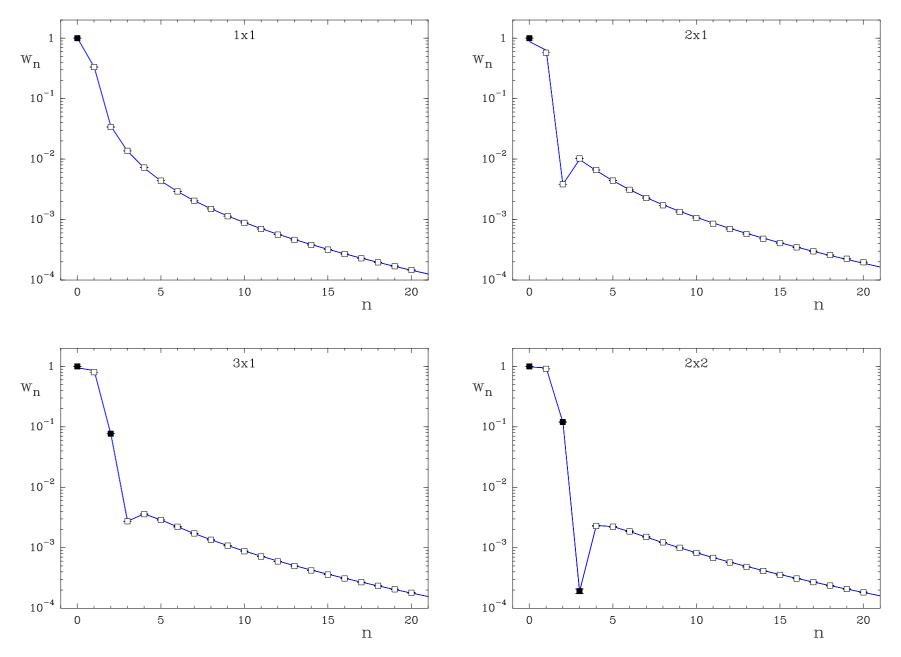


Domb-Sykes plots all look very like hyperbolae. Can get good fits with formula of the type

$$r_n = u \frac{n^2 + an + b}{n(n+s)}$$

This includes our plaquette fit, if b = 0. Fits all four Wilson loops well (errors  $<\frac{1}{2}\%$ ) for  $n \ge 5$ . plaquette: "converges"  $g^2 < 1.04$ , other loops:  $g^2 < 1.05$ .





The new fit Ansatz works surprisingly well, if we ask it to fit all the data, down to n = 0, it manages within  $\approx 5\%$ .

# **Boosted Perturbation Theory**

Perturbation theory in the lattice coupling  $g^2$  converges slowly.

Explanation:

The  $\Lambda_{QCD}$  parameter of  $g^2$  is surprisingly small  $\Lambda_{lat} = 0.035 \Lambda_{\overline{MS}}$  so  $g^2$  is not a very good expansion parameter. Remedy: Use a "boosted" coupling which has a more reasonable  $\Lambda_{QCD}$ . Popular choice

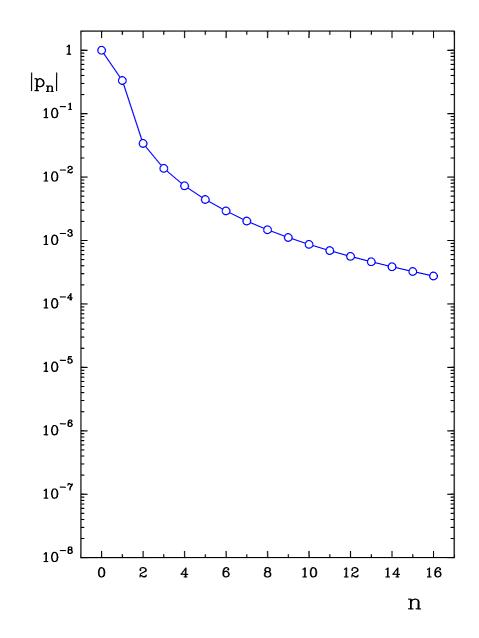
$$g_{\Box}^2 \equiv \frac{g^2}{P}$$

which has  $\Lambda_{lat} = 0.38 \Lambda_{\overline{MS}}$ 

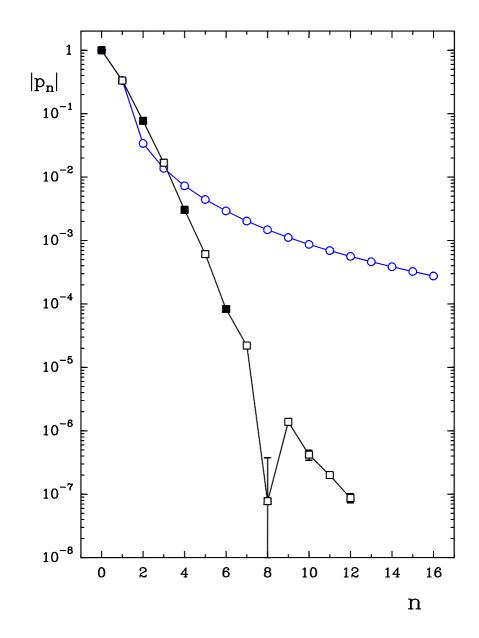
## **Boosted Perturbation Theory**

The boosted coupling is larger, but we hope that the coefficients will now decrease faster, and that we will find a reliable answer with fewer terms in the series. If we had the conventional series to very high accuracy, transforming to the new coupling would be easy. However if we only have limited accuracy in the coefficients the transformation is numerically difficult (we are looking for small differences between large quantities). Modify Stochastic Perturbation theory so that it works directly with a boosted coupling.

## **Boosted Perturbation Theory**



## **Boosted Perturbation Theory**



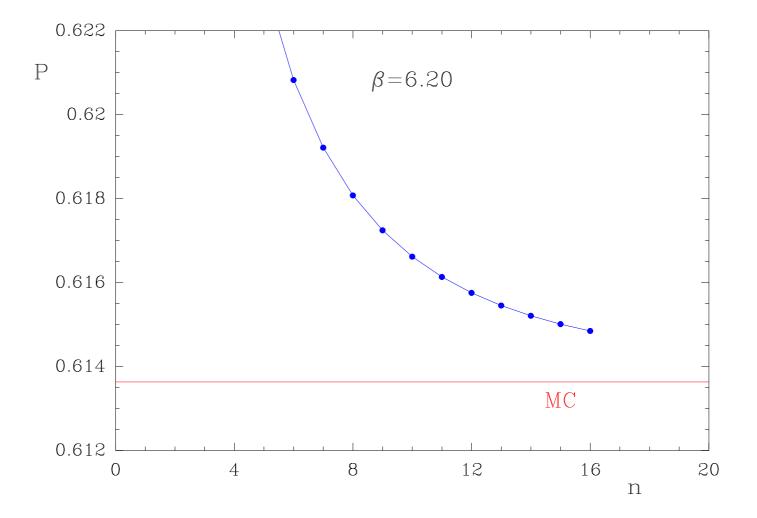
## **Boosted Perturbation Theory**

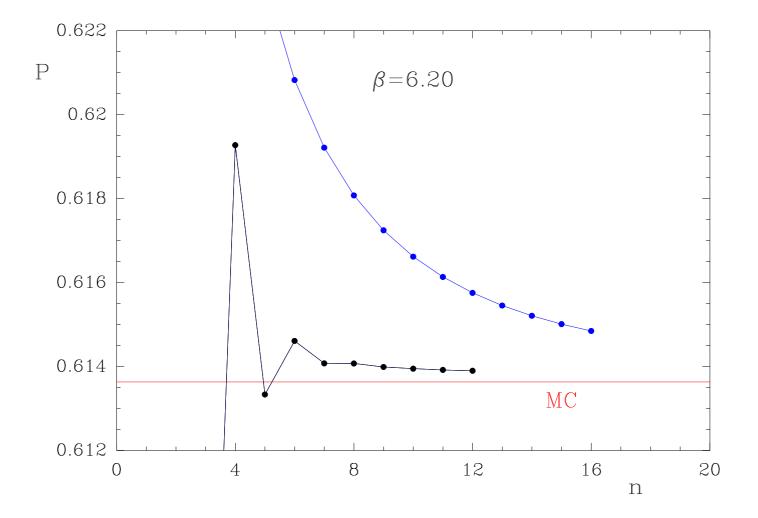
Boosted series falls off rapidly.

Even though we are interested in regions where the boosted coupling is 50% to 100% larger the change is worth while.

We are now ready to compare perturbation theory with Monte Carlo data, to see if there is any remainder which we could attribute to a condensate.

If there is, how does the condensate scale?





Interesting quantity:

Difference between perturbation theory and MC data.

 $\Delta P \equiv P_{pert} - P_{MC}$ 

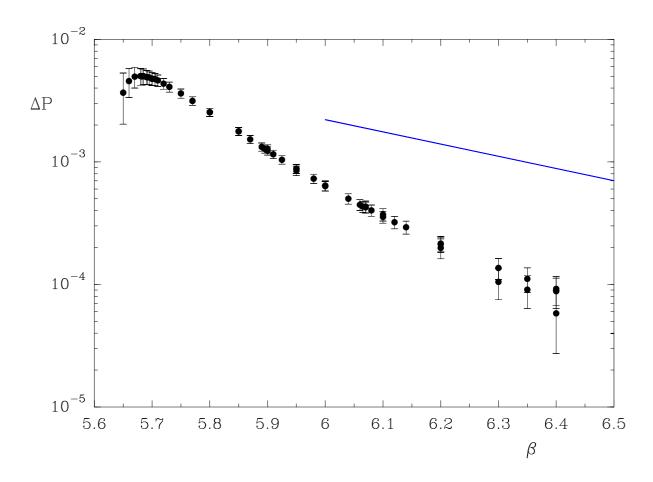
 $10^{-2}$ ΔP  $10^{-3}$ Ī  $10^{-4}$  $10^{-5}$ 5.9 6.2 6.3 5.6 5.7 5.8 6 6.1 6.4 6.5 β

 $\Delta P \equiv P_{pert} - P_{MC}$ 

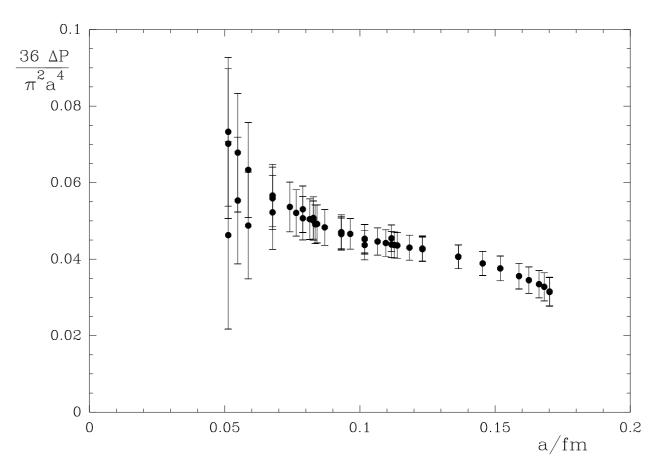
Signal below  $\beta \approx 6.4$ .

 $\Delta P$  is considerably smaller than original estimate of Burgio et al.

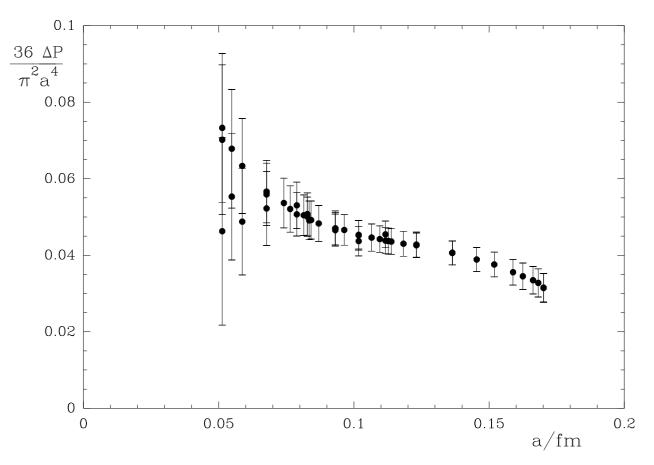
Steeper:  $a^p$  with  $p \approx 3.7(3)$ .



$$\frac{12N_c}{\pi^2 a^4} \Delta P = \left[\frac{-b_0 g^3}{\beta(g)}\right] \left\langle \frac{\alpha}{\pi} G G \right\rangle$$



Gluon condensate  $\approx 0.04 \text{ GeV}^4$ .



## Conclusions

- Stochastic Perturbation theory gives long series for some simple quantities, allows us to test ideas that we couldn't attempt conventionally.
- See no signs of series divergence yet.
- Boosted perturbation theory does accelerate convergence.
- Condensate in the plaquette seems to be dimension 4, not dimension 2.
- Value of gluon condensate  $\approx 0.04 \text{ GeV}^4$ , but with large uncertainties.