

Charged-Current Deep Inelastic Scattering at α_s^3

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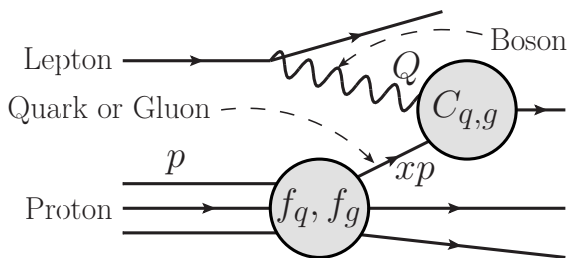
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INTRODUCTION

Deep Inelastic Scattering – lepton scatters from a proton:



The Boson can be γ, H, Z^0 (NC) or W^\pm (CC).

Cross-section $\sigma \sim \sum_a F_a = \sum_a [C_{a,q} \otimes f_q + C_{a,g} \otimes f_g]$.

$C_{a,j}$ – “Coefficient Function”
 f_j – “Parton Distribution Function”
 F_a – “Structure Function”

KNOWN SO FAR

Natural to consider *combinations* $F_a^{\nu P + \bar{\nu} P}$ and $F_a^{\nu P - \bar{\nu} P}$.

Known, to third-order (massless DIS):

- ▶ γ probe: $C_{2,q}, C_{2,g}, C_{L,q}, C_{L,g}$. [hep-ph/0504.242]
- ▶ H probe: $C_{\phi,q}, C_{\phi,g}$. [hep-ph/0912.0369]
- ▶ W^\pm probe (+): $C_{2,q}, C_{2,g}, C_{L,q}, C_{L,g}, C_{3,q}$. [hep-ph/0812.4168]

Unknown/Incomplete:

- ▶ Z^0 probe: $C_{2,q}, C_{2,g}, C_{L,q}, C_{L,g}, C_{3,q}$.
- ▶ W^\pm probe (-): $C_{2,q}, C_{L,q}, C_{3,q}$. (No gluon dependence!)

Useful in neutrino detection, for e.g., IceCube.

Probes at small- x . NLO bad here. [hep-ph/1505.06583]

NuTeV anomaly, $\sin^2 \theta_W$.

CHARGED-CURRENT DIS

Previous work at α_s^3 . Moch, Rogal:

[[hep-ph/0704.1740](https://arxiv.org/abs/hep-ph/0704.1740)]

- ▶ $N = 3, 5, \dots, 11$ of $C_{2,q}, C_{L,q}$.
- ▶ $N = 2, 4, \dots, 10$ of $C_{3,q}$.

Used to produce approximations of analytic result.

(We compute in Mellin space. N th Mellin Moment:

$$F_a(N, Q^2) = \int_0^1 x^{N-1} F_a(x, Q^2) dx.$$

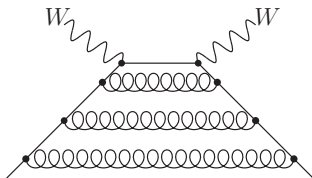
For $F_a^{\nu P - \bar{\nu} P}$: F_2, F_L live on odd moments, F_3 on even moments.)

THE OPTICAL THEOREM

Relates structure functions to the *forward scattering amplitude*:

$$\left| \begin{array}{c} \text{wavy line} \\ \text{diagonal line} \\ \text{dot} \\ \text{diagonal line} \\ \text{dot} \\ \text{horizontal line} \\ \text{dot} \\ \text{diagonal line} \end{array} \right|^2 \sim \text{Im} \begin{array}{c} \text{wavy line} \\ \text{diagonal line} \\ \text{dot} \\ \text{diagonal line} \\ \text{dot} \\ \text{horizontal line} \\ \text{dot} \\ \text{diagonal line} \\ \text{dot} \\ \text{diagonal line} \\ \text{wavy line} \end{array}$$

Calculations proceed via loop integrals rather than phase-space integrals.



+ 208 more (F_2, F_L) + 196 more (F_3), after applying symmetries.

Big improvement on diagram database of Moch, Rogal (1300+)

Check: gauge parameter in propagators, ensure cancellation.

RENORMALIZATION

Forward scattering amplitude T_a will contain poles in Dimensional-Regularization parameter ε (where $D = 4 - 2\varepsilon$), due to loop integrals (UV) and initial state collinear momenta.

Structure functions F_a are physically measurable.

- ▶ Poles are *unphysical*: renormalize.
- ▶ UV Poles: coupling constant renormalization.

T_a can then be *mass factorized*. That is, we can write

$$F_a = T_a \left(\frac{1}{\varepsilon}, \varepsilon \right) \hat{f}_q = C_a(\varepsilon) Z \left(\frac{1}{\varepsilon} \right) \hat{f}_q = C_a(\varepsilon) f_q$$

The poles can then be absorbed into a redefinition of the non-perturbative parton distribution function \hat{f}_q .

Z is a function of the (already known) anomalous dimension γ_{qq} . Recovering these is a strong consistency check.

TWO WAYS TO PROCEED...

- ▶ Extend moment-based calculation of SM, MR and reconstruct all- N expressions in the style of Polarized Splitting Function calculation of Moch, Vogt, Vermaseren.

[[hep-ph/1409.5131](https://arxiv.org/abs/hep-ph/1409.5131)]

- ▶ Direct all- N calculation. FORM code to do integral reduction exists, but did not run with modern versions of FORM.

MOMENT-BASED RECONSTRUCTION

Compute moments:

- ▶ $N = 3, 5, \dots, 29$ of $C_{2,q}, C_{L,q}$.
- ▶ $N = 2, 6, \dots, 30$ of $C_{3,q}$. (Not easy! Total 14 years CPU time.
Hardest diagram: 7TB intermediate expressions.
Each additional N takes $\sim 2x$ runtime and disk space of previous.)

Choose a basis of functions:

- ▶ All- N results in (massless) DIS can be written in terms of *Harmonic Sums* combined with simple denominators $1/(N \pm c)$.

$$S_m(N) = \sum_{i=1}^N \frac{1}{i^m}, \quad S_{-m}(N) = \sum_{i=1}^N \frac{(-1)^i}{i^m}, \quad S_{n,m}(N) = \sum_{i=1}^N \frac{1}{i^n} S_m(i)$$

Solve! (LLL Lattice Basis Reduction)

MOMENT-BASED RECONSTRUCTION

Eg: α_s^1 contribution to $C_{2,q}$. Choose basis (up to weight 2):

$$\left(A_1 + \frac{A_2}{N} + \frac{A_3}{N+1} + \frac{A_4}{N^2} + \frac{A_5}{(N+1)^2}\right) + \left(A_6 + \frac{A_7}{N} + \frac{A_8}{N+1}\right) S_1(N) + A_9 S_{1,1}(N) + A_{10} S_2(N)$$

Compute moments: ($N = 3, 5$)

$$A_1 + \frac{A_2}{3} + \frac{A_3}{4} + \frac{A_4}{9} + \frac{A_5}{16} + \frac{11A_6}{6} + \frac{11A_7}{18} + \frac{11A_8}{24} + \frac{85A_9}{36} + \frac{49A_{10}}{36} = \frac{29}{12}$$

$$A_1 + \frac{A_2}{5} + \frac{A_3}{6} + \frac{A_4}{25} + \frac{A_5}{36} + \frac{137A_6}{60} + \frac{137A_7}{300} + \frac{137A_8}{360} + \frac{12019A_9}{3600} + \frac{5269A_{10}}{3600} = \frac{589}{90}$$

+ 2 more, *Diophantine* system yields the solution:

$$\{A_1, A_2, A_3, A_4, A_5, A_6, A_7, A_8, A_9, A_{10}\} = \{-9, 3, 4, 2, 0, 3, -2, 2, 4, -4\} \checkmark$$

MOMENT-BASED RECONSTRUCTION

For terms proportional to n_f , this method worked. Basis is not too big (n_f reduces weight).

n_f^0 terms: have ~ 200 unknowns

- ▶ smaller than the maximal set – have applied constraints

System of 13 (+1) Diophantine equations.

Tried various implementations, no success.

- ▶ Basis is too big!

(Andreas used 8 equations for 41 unknowns in Polarized Splitting Function calculation.)

DIRECT ALL-N CALCULATION

Meanwhile... Andreas was speaking to Jos Vermaseren regarding FORM 4.1 and running the all- N code.

Success!

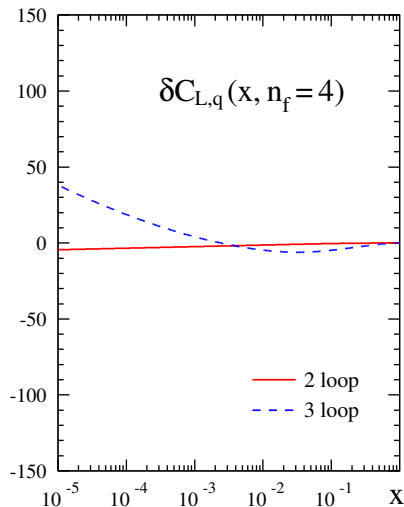
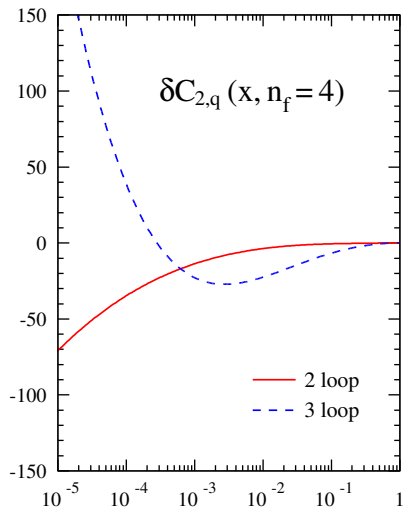
- ▶ After various bug-fixes/tweaks, the code runs.
- ▶ We have a result for n_f^0 terms!

The computed moments have not been a waste of time/resources

- ▶ Provide a verification that the code is running *correctly*.
- ▶ Techniques developed will be used in future calculations where a direct all- N calculation will *not* be possible.

RESULTS

Difference between $F_{2,L}^\gamma$ and $F_{2,L}^{\nu P-\bar{\nu}P}$ Coefficient Functions:



PASCHOS-WOLFENSTEIN RELATION

Ratio between NC and CC cross-sections:

$$R^- = \frac{\sigma(\nu P \rightarrow \nu X) - \sigma(\bar{\nu} P \rightarrow \bar{\nu} X)}{\sigma(\nu P \rightarrow \mu^- X) - \sigma(\bar{\nu} P \rightarrow \mu^+ X)}$$

Determines $\sin^2 \theta_W$.

Can now provide an α_s^3 correction (uses 2nd moments of $\delta C_{2,q}$, $\delta C_{L,q}$):

$$R^- = \left(\frac{1}{2} - \sin^2 \theta_W \right) + \frac{u^- - d^- + c^- - s^-}{u^- + d^-} \left\{ 1 - \frac{7}{3} \sin^2 \theta_W + \left(\frac{1}{2} - \sin^2 \theta_W \right) \frac{8\alpha_s}{9\pi} \left[1 + 1.689\alpha_s + 3.661\alpha_s^2 \right] \right\}$$

16% correction of $[\]$, 1% correction of $\{ \}$ over α_s^2 contribution.

CONCLUSION AND FUTURE WORK

Remaining α_s^3 corrections to CC-DIS Coefficient Functions computed.

Future work:

- ▶ Large- x , small- x resummation of leading logarithms in the new Coefficient Functions (last year's talk)
- ▶ Polarized CC-DIS
- ▶ Complete NC-DIS (Z^0 exchange)
- ▶ Photon structure with polarization (reconstruction from moments)