Quasi-Realistic Heterotic String Vacua Left Right Symmetric Model

Glyn Harries

In collaboration with Alon Faraggi & Hasan Sonmez

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Glyn Harries (UoL)

Quasi-Realistic Heterotic String Vacua

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Outline

- Introduction
- Free Fermionic Construction
- ABK Rules and GSO Projections
- Current project and results
- Conclusion

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- The motivation of this project is to create quasi-realistic string vacua
- This project uses the free fermionic construction of heterotic superstring theory
- The basis vectors chosen produce a Left Right symmetric model
- Therefore the visible gauge group at the string scale is $SU(3) \times U(1) \times SU(2)_L \times SU(2)_R$

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- $\mathcal{N} = 1$ Supersymmetry
- 3 Chiral Generations of Matter

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- Instead of associating the degrees of freedom needed to cancel the conformal anomaly as spacetime dimensions, we can interpret them as free fermions which propagate on the string worldsheet
- The string worldsheet can be mapped to a genus g Riemann surface
- We are interested in the partition function which is the integrand of the vacuum to vacuum amplitude
- Therefore we are considering a g = 1 Riemann surface *i.e* a torus

When the fermions are propagated around the two incontractible loops they pick up a phase

$$f \to -e^{i\pi\alpha(f)}f$$
 where $\alpha(f) \in (-1,1]$



Assigning different boundary conditions to each of the fermions around these loops results in different models

Free Fermionic Construction

The states on the worldsheet are

	Label	Description
Left-moving	X^{μ}	Bosonic coordinates with spacetime index, $\mu = 0, \dots, 3$
	ψ^{μ}	Majorana–Weyl superpartners of the bosonic coordinates with spacetime index
	$\chi^{1,,6}$	Majorana–Weyl superpartners to the six compactified di- mensions
	$y^{1,,6}, w^{1,,6}$	Majorana–Weyl fermions that correspond to the bosons describing the six compactified dimensions in the bosonic formulation
Right-moving	\overline{X}^{μ}	Bosonic coordinates with spacetime index
	$\overline{y}^{1,\dots,6}, \overline{w}^{1,\dots,6}$	Majorana–Weyl fermions that correspond to the bosons describing the six compactified dimensions in the orbifold formulation
	$\overline{\psi}^{1,\dots,5},\overline{\eta}^{1,2,3}$	Complex fermions that describe the visible gauge sector
	$\overline{\phi}^{1,,8}$	Complex fermions that describe the hidden gauge sector

There are 18 free fermions in the left moving supersymmetric sector and 44 free fermions in the right moving sector

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A model is defined by specifying two ingredients

- A set of boundary condition basis vectors
- The one loop phases $C \begin{pmatrix} b_i \\ b_j \end{pmatrix}$ for all pairs of the basis vectors

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- The basis vectors and one loop coefficients must satisfy the ABK rules
- These are derived from modular invariance conditions

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The Spacetime Spin Statistics Index is

$$\delta_{b_i} = e^{i\pi b_i(\psi^{\mu})} = \begin{cases} -1 & b_i(\psi^{\mu}) = 1\\ +1 & b_i(\psi^{\mu}) = 0 \end{cases}$$

- If the basis vector specifies ψ^{μ} is periodic then $\delta_{b_i}=-1$
- If the basis vector specifies ψ^{μ} is anti-periodic then $\delta_{b_i}=0$

• The equation for the GSO projection is

$$e^{i\pi b_i\cdot F_\alpha}\left|s\right\rangle_\alpha = \delta_\alpha C \begin{pmatrix} \alpha\\ b_i \end{pmatrix}^* \left|s\right\rangle_\alpha$$

- This selects the states that are either kept in or projected out of the spectrum
- If the equation is satisfied by a state then it is kept, else it is projected out.

Current Project

The current project has the basis vectors

$$\begin{split} &\mathbb{I} = \{\psi_{1,2}^{\mu}, \chi^{12}, \chi^{34}, \chi^{56}, y^{12}, y^{34}, y^{56}, w^{12}, w^{34}, w^{56} \mid \bar{y}^{12}, \bar{y}^{34}, \bar{y}^{56}, \bar{w}^{12}, \\ &\bar{w}^{34}, \bar{w}^{56}, \bar{\psi}^{1,\dots,5}, \bar{\eta}^{1,2,3}, \bar{\phi}^{1,\dots,8} \} \\ &S = \{\psi_{1,2}^{\mu}, \chi^{12}, \chi^{34}, \chi^{56} \} \\ &e_i = \{y^i, w^i \mid \bar{y}^i, \bar{w}^i \} \\ &b_1 = \{\chi^{34}, \chi^{56}, y^{34}, y^{56} \mid \bar{y}^{34}, \bar{y}^{56}, \bar{\psi}^{1,\dots,5}, \bar{\eta}^1 \} \\ &b_2 = \{\chi^{12}, \chi^{34}, y^{12}, y^{56} \mid \bar{y}^{12}, \bar{y}^{56}, \bar{\psi}^{1,\dots,5}, \bar{\eta}^2 \} \\ &z_1 = \{\bar{\phi}^{1,\dots,4} \} \\ &z_2 = \{\bar{\phi}^{5,\dots,8} \} \\ &z_3 = \{\bar{\phi}^{1,2}, \bar{\phi}^{7,8} \} \\ &\alpha = \{\bar{\psi}^{1,2,3} = \frac{1}{2}, \bar{\eta}^{1,2,3} = \frac{1}{2}, \bar{\phi}^{1,2} = \frac{1}{2} \} \end{split}$$

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Matter Spectrum B_{pqrs}

- The observable matter spectrum can be calculated by performing the GSO projections on B_{pqrs} which is a linear combination of basis vectors
- For example

$$B_{pqrs}^{(1)} = S + b_1 + pe_3 + qe_4 + re_5 + se_6$$

= $\{\psi^{\mu}, \chi^{1,2}, (1-p)y^3 \bar{y}^3, pw^3 \bar{w}^3, (1-q)y^4 \bar{y}^4, qw^4 \bar{w}^4, (1-r)y^5 \bar{y}^5, rw^5 \bar{w}^5, (1-s)y^6 \bar{y}^6, sw^6 \bar{w}^6, \bar{\eta}^1, \bar{\psi}^{1,\dots,5}\}$

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- Under the GSO projection of the basis vector b_1 the fermions $\{\bar{\eta}^1,\bar{\psi}^{1,\dots,5}\}$ are isolated
- Under b_2 this splits to give $\{\bar{\eta}^1\}, \{\bar{\psi}^{1,\dots,5}\}$
- Performing the α GSO projection splits this into the Left Right Symmetric model

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Matter Spectrum $B_{pqrs}^{(1)}$



• The coding is written in Java

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- It performs the GSO projections and scans for vacua which are consistent with the contraints specified
- Currently the observable matter spectrum is being written
- The program can also check for vacua criteria such as light or heavy Higgs, exotic matter states *etc.*

- The choice of basis vectors generates string models which are left right symmetric
- The program currently does give models with ${\cal N}=1$ supersymmetry and 3 chiral generations of matter
- Further work to be completed is to fully complete the section of the program which tests the matter spectrum
- The program must also be extended to correctly test for light and heavy Higgs particles, exotic states and gauge group enhancements

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Conclusions

Thank you for listening

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References

ABK rules:

I. Antoniadis and C. Bachas. 4d fermionic superstrings with arbitrary twists. Nuclear Physics B, 298(3):586 - 612, 1988.
I. Antoniadis and C. Bachas, and C. Kounnas. Four-dimensional

superstrings. Nuclear Physics B, 289(0):87 - 108, 1987

• Figures 1 and 2:

"Light U(1)'s in Heterotic-string Models' - Viraf M. Mehta - September 2013

- "Classification of the Flipped SU(5) Heterotic-String Vacua" Hasan Sonmez String Phenomenology 2014 Trieste Talk
- "Semi-Realistic Heterotic-String Vacua" Johar M. Ashfaque String Theory Seminar May 2015

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