Dirac Neutrinos and a Vanishing Higgs at the LHC

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with Athanasios Dedes and David Cerdeño
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also with Frank Krauss and Terrance Figy
hep-ph/to appear

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Introduction

- Minimal Lepton Number Conserving Phantom Sector
- "Phantom" $\rightarrow$ singlet under the Standard Model gauge group SU(3)$_c$×SU(2)$_L$×U(1)$_Y$
- Simple model leading to interesting phenomenology:

  - Dirac Neutrino Masses
  - Dirac Leptogenesis
  - Higgs Phenomenology

Outline

- Dirac Neutrino Masses
- Dirac Leptogenesis
- Higgs Phenomenology

Model building

- Just 2 openings in the SM for renormalisable operators coupling $SU(3)_c \times SU(2)_L \times U(1)_Y$ singlet fields to SM fields[1]

- Higgs mass term: $H^\dagger H\,?^*?$
- Lepton-Higgs Yukawa interaction: $\bar{L}\,\widetilde{H}\,?_R$

- What would happen if we filled in the gaps?
- But, no evidence for $B-L$ violation yet, so could try to build a $B-L$ conserving model
- Will try to be "natural" in the 't Hooft and the aesthetic sense - couplings either $\mathcal{O}(1)$ or strictly forbidden

[1] B. Patt and F. Wilczek, hep-ph/0605188

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- Augment the SM with two SU(3)$_c$×SU(2)$_L$×U(1)$_Y$ singlet fields
  - a complex scalar Φ
  - a Weyl fermion s_R

$$-\mathcal{L}_{\rm link} = \left(h_\nu\, \overline{l_L} \cdot \widetilde{H}\, s_R \,+\, \text{H.c.}\right) - \eta\, H^\dagger H\;\Phi^*\Phi$$

$\widetilde{H} = i\sigma_2 H^*$,
h_ν and η will be $\mathcal{O}(1)$,
s_R carries lepton number $L = 1$.

- But, this model is no good $\rightarrow$ neutrinos would have large, electroweak scale masses

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- Solution: Postulate the existence of a purely gauge singlet sector; add ν_R and s_L.

$$-\mathcal{L}_{\rm p} = h_{\rm p}\, \Phi\, \overline{s_L}\, \nu_R \;+\; M\, \overline{s_L}\, s_R \;+\; \text{H.c.}$$

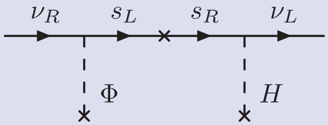
- Forbid other terms by imposing a "phantom sector" global $\mathrm{U(1)_D}$ symmetry, such that only

$$\nu_R \to e^{i\alpha}\, \nu_R \quad , \quad \Phi \to e^{-i\alpha}\, \Phi$$

transform non-trivially
- If we require small Dirac neutrino masses this is the simplest choice for the phantom sector

$$\mathcal{L} = \mathcal{L}_{\rm SM} \;+\; \mathcal{L}_{\rm link} \;+\; \mathcal{L}_{\rm p}$$

Small effective Dirac neutrino masses – Dirac See-Saw



- Spontaneous breaking of both SU(2)$_L\times$U(1)$_Y$ and U(1)$_\mathrm{D}$ will result in the effective Dirac mass terms

$$-\mathcal{L} \supset \overline{\nu'_L}\,\mathbf{m}_\nu\,\nu'_R + \overline{s'_L}\,\mathbf{m_N}\,s'_R$$

assuming $M \gg v$ and where

$$\mathbf{m}_\nu = -v\,\sigma\,\mathbf{h}_\nu\,\hat{\mathbf{M}}^{-\mathbf{1}}\,\mathbf{h}_\mathrm{p} \qquad\qquad \mathbf{m_N} = \hat{\mathbf{M}}$$

with $\sigma \equiv \langle\Phi\rangle$ and $v \equiv \langle H\rangle = 175$ GeV.

Essentially the Froggatt-Nielsen mechanism!

C. D. Froggatt and H. B. Nielsen, NPB**147**(1979)277.

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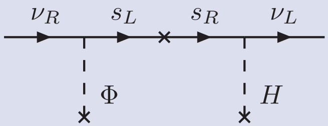

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M. Roncadelli and D. Wyler, PLB**133**(1983)325

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- Higgs Phenomenology

We can measure the baryon asymmetry of the universe but do we understand where it came from?

Sakharov's famous conditions

- Baryon number violation
- C and CP violation
- Conditions out of thermal equilibrium

Leptogenesis is commonly cited as a possible explanation

- In the SM, $B + L$ violation occurs at high temperatures allowing a lepton asymmetry to be partially converted to a baryon asymmetry
- In the Majorana see-saw, lepton number and CP are generally violated in the decays of the heavy Majorana neutrinos
- These decays can occur out of thermal equilibrium

M. Fukugita and T. Yanagida, PLB**174**(1986)45

This model exactly conserves $B - L$, so it seems we cannot create a lepton asymmetry in the same way.
However

- $B + L$ violation in the SM does not directly affect right handed gauge singlet particles
- Small effective Yukawa couplings between the left and right handed neutrinos could prevent asymmetries in this sector from equilibrating
  - L_{ν_R} could "hide" from the rapid $B + L$ violating processes

V. A. Kuzmin, hep-ph/9701269
K. Dick, M. Lindner, M. Ratz and D. Wright, PRL**84**(2000)4039

see also: H. Murayama and A. Pierce, PRL**89**(2002)271601
S. Abel and V. Page, JHEP**0605**(2006)024
B. Thomas and M. Toharia, PRD**73**(2006)063512

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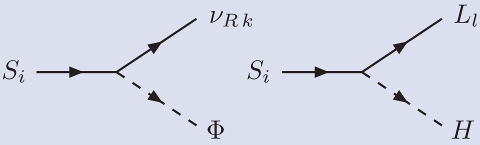
Generation of the L_{ν_R} ($L_{\rm SM}$) asymmetry



$S \equiv s_L + s_R$

- Heavy particle decay – similar to Majorana leptogenesis
- In analogy with Davidson and Ibarra, the CP-asymmetry is bounded

$$|\delta_{R1}| \lesssim \frac{1}{16\pi}\,\frac{M_1}{v\,\sigma}\,(m_{\nu_3} - m_{\nu_1})$$

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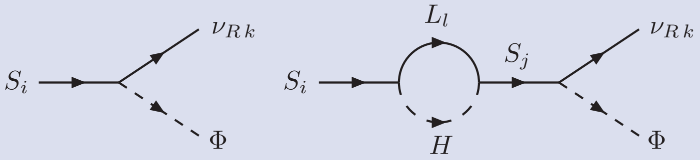

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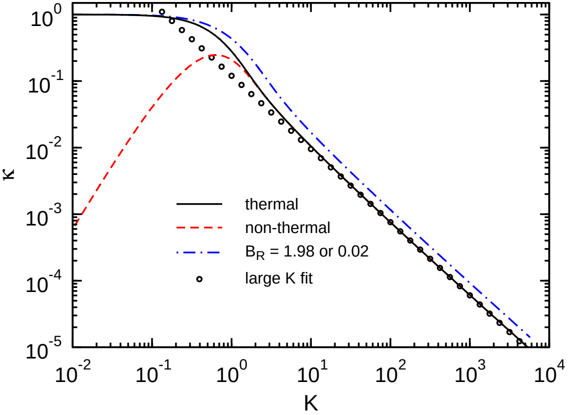
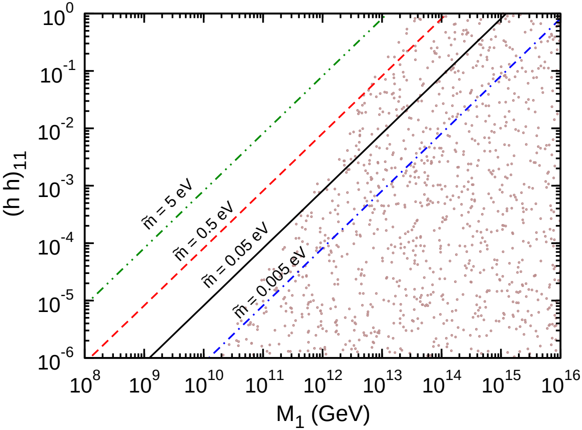


Leptogenesis efficiency, κ, versus K for thermal and zero initial abundance of S_1 ($\bar{S}_1$). Also shown is the efficiency for differing left-right branching ratios.



Area in the M_1, $(\mathbf{h}^\dagger \mathbf{h})_{11}$ parameter space allowed by successful baryogenesis when $(\mathbf{h}_\nu^\dagger \mathbf{h}_\nu)_{11} = (\mathbf{h}_\mathrm{p} \mathbf{h}_\mathrm{p}^\dagger)_{11}$ and $\sigma = v = 175$ GeV.

- If we take a 'natural' scenario with $(\mathbf{h}_\nu^\dagger \mathbf{h}_\nu)_{11} = (\mathbf{h}_\mathrm{p} \mathbf{h}_\mathrm{p}^\dagger)_{11} \simeq 1$ and $\widetilde{m} = 0.05$ eV (hierarchical light neutrinos) we can use the bound on the CP-asymmetry and the observed baryon asymmetry to put a bound on σ

$$\sigma \gtrsim 0.1 \text{ GeV}$$

- If we require that S_1 be produced thermally after inflation there exists an approximate bound $M_1 \lesssim T_{RH}$.
- Given the same reasonable assumptions, this implies an approximate upper bound on σ

$$0.1 \text{ GeV} \;\lesssim\; \sigma \;\lesssim\; 2 \text{ TeV} \left(\frac{T_{RH}}{10^{16} \text{ GeV}} \right)$$

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The potential

$$V = \mu_H^2 H^* H + \mu_\Phi^2 \Phi^* \Phi + \lambda_H (H^* H)^2 + \lambda_\Phi (\Phi^* \Phi)^2 - \eta H^* H \Phi^* \Phi$$

where $H \equiv H^0$

- After spontaneous breaking of $\mathrm{U(1)_D}$, Φ develops a non-zero vev. This, through the η term, would trigger electroweak SU(2)$_L \times$U(1)$_Y$ symmetry breaking
- Expanding the fields around their minima

$$H = v + \frac{1}{\sqrt{2}}(h + iG) \quad , \quad \Phi = \sigma + \frac{1}{\sqrt{2}}(\phi + iJ)$$

- We have
  - the Goldstone bosons: G (eaten as usual) and J
  - h and ϕ mix (due to the η term) and become two massive Higgs bosons H_1 and H_2

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$$\begin{pmatrix} H_1 \\ H_2 \end{pmatrix} = O \begin{pmatrix} h \\ \phi \end{pmatrix} \quad \text{with} \quad O = \begin{pmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{pmatrix}$$

and the mixing angle

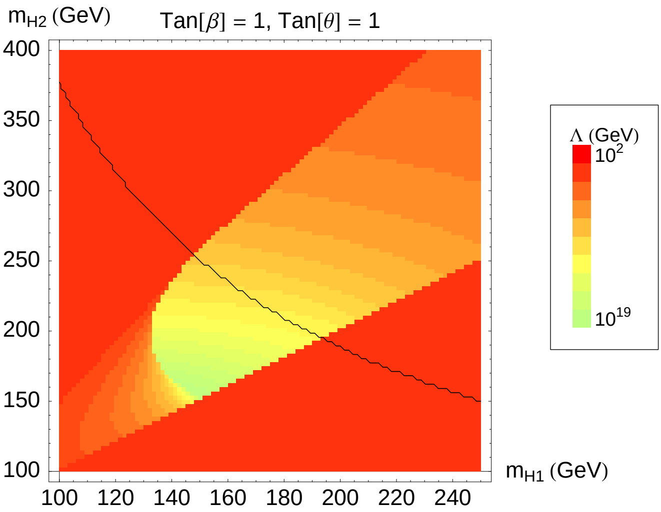
$$\tan 2\theta = \frac{\eta v \sigma}{\lambda_\Phi \sigma^2 - \lambda_H v^2}$$

- The limits $v \ll \sigma$ and $\sigma \ll v$ both lead to the SM with an isolated hidden sector
- These limits need an unnaturally small η, and would present problems with baryogenesis and small neutrino masses.
- A 'natural' choice of parameters (with e.g. $\eta \sim 1$) would lead to

$$\tan\theta \sim 1 \quad , \quad \tan\beta \equiv v/\sigma \sim 1$$

Triviality and Positivity

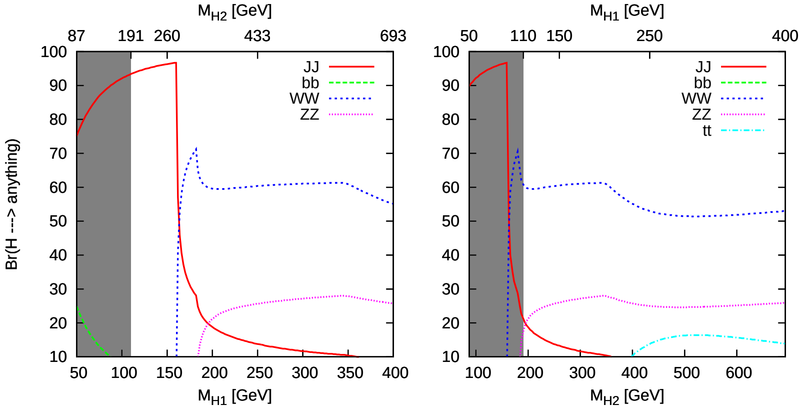
- We require that the parameters λ_H, λ_Φ and η do not encounter Landau poles at least up to the scale where we encounter “new physics”.
- We also require that the potential remain positive definite everywhere, at least up to the scale of “new physics”.
- After solving 1-loop RGEs, we can plot the maximum scale up to which our effective theory satisfies the above constraints.


m_H2 (GeV)
Tan[β] = 1, Tan[θ] = 1
400
350
300
250
200
150
100
100 120 140 160 180 200 220 240
m_H1 (GeV)
Λ (GeV)
10^2
10^19

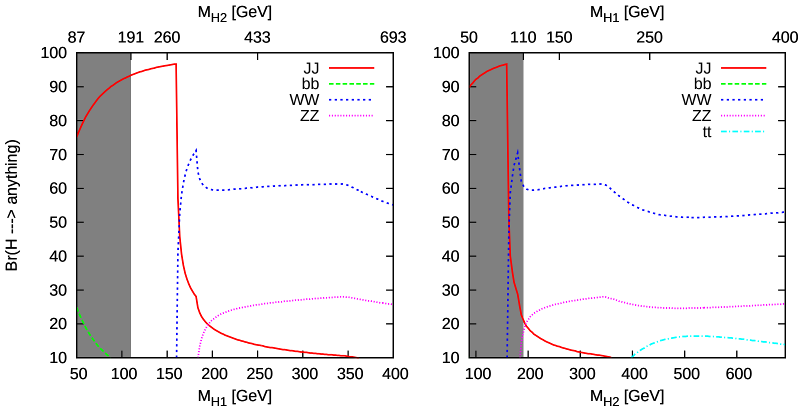
- Couplings of the H_i to SM fermions and gauge bosons will be reduced by a factor O_{i1} (relative to the SM)
- H_i will also couple to the massless Goldstone pair JJ
- In the SM, for light Higgs masses $\lesssim$ 160 GeV, $H \rightarrow b\bar{b}$ dominates. Here we find:

$$
\begin{aligned}
\frac{\Gamma(H_1 \rightarrow JJ)}{\Gamma(H_1 \rightarrow bb)} &= \frac{1}{12}\left(\frac{m_{H1}}{m_b}\right)^2 \tan^2\beta \; \tan^2\theta \\
\frac{\Gamma(H_2 \rightarrow JJ)}{\Gamma(H_2 \rightarrow bb)} &= \frac{1}{12}\left(\frac{m_{H2}}{m_b}\right)^2 \tan^2\beta \; \cot^2\theta
\end{aligned}
$$

- In this model a 'light' Higgs boson will decay dominantly into invisible JJ as long as it is heavier than 60 GeV.



Dominant branching ratios of the two Higgs bosons H_1 (left) and H_2 (right) for the parameters $\theta = \beta = \pi/4$, with couplings equal to one. The shaded area is excluded by LEP.

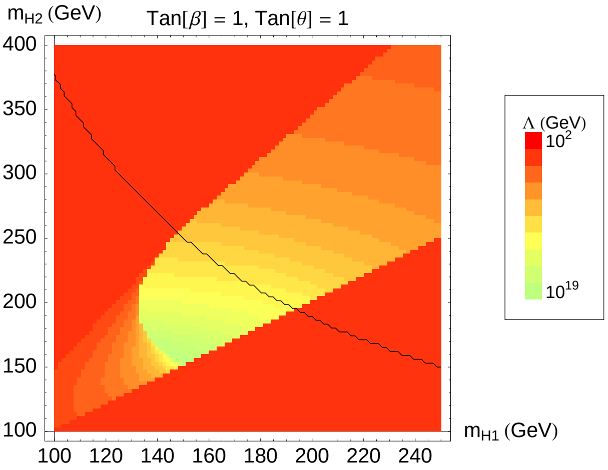


- LEP excludes a light invisible Higgs with a mass $m_{H1} \lesssim$ 110 GeV.
- It therefore sets a lower bound on the heavier Higgs $m_{H2} \gtrsim$ 191 GeV.

- Let us compare the number of Higgs events at the LHC in this model vs. the SM (for an identical Higgs mass)
- Compare numbers of visible events, in the narrow width approximation and assuming that the vector bosons produced in Higgs decays are on-shell.

Define a parameter $\mathcal{R}_i$

$$\mathcal{R}_i \equiv \frac{\sigma(pp \to H_i\, X)\,\mathrm{Br}(H_i \to YY)}{\sigma(pp \to H_{\mathrm{SM}}\, X)\,\mathrm{Br}(H_{\mathrm{SM}} \to YY)}$$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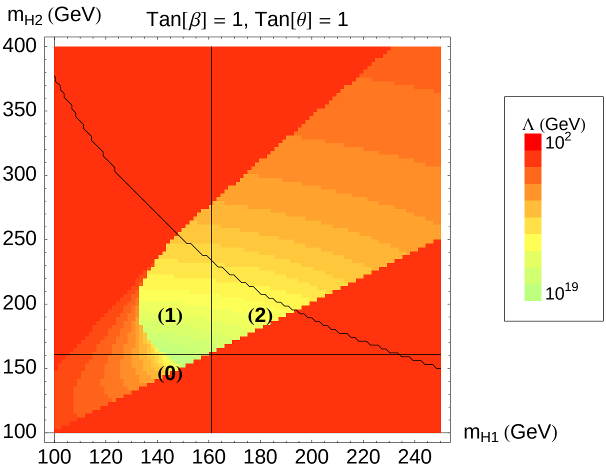
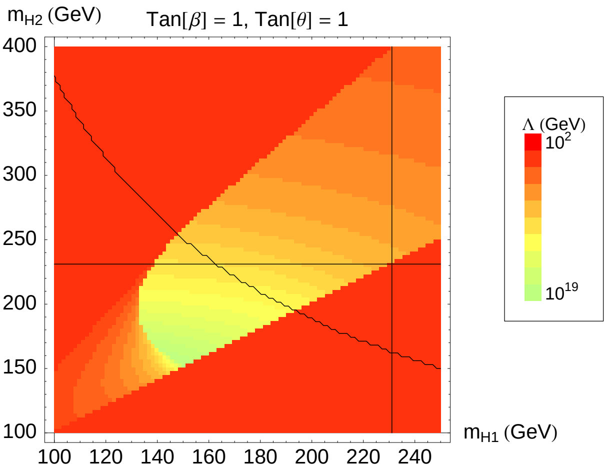
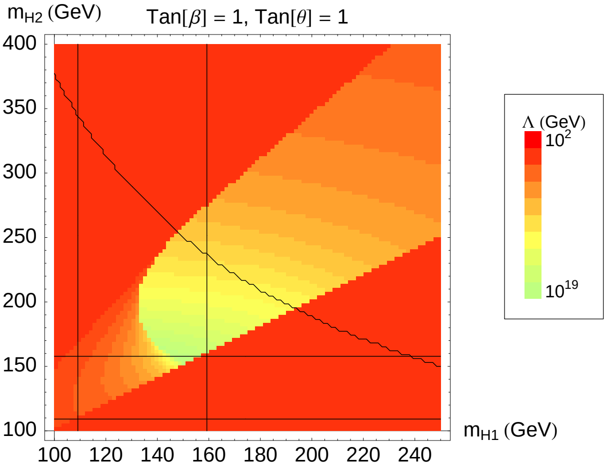


$\mathcal{R}_i = 0.1$



$\mathcal{R}_i = 0.3$



$\mathcal{R}_i = 0.01$

- There is a mass region where one, or both H_i decay to invisible JJ with $\mathrm{Br}(H_i \to JJ) > 90\%$.
- How could this Higgs be found at the LHC?

S. G. Frederiksen, N. Johnson, G. L. Kane and J. Reid, PRD**50**(1994)4244
R. M. Godbole, M. Guchait, K. Mazumdar, S. Moretti and D. P. Roy, PLB**571**(2003)184
K. Belotsky, V. A. Khoze, A. D. Martin and M. G. Ryskin, EPJC**36**(2004)503
H. Davoudiasl, T. Han and H. E. Logan, PRD**71**(2005)115007

- Strategies:
  - $Z + H_1$
  - W-boson fusion
  - central exclusive diffractive production

$Z(\to l^+l^-) + H_{\rm inv}$

using H. Davoudiasl, T. Han and H. E. Logan, PRD**71**(2005)115007

- multiply $S/\sqrt{B}$ by $1/2$ because of mixing
- assume LHC integrated luminosity of 30fb^{-1}

Signal significance for discovering the invisible H_1 is

- m_{H1} = 120 GeV 4.9σ
- m_{H1} = 140 GeV 3.6σ
- m_{H1} = 160 GeV 2.7σ

- Although this applies to $\theta = \pi/4$, the situation is rather generic in this region
- Note that for $m_{H1} \lesssim$ 140 GeV, the $H_1 \to \gamma\gamma$ channel may still be usable.

Simulation for High Energy Reactions of PArticles

- We have implemented this model in the matrix element monte carlo program SHERPA[1]
- SHERPA is built to make it “easy” to implement new physics models in a monte carlo simulation – essential for being able to talk about realistic LHC phenomenology
- Will the invisible Higgs remain invisible?

[1] F. Krauss *et al*

Summary

- Proposed a minimal, L conserving, phantom sector of the SM leading to
  - Viable Dirac neutrino masses
  - Successful baryogenesis (through Dirac leptogenesis)
  - Interesting 'invisible' Higgs phenomenology for the LHC
- In this model, $\mathcal{O}(1)$ couplings, correct neutrino masses and baryogenesis seem to suggest an electroweak scale vev in the minimal phantom sector

Other Astro/Cosmo Constraints

H_i couples to JJ as

$$-\mathcal{L}_J \supset \frac{(\sqrt{2}G_F)^{1/2}}{2} \tan\beta \, O_{i2} \, m_{H_i}^2 \; H_i \, JJ$$

- After electroweak/U$(1)_{\mathrm{D}}$ symmetry breaking the Js are kept in equilibrium via reactions of the sort $JJ \leftrightarrow f\bar{f}$ mediated by H_i
- A GIM-like suppression exists for these interactions from the orthogonality condition $\sum_i O_{i1}O_{i2} = 0$
- J falls out of equilibrium just before the QCD phase transition and remains as an extra relativistic species thereafter

- BBN/CMB yield a bound on the effective number of neutrino species $N_\nu = 3.24 \pm 1.2$ (90% C.L.)
- Early decoupling of J implies T_J is much lower than T_ν

$$\left(\frac{T_J}{T_\nu}\right)^4 = \left(\frac{g_*(T_\nu)}{g_*(T_D)}\right)^{4/3} \lesssim \left(\frac{10.75}{60}\right)^{4/3}$$

- The increase in the effective number of light neutrinos, due to J, at BBN ΔN_ν^J is then

$$\Delta N_\nu^J = \frac{4}{7}\left(\frac{T_J}{T_\nu}\right)^4 \lesssim 0.06$$