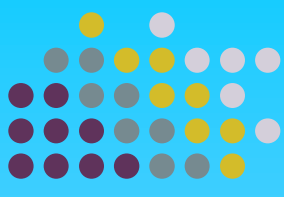


The Higgs Reduction Mechanism in Free Fermionic Models

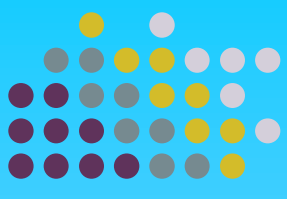


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Eur. Phys. J. C (2007)

University of Liverpool
UKBSM Workshop, 29-30 March 2007

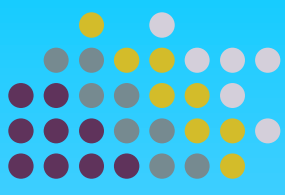


Outline



- *Motivations*
- *Free Fermionic Models (FFM)*
- *Yukawa selection mechanism*
- *Higgs doublet-triplet splitting*
- *Model with reduced Higgs spectrum*
- *Conclusions*

Motivations



In the free fermionic formulation we have realistic models

- Existence of **3 chiral generations**
- Observable gauge group : **$SU(3) \times SU(2) \times U(1)^n$**
- **N=1** Supersymmetry
- Standard **SO(10)** embedding for the weak hypercharge
- **Stability of the proton**
- ...

Heterotic string vacua with solely MSSM states are derived.

The Free Fermionic Models

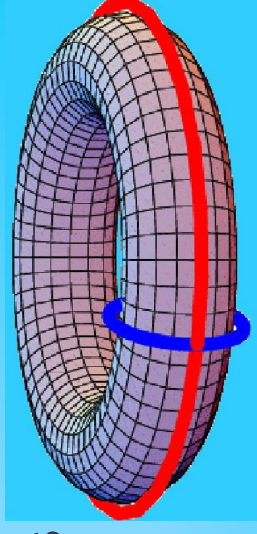
- Field content (light-cone gauge)

$$\{ \psi^{\mu}_{1,2} \chi^{1,\dots,6} y^{1,\dots,6} \omega^{1,\dots,6} | \bar{y}^{1,\dots,6} \bar{\omega}^{1,\dots,6} \bar{\Psi}^{1,\dots,5} \bar{\eta}^{1,2,3} \bar{\phi}^{1,\dots,8} \}$$

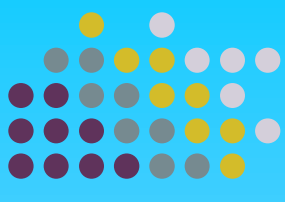
left sector (susy) right sector

- A model is constructed by specifying the *phases picked up* by the WS fermions along *non-contractible loops*

$$\mathbf{f} \rightarrow \mathbf{e}^{-i\alpha(\mathbf{f})\pi} \mathbf{f}$$



- The phases are consistent with *modular invariance*



- The phases are given in terms of *basis vectors*
 $\vec{b} = \{\alpha(\psi^{u_{12}}), \alpha(\chi^1), \alpha(\chi^2), \dots\}$
- For a given basis $B = \{\vec{b}_1, \dots, \vec{b}_n\}$
the Hilbert space $\Xi = \sum m_j \vec{b}_j$, $m_j = 1, \dots, N_{b_j} - 1$
is obtained by acting on the vacuum with bosonic and fermionic operators
- The physical spectrum is obtained by applying the *GSO projections*
- If $B = \{1, \mathbf{S}, \mathbf{b}_1, \mathbf{b}_2, \mathbf{b}_3\}$ + opportune choice of α, β, γ
b.c. basis vectors we have **4D** models with **N=1 Susy**
+ **3 chiral generations**, one from each b_i , $i=1,2,3$.

Yukawa selection mechanism

The Yukawa selection mechanism is given by the vector γ responsible of the breaking

$$SO(10) \rightarrow SU(5) \times U(1)$$

The b.c. basis vectors in γ fix the Yukawa couplings

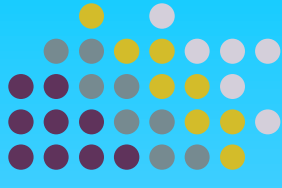
$$Q\bar{u}h, LN\bar{h} \\ Q\bar{d}h, L\bar{e}h$$

Each sector b_i gives rise to an *up-like* or *down-like* cubic level Yukawa coupl.

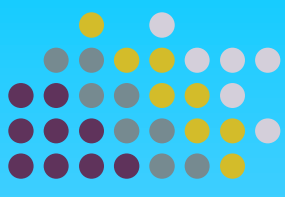
$$\Delta_i = |\gamma_L - \gamma_R| = 0, 1$$

down-quark type
Yukawa coupl.

up-quark type
Yukawa coupl.



Higgs doublet-triplet splitting



The NS sector gives 3 multiplets of Higgs states in the $\mathbf{10}$ of $\mathbf{SO}(10)$ each of them associated with one of the twisted sectors b_i .

Higgs electroweak doublets

$\overline{h}_i, \overline{h}_i$

$\chi^{i,i+1}(\Psi^{4,5})^* \eta^i$

Higgs colour triplets

$\overline{D}_i, \overline{D}_i$

$\chi^{i,i+1}(\Psi^{4,5})^* \eta^i$

The doublet-triplet splitting mechanism results from the b. c. in α which break

$$\mathbf{SO}(10) \rightarrow \mathbf{SO}(6) \times \mathbf{SO}(4)$$

$$\Delta_i = |\alpha_L - \alpha_R| = 0, 1$$

h_i projected out

h_i remains in the spectrum

We look for models such that...

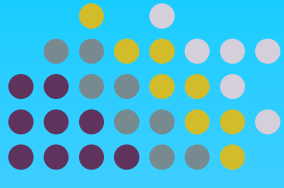
- No Higgs triplets ; one Higgs doublet

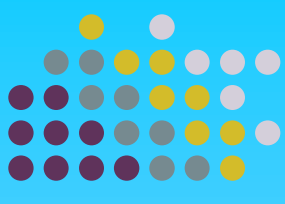
$$\Delta_3(\beta) = 1$$

- Up-quark type Yukawa couplings selected

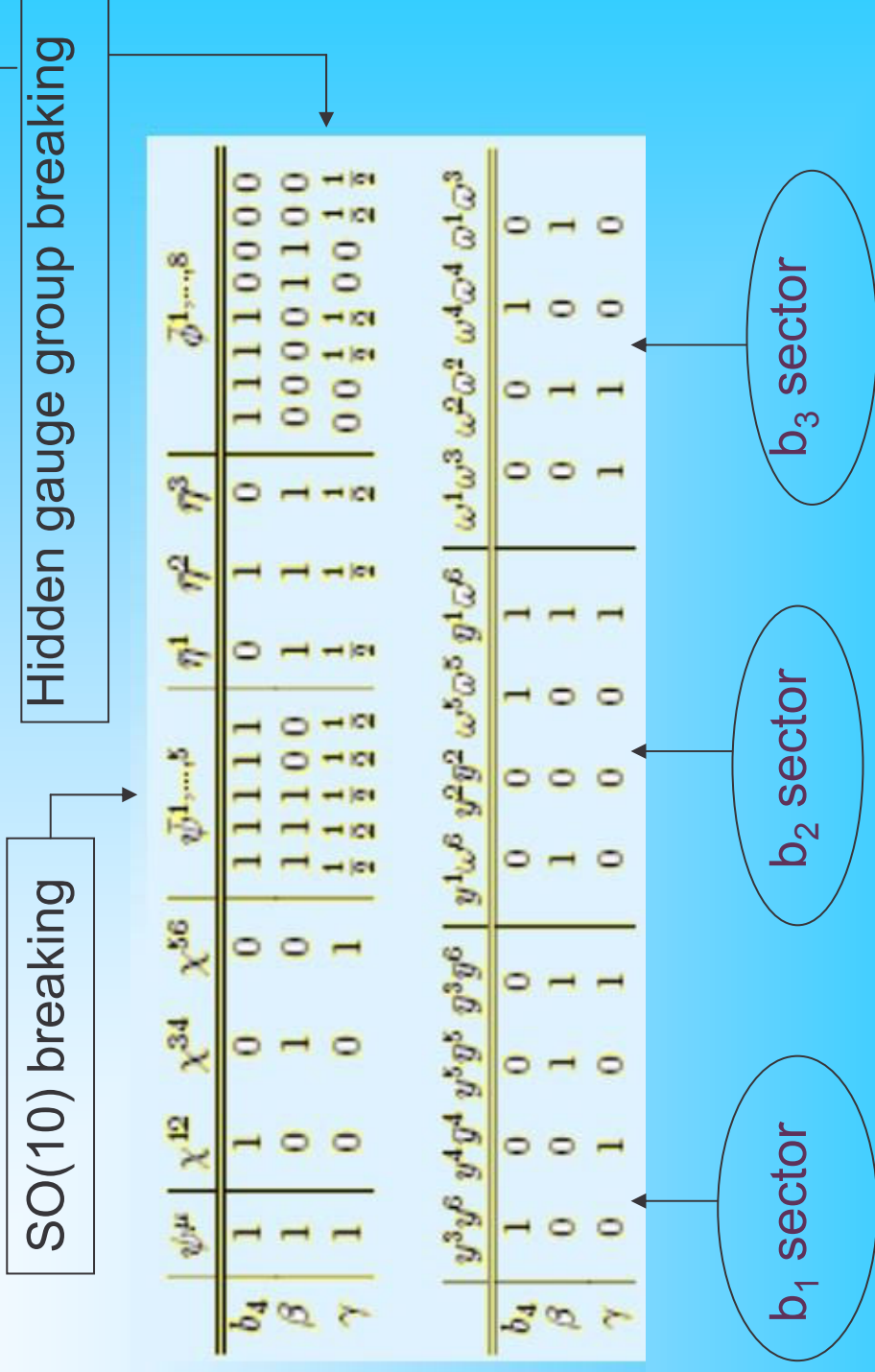
$$\Delta_3(\gamma) = 1$$

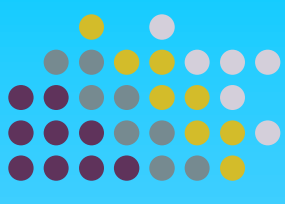
- No chiral fractionally charged exotics





Model with reduced Higgs spectrum





Model with reduced Higgs spectrum

keep h_1

keep D_2

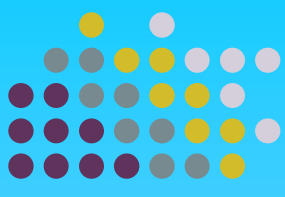
keep h_3

| | $y^3 y^6$ | $y^4 g^4$ | $y^5 g^5$ | $g^3 g^6$ | $g^1 \omega^6$ | $y^2 g^2$ | $\omega^5 \omega^5$ | $g^1 \omega^6$ | $\omega^1 \omega^3$ | $\omega^2 \omega^2$ | $\omega^4 \omega^4$ | $\omega^1 \omega^3$ |
|----------|-----------|-----------|-----------|-----------|----------------|-----------|---------------------|----------------|---------------------|---------------------|---------------------|---------------------|
| b_4 | 1 | 0 | 0 | 0 | 0 | 0 | 0 | 1 | 1 | 0 | 0 | 0 |
| β | 0 | 0 | 1 | 1 | 1 | 1 | 0 | 0 | 1 | 0 | 1 | 0 |
| γ | 0 | 1 | 0 | 1 | 0 | 0 | 0 | 0 | 1 | 1 | 0 | 0 |

Project out the untwisted vectorial repr. $\Delta_\gamma = 1$

(up-quark type Yukawa coupling is selected)

Conclusions



- Free Fermionic Models produce some of the most realistic string models to date.
The presence of 3 pairs of Higgs doublets is generally reduced by the analysis of supersymmetric flat directions.
- An alternative option for the reduction of the Higgs states is given by an opportune choice of free fermion boundary conditions at the string scale.
- The previous mechanism also provides a reduction of the moduli space.