

Long-lived charged massive particle and its effect on cosmology

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[Kawasaki, Kohri, Moroi, PLB625 \(2005\) 7](#)

[Kawasaki, Kohri, Moroi, PRD71 \(2005\) 083502](#)

[Kohri, Moroi, Yotsuyanagi, PRD73 \(2006\) 123511](#)

[Kohri, Takayama, hep-ph/0605243](#)

[Kanzaki et al, hep-ph/0609246](#)

[Kawasaki, Kohri, Moroi, hep-ph/0703122](#)



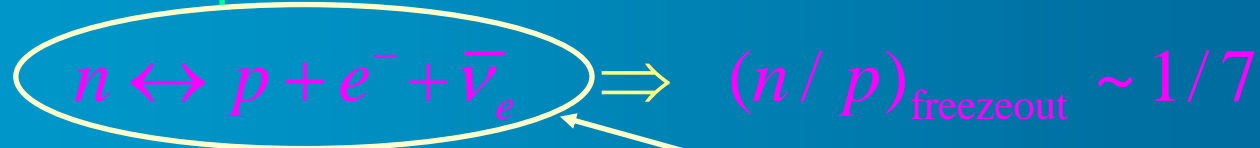
Abstract

- The Standard BBN (SBBN) **approximately** agrees with observations.
- In SUSY/SUGRA cosmology, reheating temperature after Inflation should be low, in order to solve the "**gravitino problem.**"
- We may solve the Lithium problems in SUSY cosmology scenarios
- Dark matter may be neutralinos or gravitinos which are **nonthermally** produced by decaying particles.

Big-Bang Nucleosynthesis (BBN)

- In the early universe at $T > 0.1$ MeV

Neutron to proton ratio was fixed



- Then, $T < 0.1$ MeV

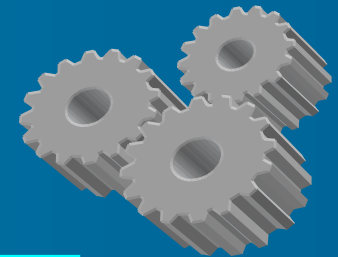
Only through weak interaction



Light element abundances



Baryon to photon ratio $\eta = n_B / n_\gamma$



Time evolution of light elements

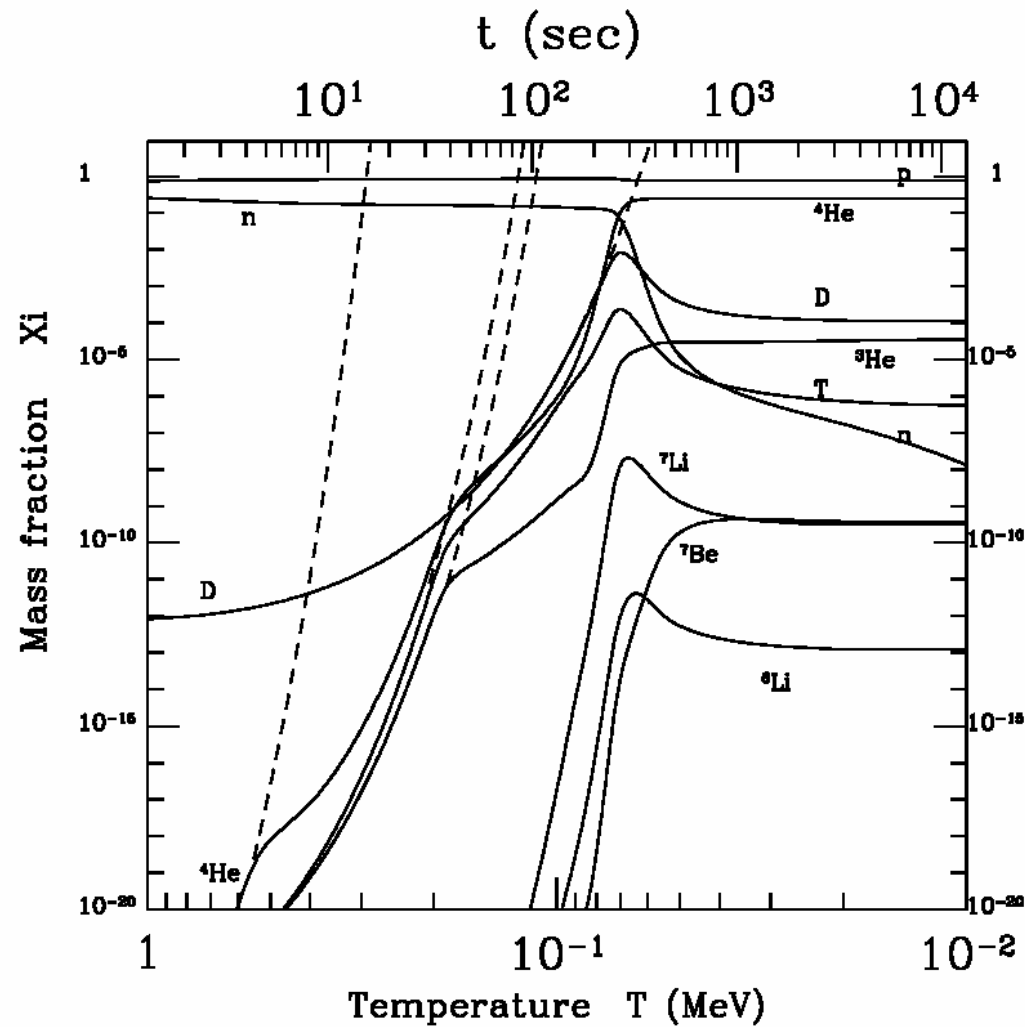


FIG. 2. Time evolution of the light element abundances at $\eta = 3.162 \times 10^{-10}$. The dashed line denotes the nuclear statistical equilibrium value of each light element.

Observational Light Element Abundances



● He4 $Y_p = 0.2516 \pm 0.004$

Fukugita, Kawasaki (2006)

Peimbert, Lridiana, Peimbert (2007)

Izotov, Thuan, Stasinska (2007)

● D/H $D/H = (2.82 \pm 0.26) \times 10^{-5}$

O'Meara et al. (2006)

● Li7/H $\log_{10} ({}^7\text{Li}/\text{H}) = -9.63 \pm 0.06 (\pm 0.3)_{\text{sys}}$

Melendez, Ramirez (2004)

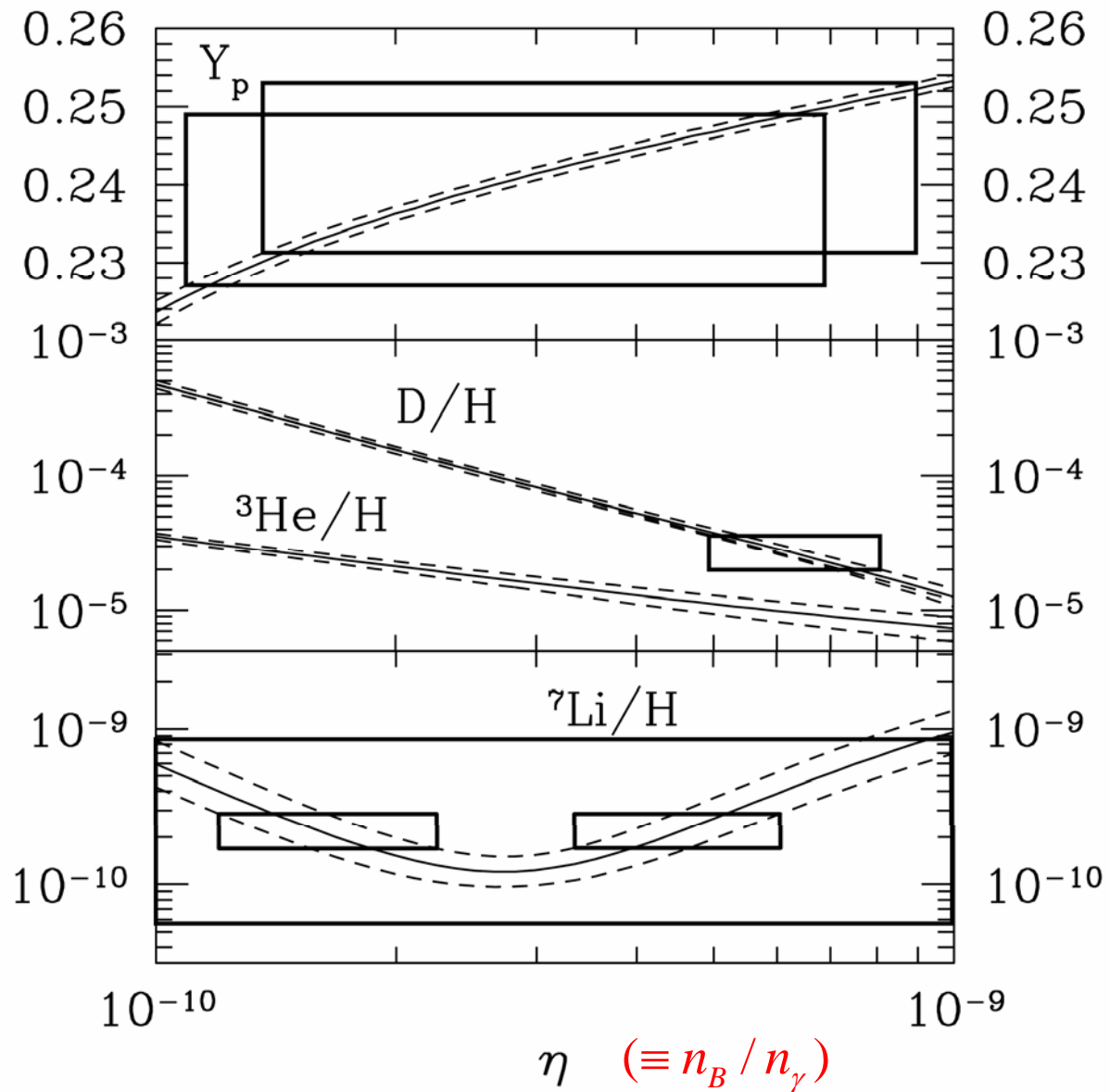
● Li6/H ${}^6\text{Li}/{}^7\text{Li} < 0.046 \pm 0.022 \pm 0.084$

Asplund et al (2006)

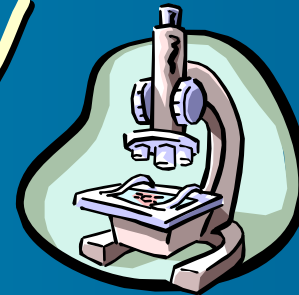
● He3/D ${}^3\text{He}/\text{D} < 0.83 + 0.27$

Geiss and Gloeckler (2003)

SBBN



Introduction of SUSY



✚ Supersymmetry (SUSY)

- Solving "Hierarchy Problem"
- Realizing "Coupling constant unification in GUT"

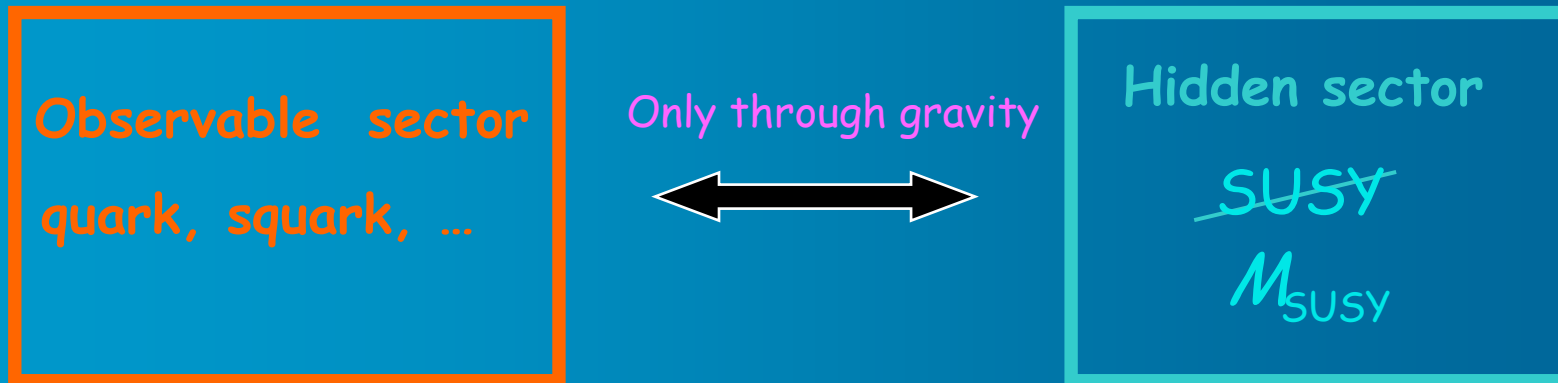


Gravitino ψ_μ is the superpartner of graviton

Spin 3/2

SUSY Breaking

◆ Gravity mediated SUSY breaking model



● Masses of squarks and sleptons

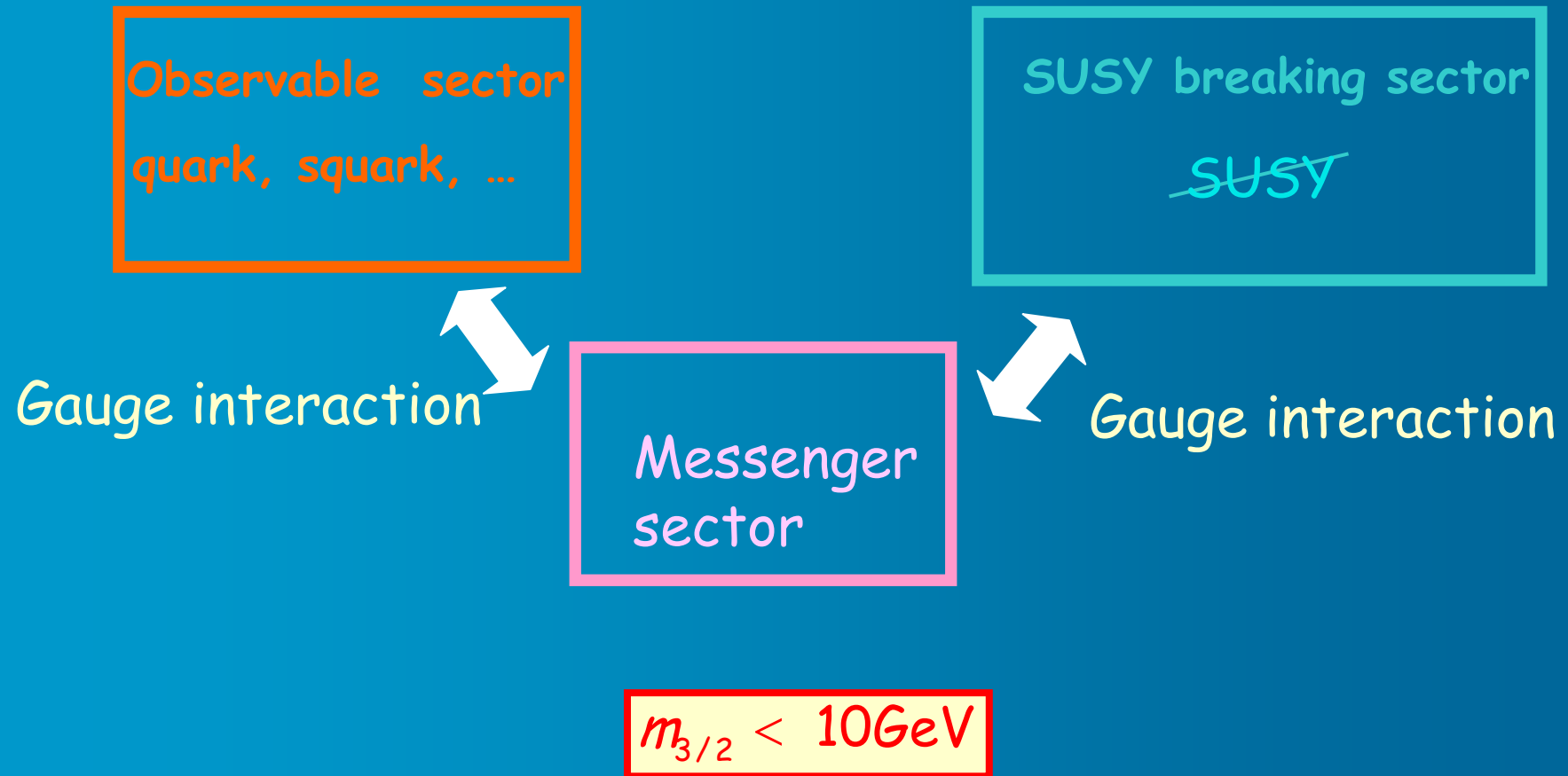
$$m_{\tilde{q}}, m_{\tilde{l}} = M_{\text{SUSY}}^2 / M_{\text{pl}} = 10^2 - 10^3 \text{ GeV}$$
$$(M_{\text{SUSY}} = 10^{10} - 10^{11} \text{ GeV})$$

● Gravitino mass

$$m_{3/2} = M_{\text{SUSY}}^2 / M_{\text{pl}} = 10^2 - 10^3 \text{ GeV}$$

SUSY Breaking II

- ◆ Gauge-mediated SUSY breaking model
(Dynamical SUSY brasking)



Gravitino problem (old version -1)



Cosmological gravitino problem

If gravitinos were abundant as photons,

$$\frac{n_{3/2}}{n_\gamma} \sim \mathcal{O}(1)$$

1) When gravitino decays after BBN epoch ($\psi_\mu \rightarrow \gamma + \tilde{\gamma}$)

$$\tau_{3/2} \simeq 4 \times 10^8 \text{ sec} \left(\frac{m_{3/2}}{100 \text{ GeV}} \right)^{-3}$$

← Long lifetime
only through gravity

Then,

$$\frac{n_B}{n_\gamma} \ll 10^{-10}$$

⇓ to avoid the problem

$$m_{3/2} \gtrsim 100 \text{ TeV}$$

(Weinberg 82)



Gravitino problem (old version -2)

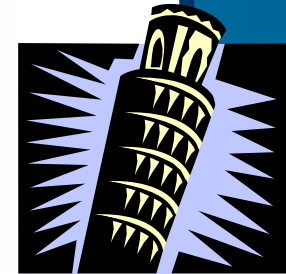
2) If gravitino is stable, $\Omega_{3/2} \lesssim 1$



$$m_{3/2} \lesssim 1 \text{keV}$$

(Pagel, Primack 82)

Gravitino with $m_{3/2} \sim 100 \text{GeV}$ is exclude ?



No. Inflation dilutes gravitinos!

Gravitino production after inflation

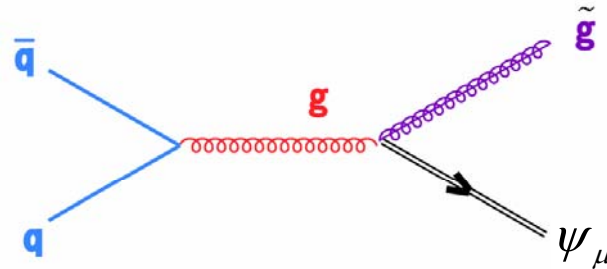


I) Inflation dilutes primordial gravitinos

$$n_{3/2} / s \rightarrow 0$$

II) Reheating process produces gravitinos again

e. g., $q + \bar{q} \rightarrow \psi_{3/2} + \tilde{g}$



$$Y_{3/2} \equiv \frac{n_{3/2}}{n_\gamma} \simeq 10^{-11} \left(\frac{T_R}{10^{10} \text{GeV}} \right)^1$$

Because,

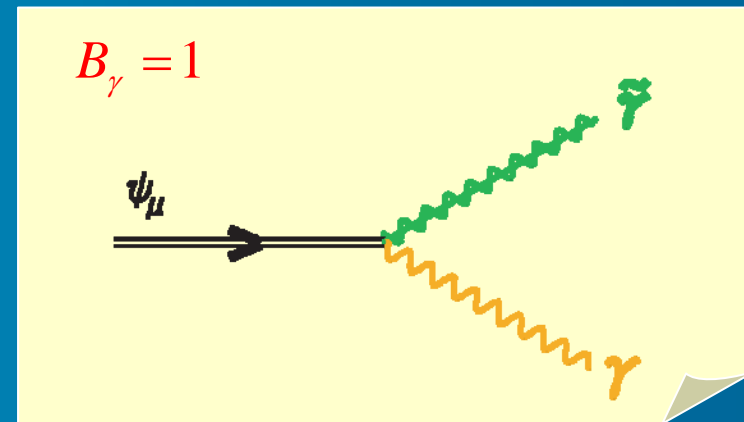
$$\begin{aligned} \frac{\Delta n_{3/2}}{n_\gamma} &\sim \Gamma \Delta t \\ &\sim \frac{1}{M_G^2} T_R^3 \times \frac{M_G}{T_R^2} \\ &\propto T_R \end{aligned}$$

Gravitino Decay and BBN

1. Gravitinos are unstable in Gravity Mediation ~~SUSY~~

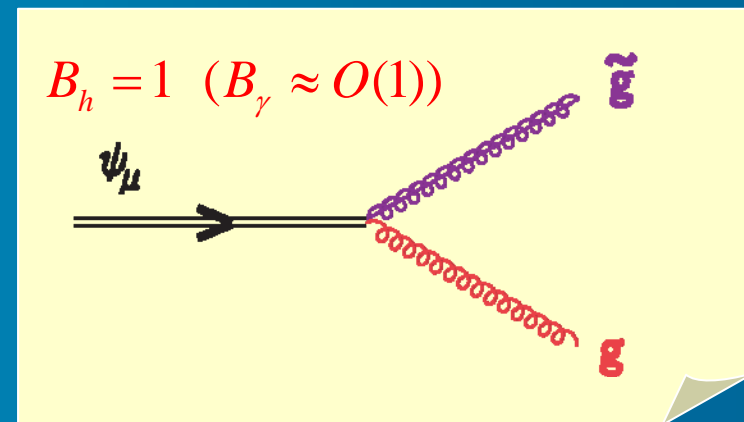
• Radiative decay

$$\tau(\psi_{3/2} \rightarrow \gamma + \tilde{\gamma}) = 4 \times 10^8 \text{ sec} \left(\frac{m_{3/2}}{100 \text{ GeV}} \right)^{-3}$$



• Hadronic decay

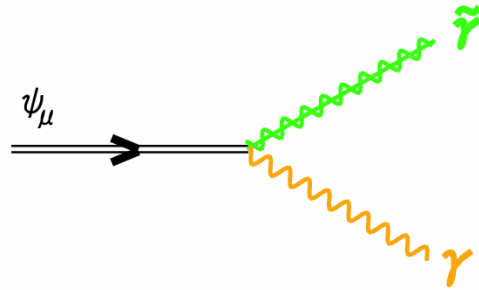
$$\tau(\psi_{3/2} \rightarrow g + \tilde{g}) = 6 \times 10^7 \text{ sec} \left(\frac{m_{3/2}}{100 \text{ GeV}} \right)^{-3}$$



Radiative Decay



Radiative decay of gravitino



1) Electro-magnetic cascade

$$\gamma + \gamma_{\text{BG}} \rightarrow e^+ + e^-$$

$$\gamma + e_{\text{BG}}^- \rightarrow \gamma + e^-, \quad e^- + \gamma_{\text{BG}} \rightarrow e^- + \gamma$$

$$\gamma + \gamma_{\text{BG}} \rightarrow \gamma + \gamma$$

2) many soft photons are produced

3) Photo-dissociation of light elements

$$\text{D} + \gamma \rightarrow p + n,$$

$${}^4\text{He} + \gamma \rightarrow {}^3\text{He} + n, \quad \text{T} + p,$$

$${}^3\text{He} + \gamma \rightarrow \text{D} + p + n, \quad \text{etc.}$$

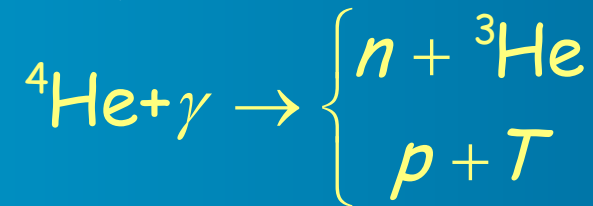
Non-thermal Li6 Production

Dimopoulos et al (1989)

Jedamzik (2000)

Kawasaki, Kohri, Moroi (2001)

i. He3 (T) production



ii. Li6 production

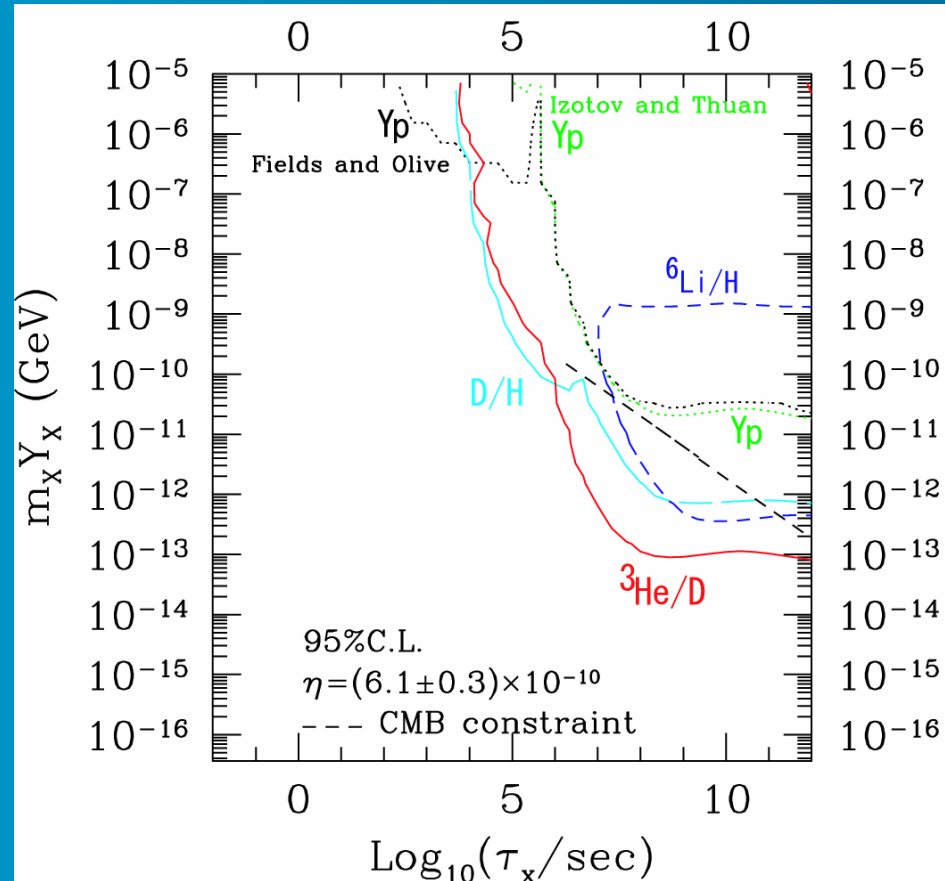


Coulomb energy loss by background electrons

$$\frac{dE}{dx} = \frac{4\pi\alpha^2 Z^2 \Lambda n_e}{m_e \beta}, \quad \Lambda \sim O(1)$$

Photodissociation

$$Y_{3/2} \equiv \frac{n_{3/2}}{n_\gamma}$$



Kawasaki, Kohri, Moroi (2001)

Relation among variables

- Yield variable and reheating temperature

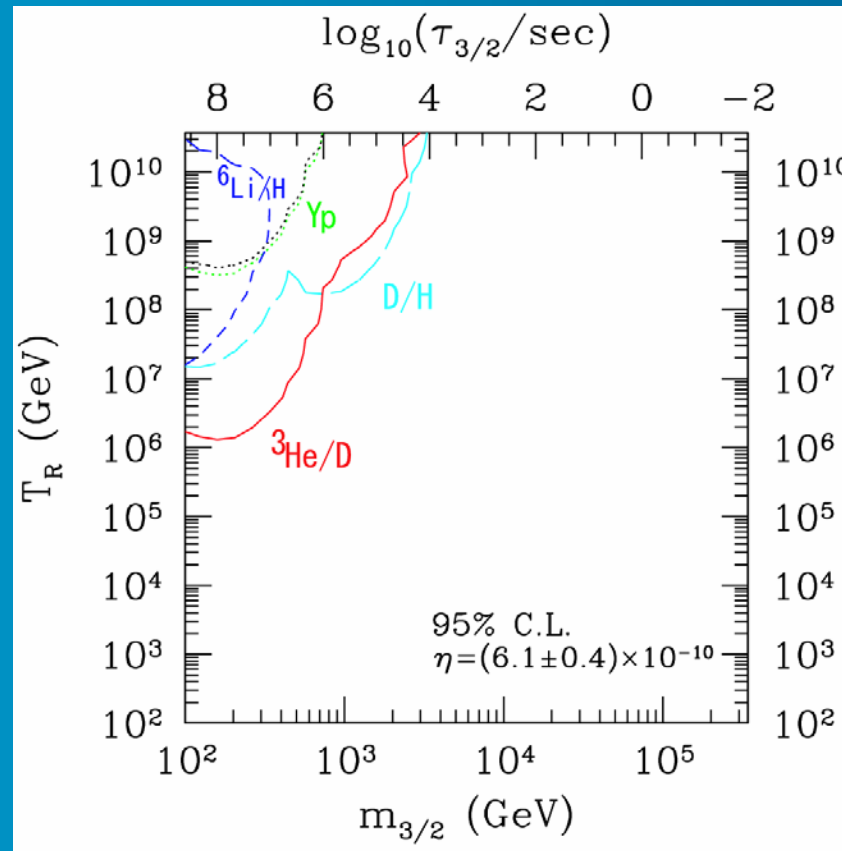
$$y_{3/2} \equiv \frac{n_{3/2}}{n_\gamma} = 1.1 \times 10^{-11} \left(\frac{T_R}{10^{10} \text{ GeV}} \right)$$

- Lifetime and mass

$$\tau(\psi_{3/2} \rightarrow \gamma + \tilde{\gamma}) = 4 \times 10^8 \text{ sec} \left(\frac{m_{3/2}}{100 \text{ GeV}} \right)^{-3}$$

Upper bound on reheating temperature

Kawasaki, Kohri, Moroi (2001)



$$T_R = 10^9 \text{ GeV} (y_{3/2} / 10^{-12})$$

$$m_{3/2} = 10^3 \text{ GeV} (\tau_{3/2} / 4 \times 10^5 \text{ sec})^{-1/3}$$



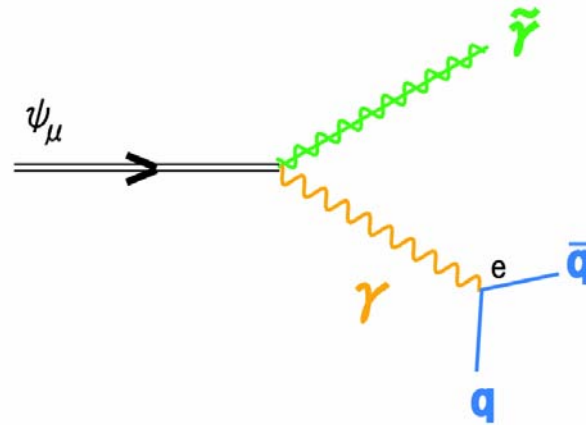
Hadronic Decay



Hadronic decay

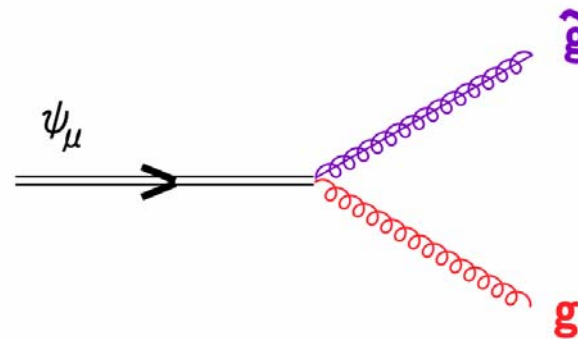
Reno, Seckel (1988)

S. Dimopoulos et al.(1989)



Two hadron jets with
 $E_{\text{jet}} = m_X / 3$

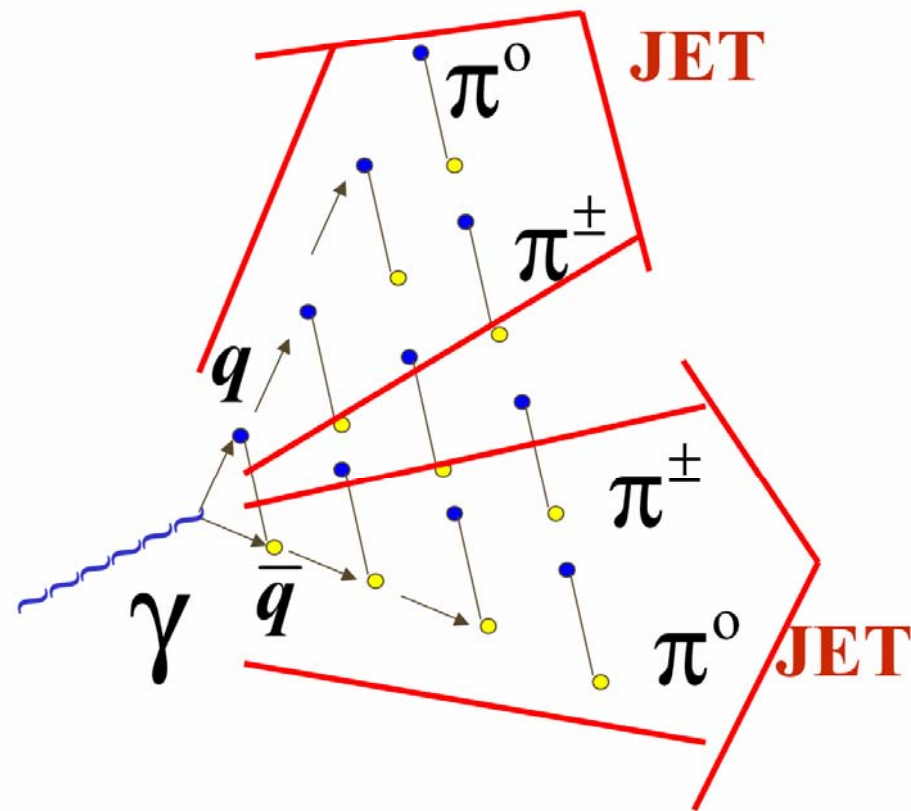
$$B_h \approx \alpha / 4\pi \approx 10^{-3}$$



One hadron jet with
 $E_{\text{jet}} = m_X / 2$

$$B_h = 1$$

Quark Emission and Hadron Jets



$$q = u, d, s, \dots$$

Hadron-fragmentation Monte Carlo event generator, JETSET 7.4 (PYTHIA 5.7) Sjostrand (1994)

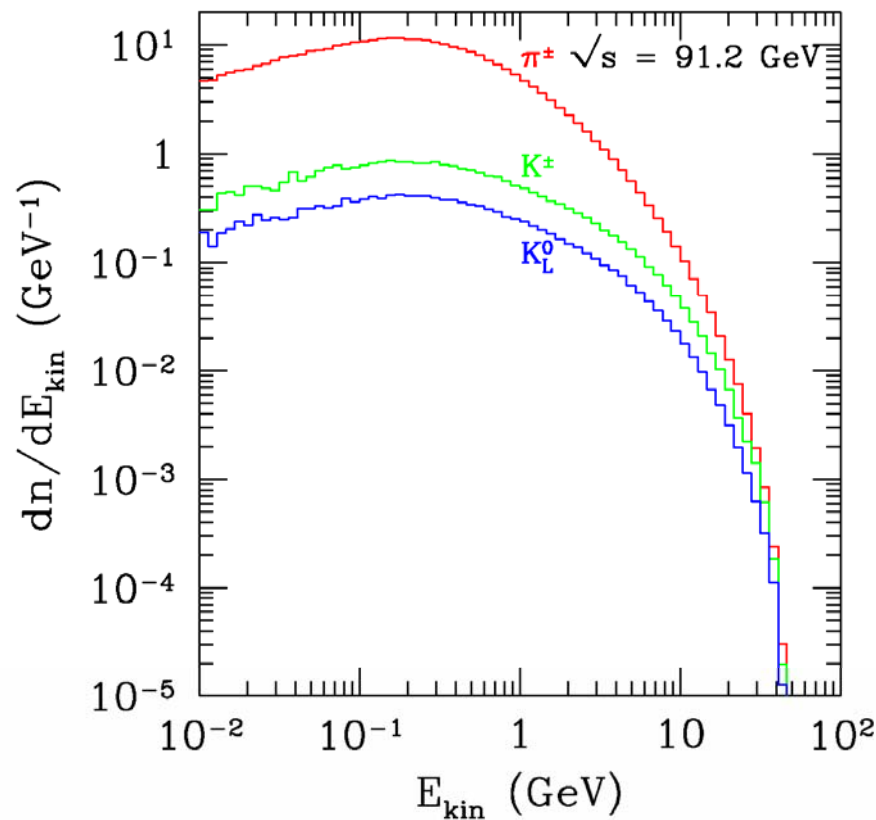
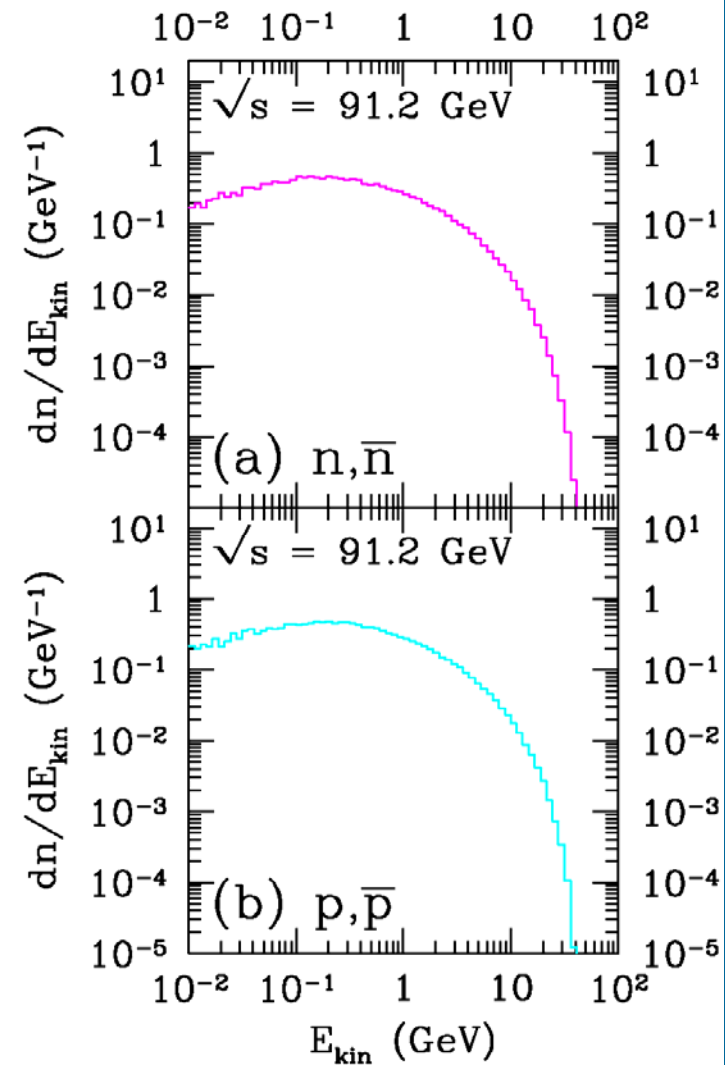


FIG. 3. Plot of the spectrum of the produced mesons ($\pi^+ + \pi^-$, $K^+ + K^-$, and K_L^0) as a function of the kinetic energy E_{kin} . This is the case that the center of mass energy is $\sqrt{s} = 91.2$ GeV which corresponds to the Z^0 resonance. They are computed by using the JETSET 7.4 Monte Carlo event generator.



Kohri (2001)

Electromagnetic stopping of emitted hadrons?

There is a question if high energy hadrons are stopped, until they scatter off the ambient nuclei in the electromagnetic plasma of e^\pm and γ .

$$R_{\text{stop}}(E_i, T, Y_p) \equiv n_N \int_{E_i}^{E_{\text{th}}} \langle \sigma \beta \rangle \left(\frac{dE}{dt} \right)^{-1} dE$$

where,

(I) If $R_{\text{stop}}(E_i, T, Y_p) \ll 1$,

$$\left(\frac{dE}{dt} \right) = \left(\frac{dE}{dt} \right)_{\text{Coulomb}} + \left(\frac{dE}{dt} \right)_{\text{Compton}} + \dots$$

Hadrons are stopped electromagnetically

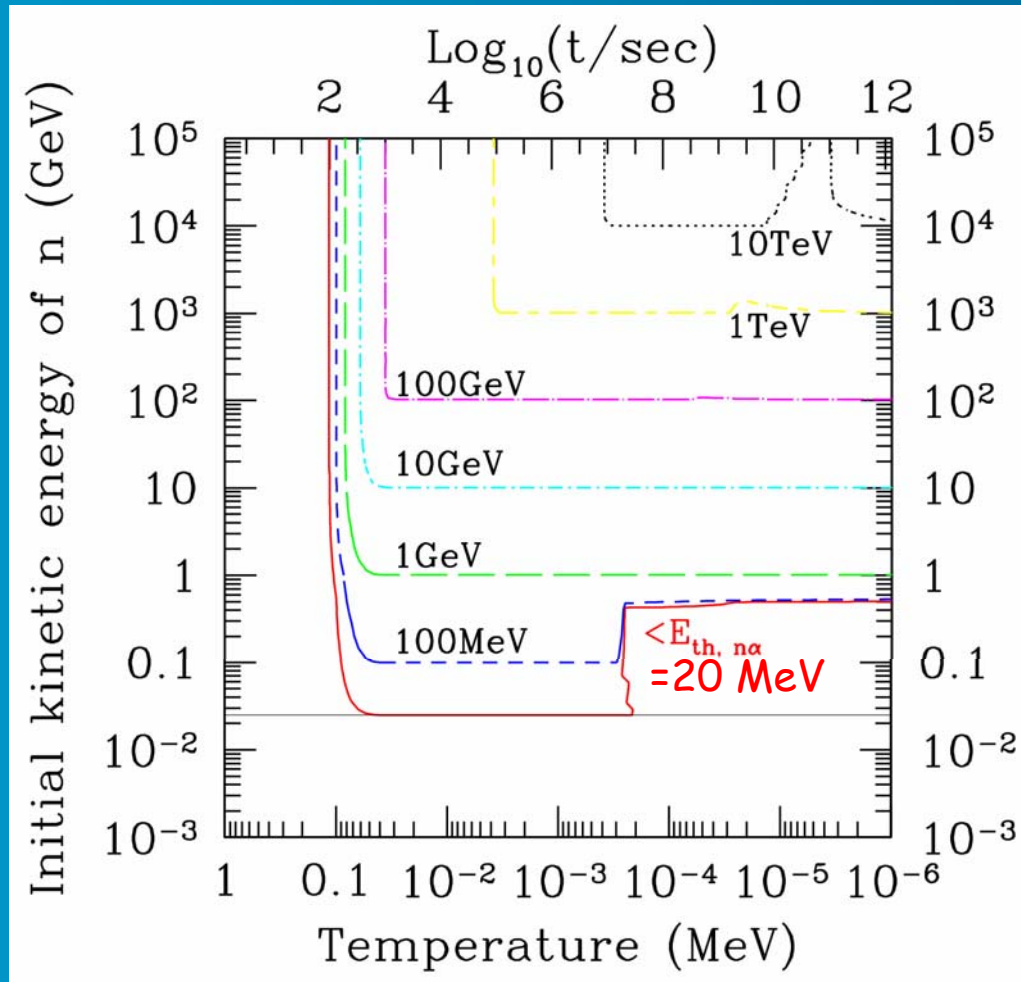
(II) If $R_{\text{stop}}(E_i, T, Y_p) \gg 1$,

Hadrons scatter off the background nuclei with the energy

E_f which is obtained by

$$n_N \int_{E_i}^{E_f} \langle \sigma \beta \rangle \left(\frac{dE}{dt} \right)^{-1} dE = 1$$

Contours of final energy of energetic neutron just before its collision



(I) If $R_{\text{stop}}(E_i, T, Y_p) \ll 1$,

Inter-conversion occurs between n and p by stopped hadrons

Effective at early stage of BBN

Reno and Seckel (1988)

Kohri (2001)

(II) If $R_{\text{stop}}(E_i, T, Y_p) \gg 1$,

Destruction of Helium4 occurs by energetic hadrons

Effective at late stage of BBN

S. Dimopoulos et al(1989)

Kawasaki, Kohri, Moroi(2004)

(I) Early stage of BBN (before/during BBN)

Reno and Seckel (1988) Kohri (2001)

Extraordinary inter-conversion reactions between n and p



$$\Gamma_{n \leftrightarrow p} = \Gamma_{n \leftrightarrow p}^{\text{weak}} + \Gamma_{n \leftrightarrow p}^{\text{strong}}$$

Hadron induced exchange

$$\Gamma_{n \leftrightarrow p} \uparrow \Rightarrow n/p \uparrow$$

Even after freeze-out of n/p in SBBN



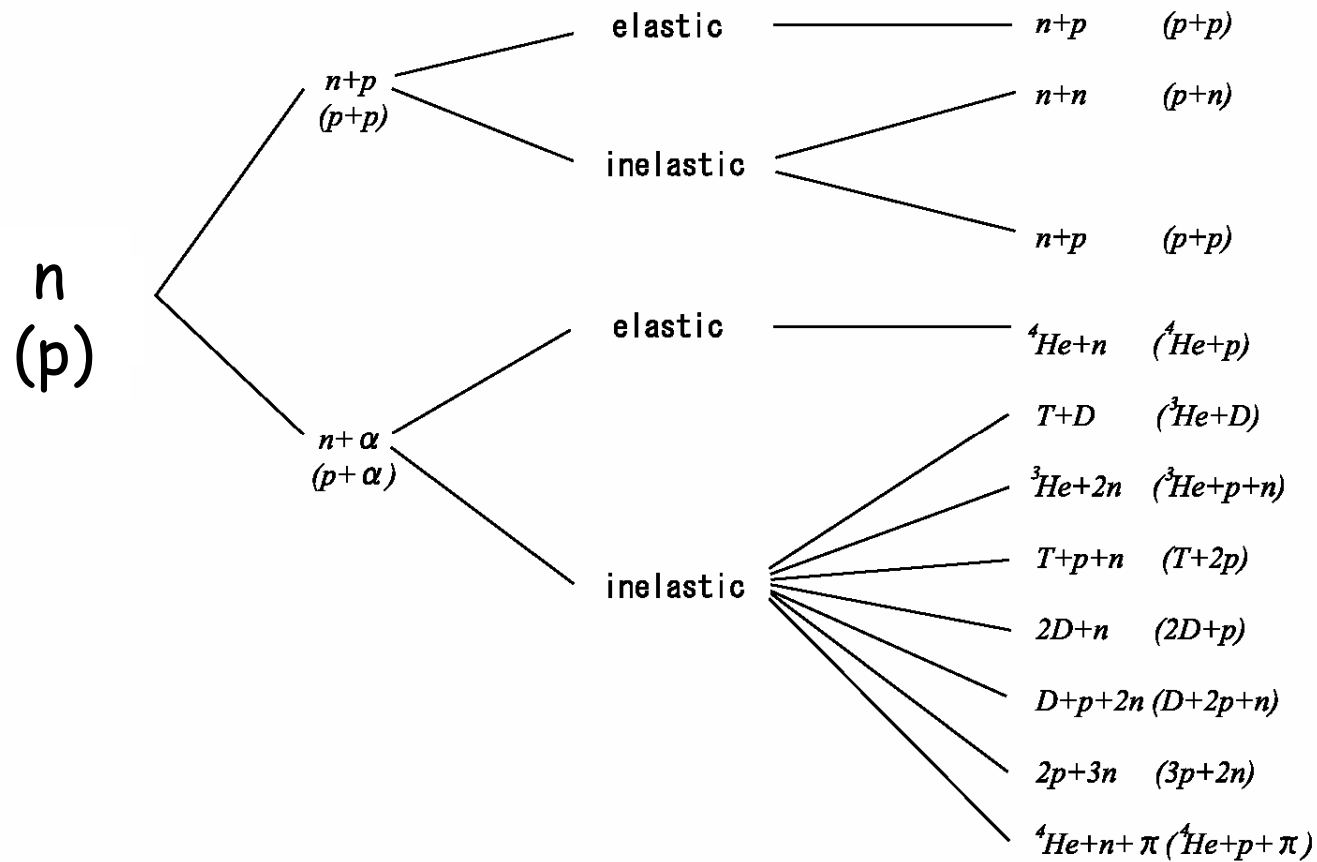
More He4, D, Li7 ...

(II) Late stage of BBN

Hadronic showers and "Hadro-dissociation"

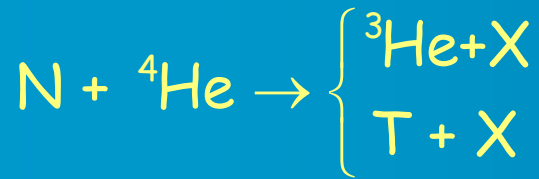
S. Dimopoulos et al. (1988)

Kawasaki, Kohri, Moroi (2004)

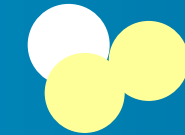


Non-thermal Li, Be Production by energetic hadrons

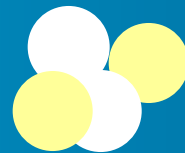
Dimopoulos et al (1989)



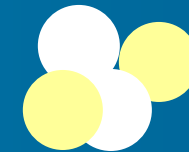
T, ${}^3\text{He}$



${}^4\text{He}$



Energy loss



${}^4\text{He}$

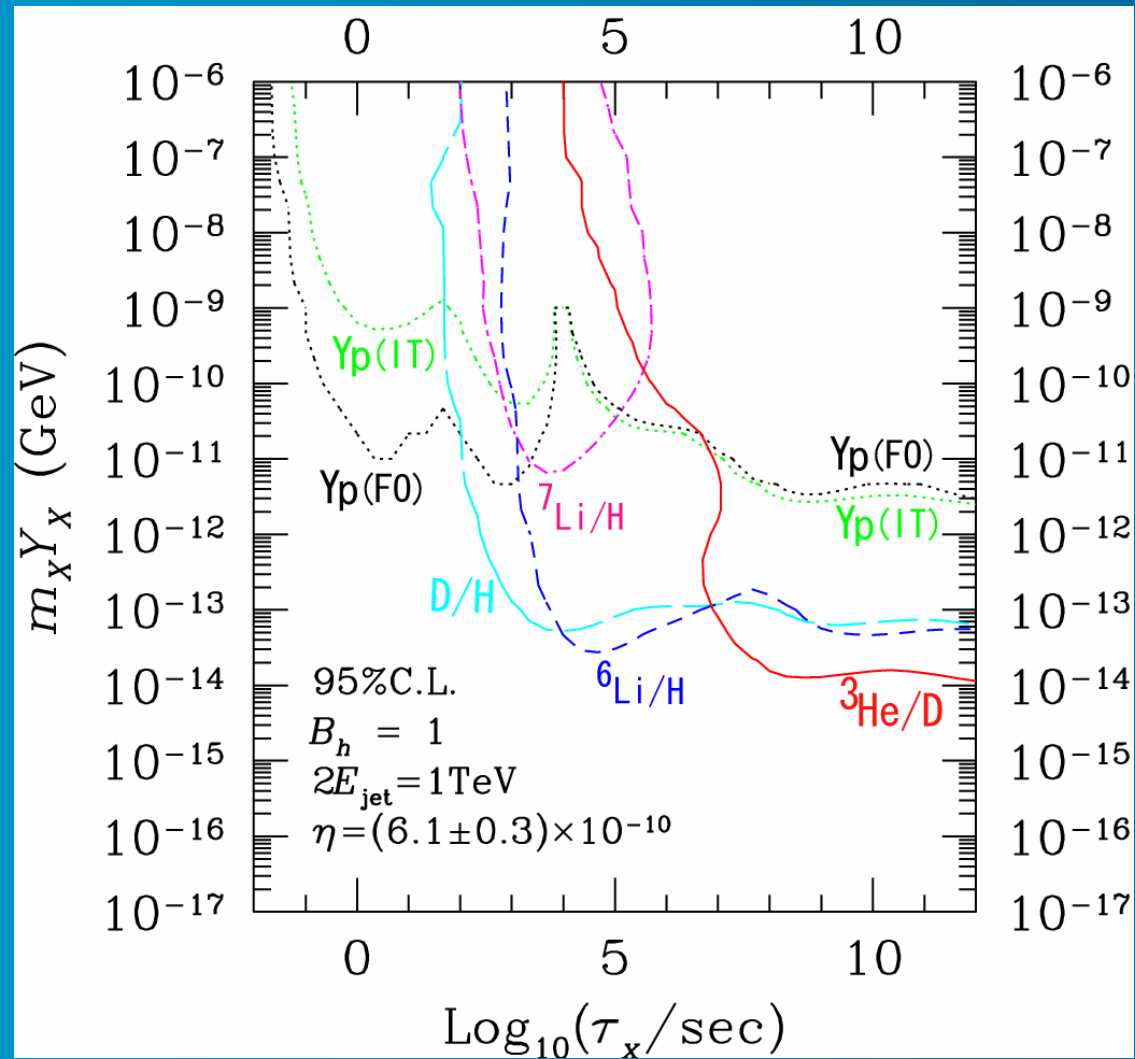
① T(He3) - He4 collision



② He4 - He4 collision



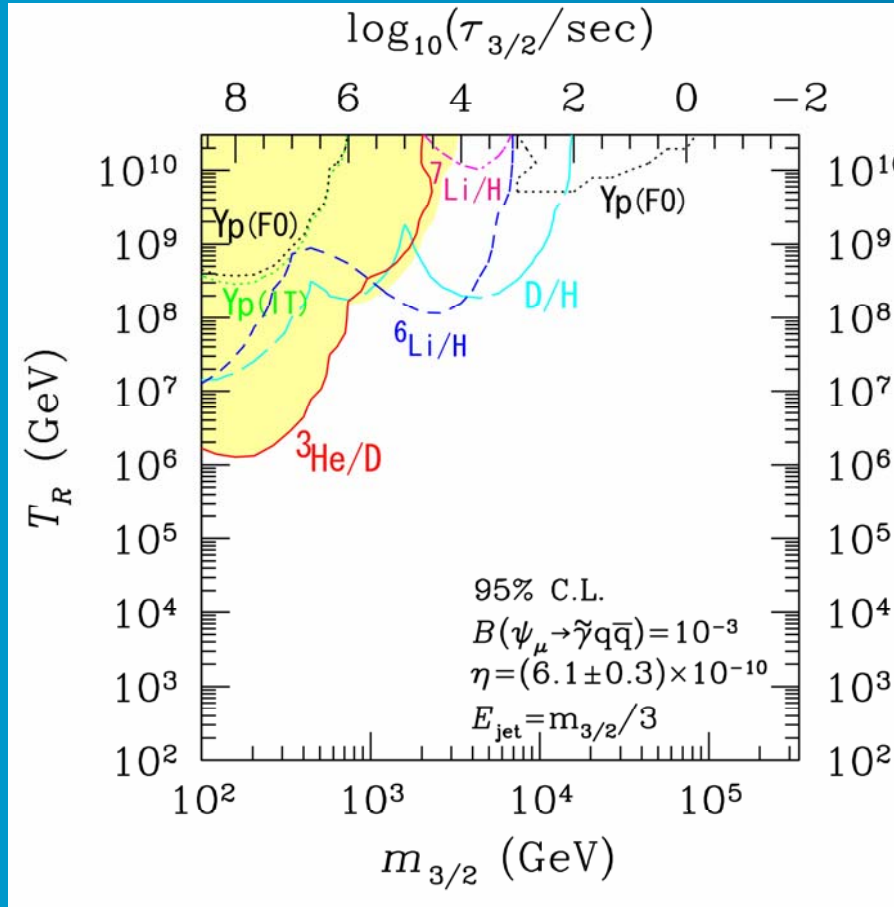
For massive particle X



Contours of light elements in $(m_x Y_x, \tau_x)$ plane
in "hadrodissociation" scenario

Upper bound on reheating temperature

Kawasaki, Kohri, Moroi (2004)



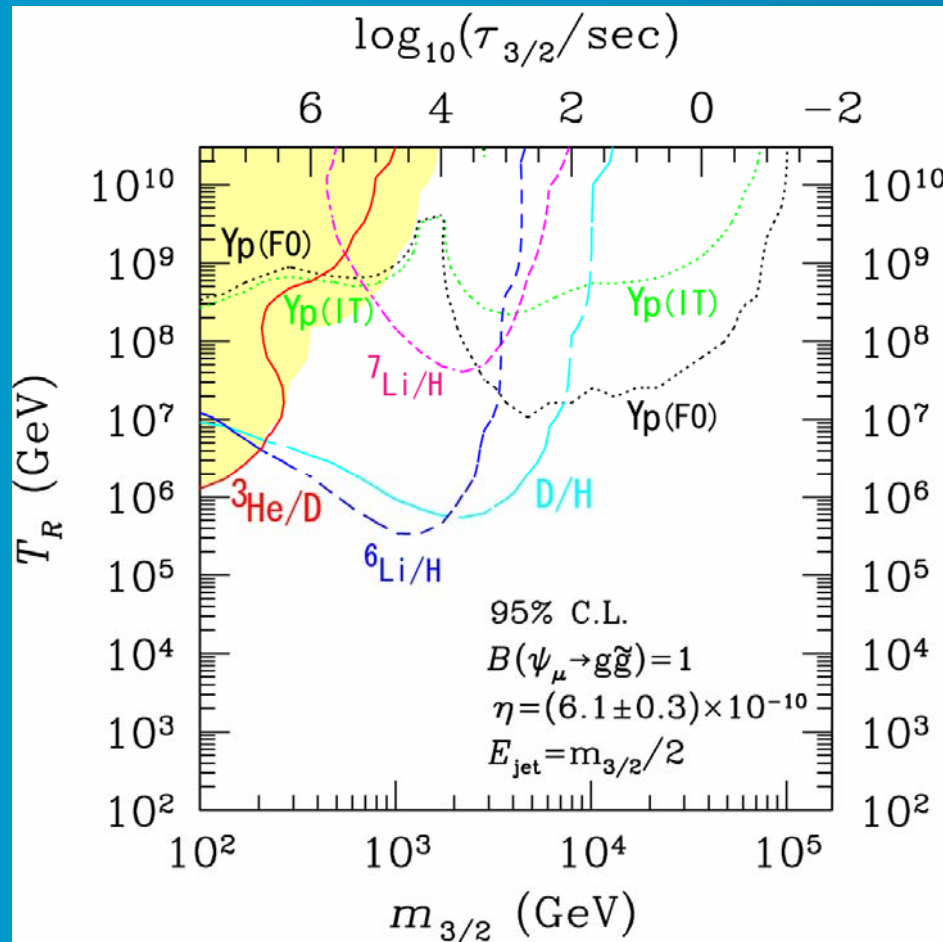
$$B_h(\psi_{\mu} \rightarrow \tilde{\gamma} + q + \bar{q}) = 10^{-3}$$

$$T_R = 10^9 \text{ GeV} (y_{3/2} / 10^{-12})$$

$$m_{3/2} = 10^3 \text{ GeV} (\tau_{3/2} / 4 \times 10^5 \text{ sec})^{-1/3}$$

Upper bound on reheating temperature

Kawasaki, Kohri, Moroi (2004)

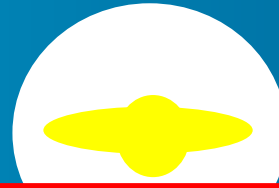


$$B_h(\psi_\mu \rightarrow g + \tilde{g}) = 1$$

$$T_R = 10^9 \text{ GeV} (y_{3/2} / 10^{-12})$$

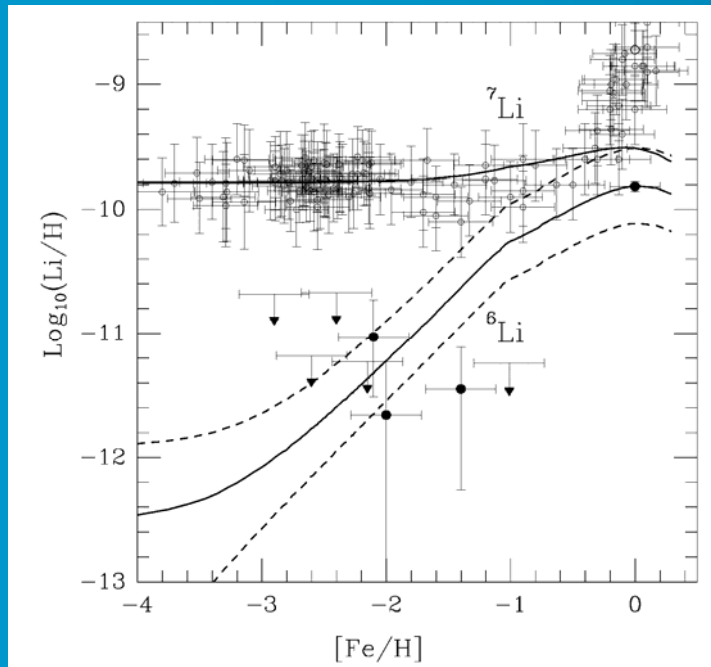
$$m_{3/2} = 500 \text{ GeV} (\tau_{3/2} / 4 \times 10^5 \text{ sec})^{-1/3}$$

Lithium 7



a factor of two or three smaller !!!

- Expected that there is little depletion in stars.



Lemoine et al., 1997

$${}^7\text{Li}/\text{H} = 2.19^{+2.2}_{-1.1} \times 10^{-10} \quad (1\sigma)$$

Bonifacio et al.(2002)

Melendez,Ramirez(2004)

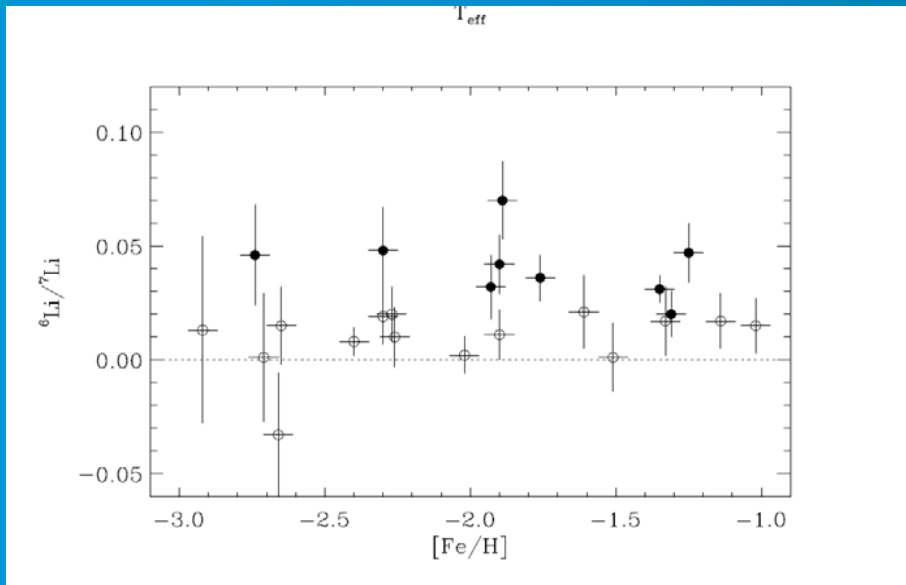
$${}^7\text{Li}/\text{H} = 1.23^{+0.68}_{-0.32} \times 10^{-10} \quad (1\sigma)$$

Ryan et al.(2000)

Lithium 6

Asplund et al.(2006)

- Observed in metal poor halo stars in Pop II
- ${}^6\text{Li}$ plateau?



$${}^6\text{Li} / {}^7\text{Li} = 0.01 - 0.09$$

${}^7\text{Li}/\text{H} \approx (1.1 - 1.5) \times 10^{-10}$
still disagrees with SBBN

Astrophysically, factor-of-two depletion of $\text{Li}7$ needs a factor of $O(10)$ $\text{Li}6$ depletion (Pinsonneault et al '02)

We need more primordial $\text{Li}6$?

Gravitino Dark Matter Scenario in Gauge Mediated SUSY Breaking

$$m_{3/2} < 10\text{GeV}$$

NLSP would be Slepton (stau, sneutrino)
or Neutralino (Bino)

CHArged Massive Particle (CHAMP)

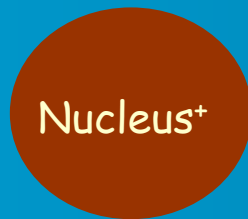
Kohri and Takayama, hep-ph/0605243

Many candidates of long-lived CHAMP

stau, ...



More massive elements capture CHAMP earlier



$$T_c \sim E_{\text{bin}}/40 \sim 10\text{keV}$$

$$(E_{\text{bin}} \sim \alpha^2 m_i \sim 100\text{keV})$$

CHAMP captured-nuclei change the nuclear reaction rates

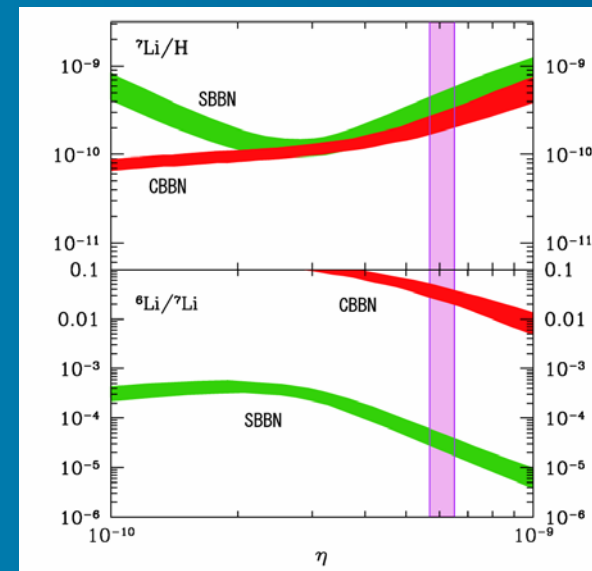
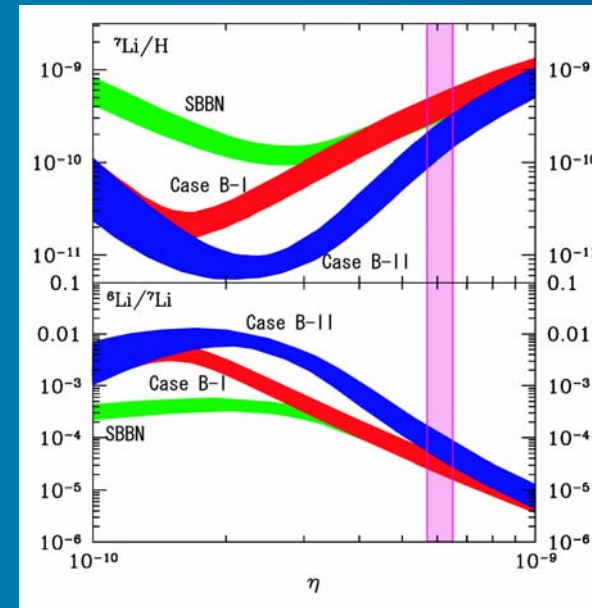
CHAMP BBN (CBBN) may solve Lithium problem?

Short lifetime ($< 10^6$ sec)

- Only $\text{Be}7$ and $\text{Li}7$ captures CHAMP
- $\text{Be}7(n,\alpha)\text{He}4$ and $\text{Li}7(p,\alpha)\text{He}4$ are enhanced

Long lifetime ($> 10^6$ sec)

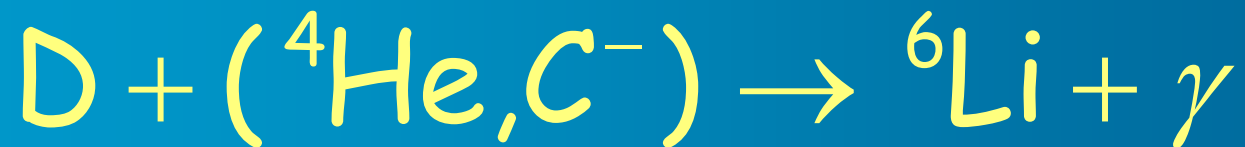
- proton, D, and T are captured
- $\text{He}4(d,g)\text{Li}6$ and $\text{Be}7(d,p)\alpha\text{He}4$ are enhanced



Pospelov's effect

Pospelov (2006), hep-ph/0605215

- CHAMP bound state with ${}^4\text{He}$ can enhance the rate



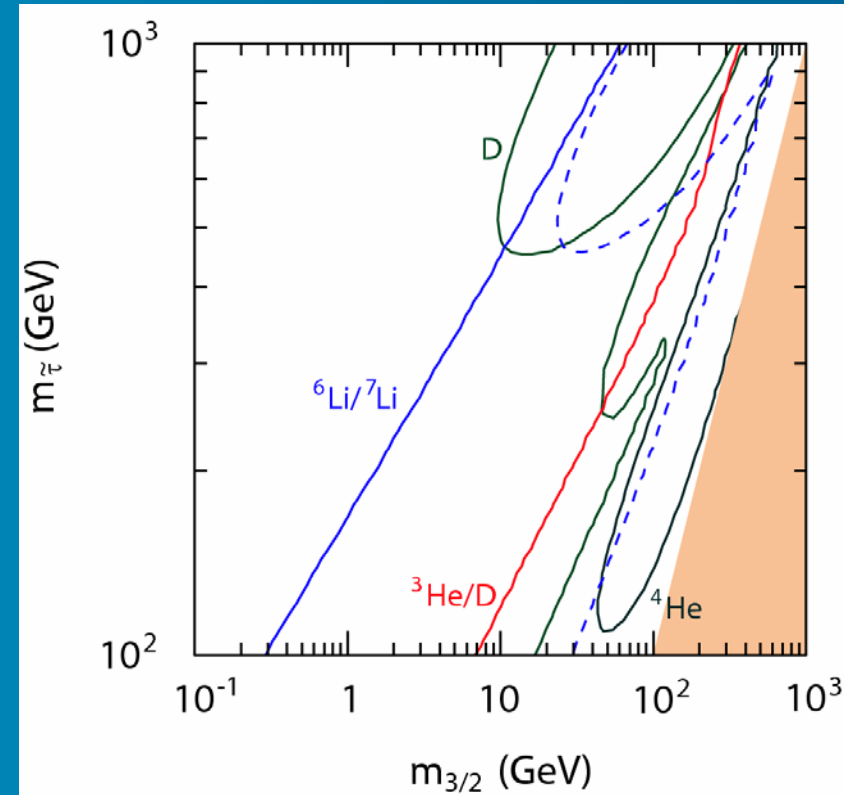
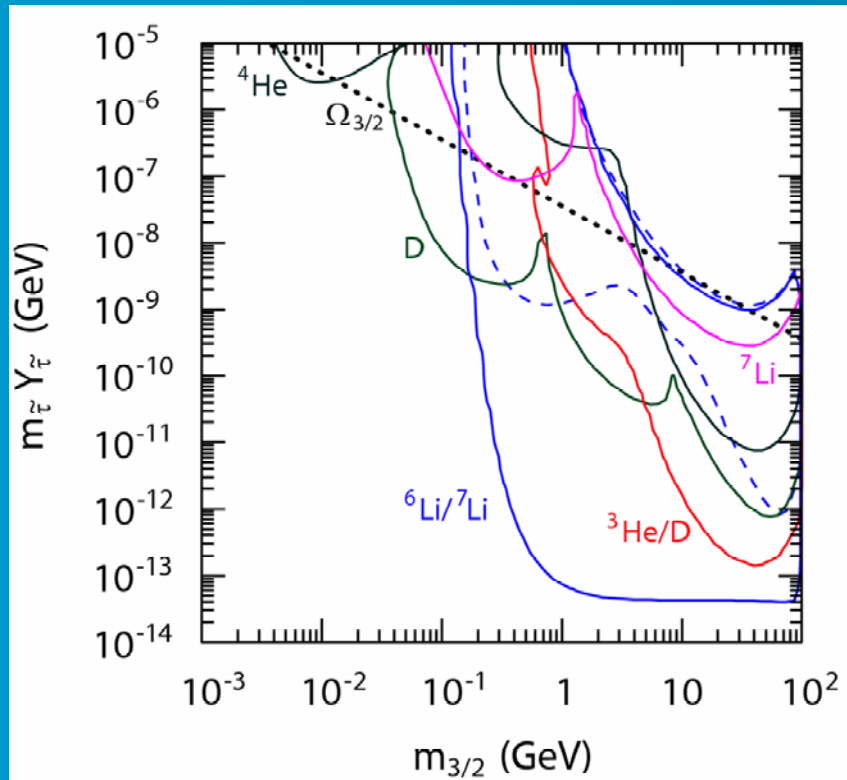
- Enhancement of cross section

$$\sim (\lambda_\gamma / a_{\text{Bohr}})^5 \sim (30)^5 \sim 10^8$$

Confirmed by Hamaguchi et al (07), hep-ph/0702274

BBN in stau NLSP and gravitino LSP Scenario

Kawasaki, Kohri, Moroi (07)



Discussion and Conclusion

- The radiative and hadronic decay-products *destroy* He4 , by which D , He3 , Li6 are overproduced.
- The constraint on reheating temperature after primordial inflation becomes very stringent in Hadronic decay scenario in gravity mediated SUSY breaking scenario.

$$T_R \leq 3 \times 10^5 \text{ GeV} - 10^7 \text{ GeV}$$

$$\text{(for } m_{3/2} = 100 \text{ GeV} - 10 \text{ TeV)}$$

- CHAMP BBN is attractive (Kohri and Takayama '06) or might be dangerous? (Pospelov hep-ph/0605215). Then DM should be a stable gravitino in gauge-mediated SUSY breaking scenario.